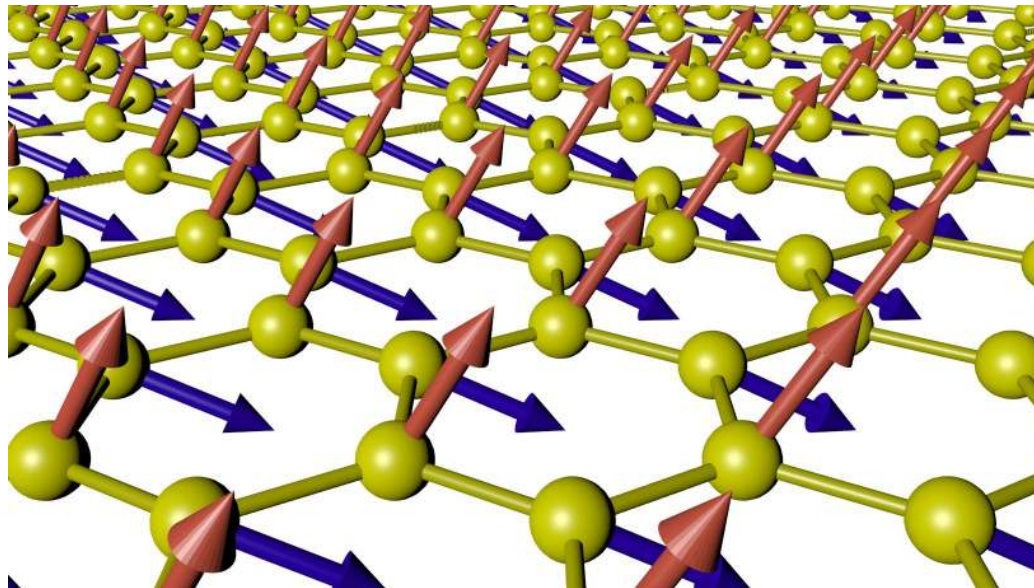


Magnetism and electronics in quantum Hall bars

Jose Lado and Joaquin Fernandez Rossier

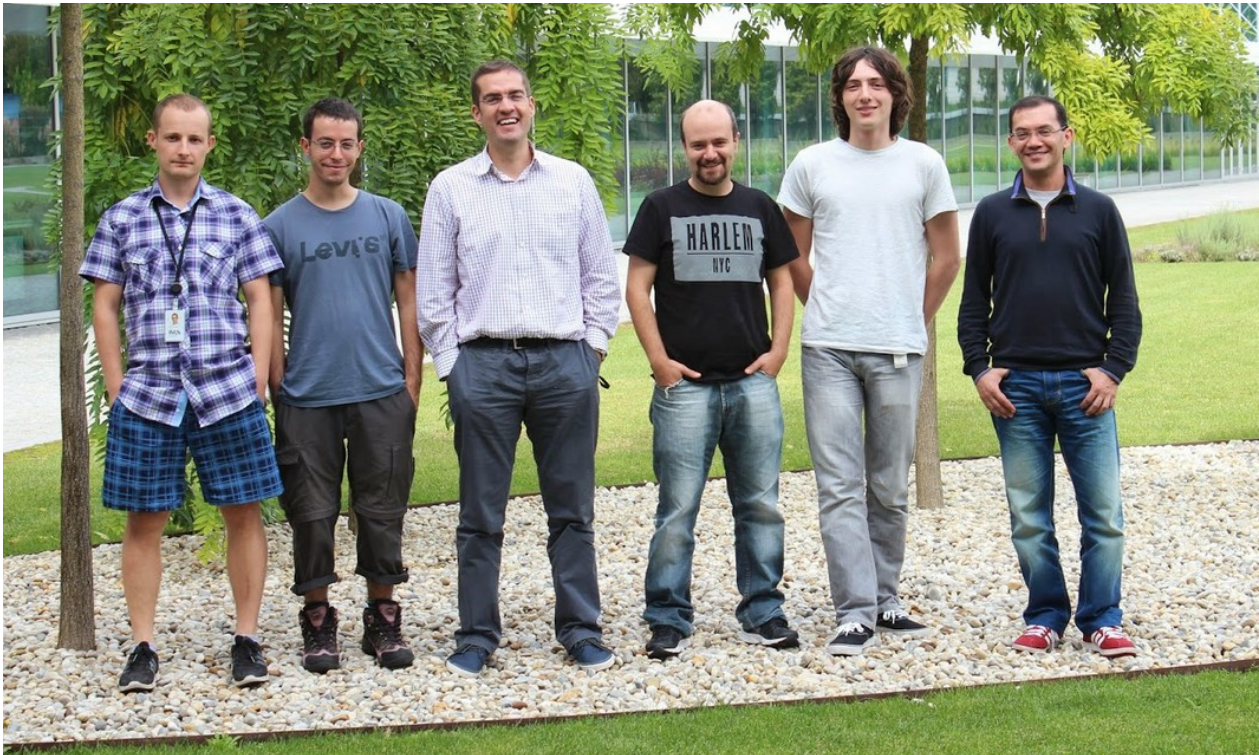
International Nanotechnology Laboratory (INL), Braga, Portugal



Aachen 2014, 9 December

The team

Theory group at INL, Braga, Portugal

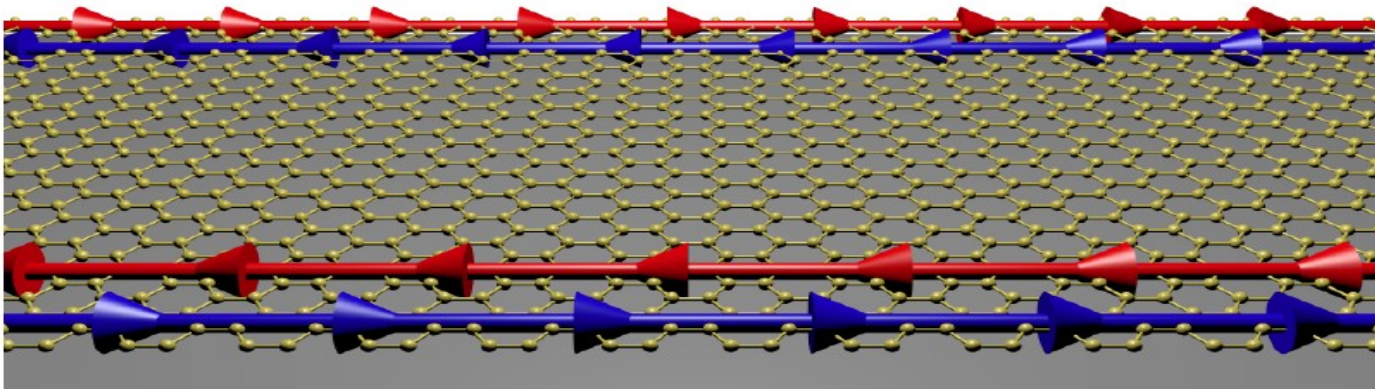
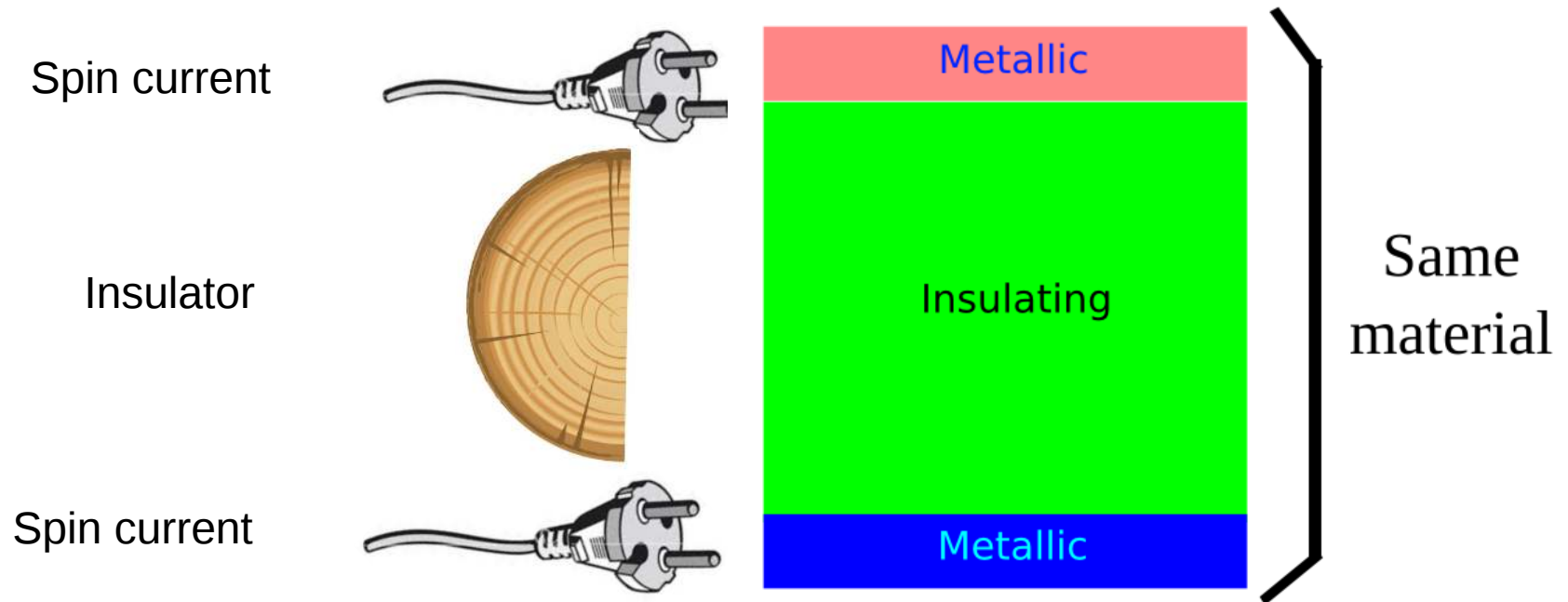


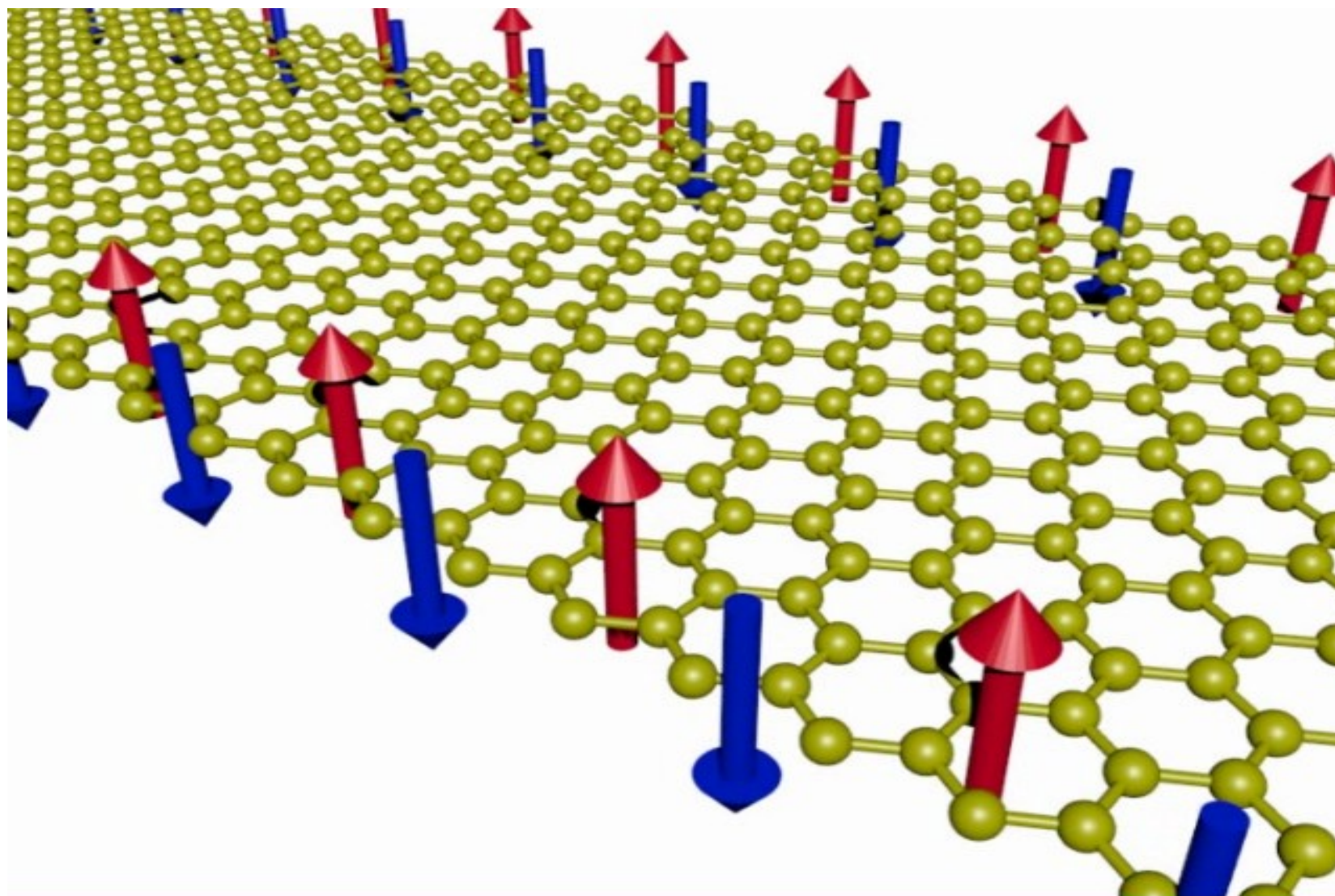
Jose Lado



Joaquin
Rossier

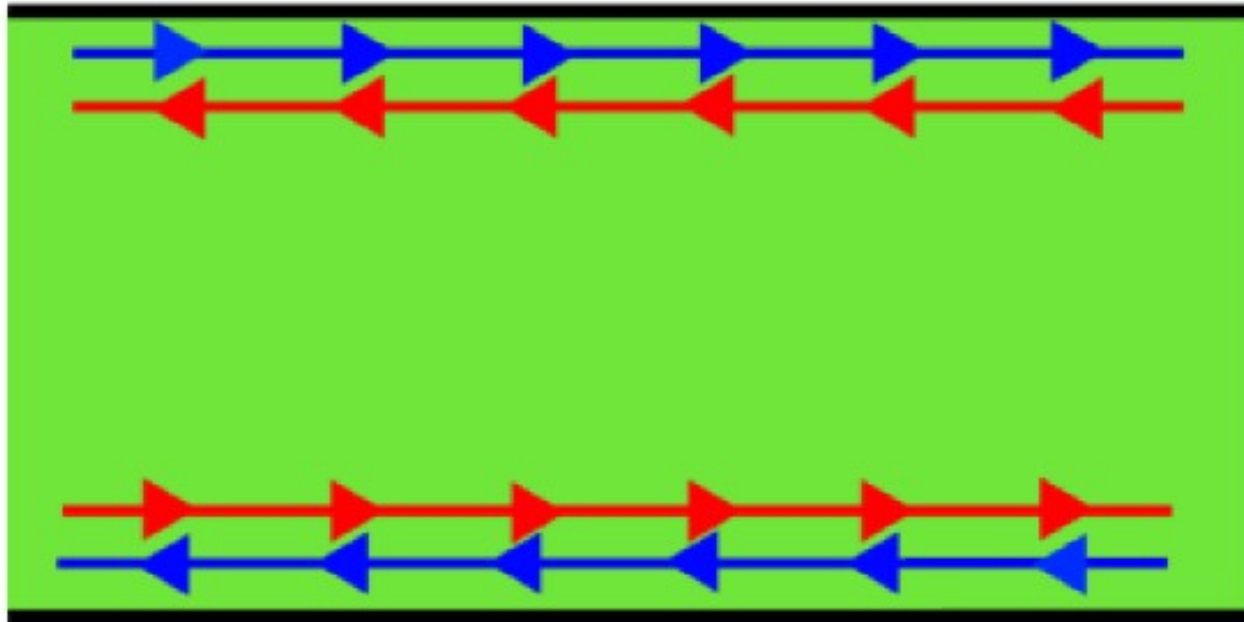
Topological state





What is the quantum spin Hall effect

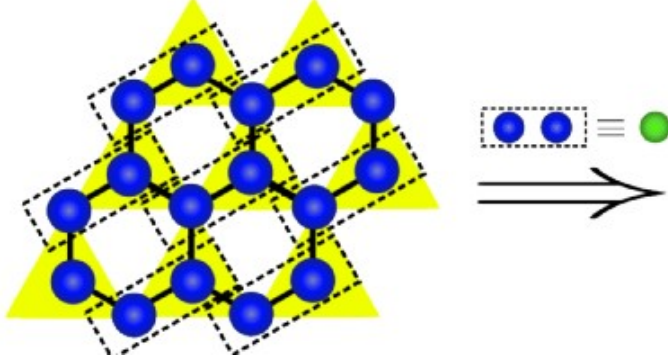
Counterpropagating spin polarized edge states



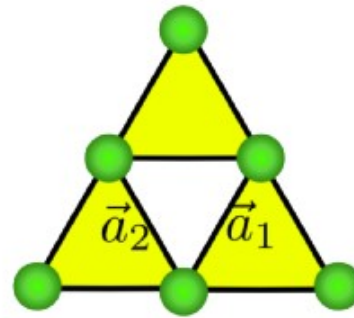
Spin current but no net current

What is graphene

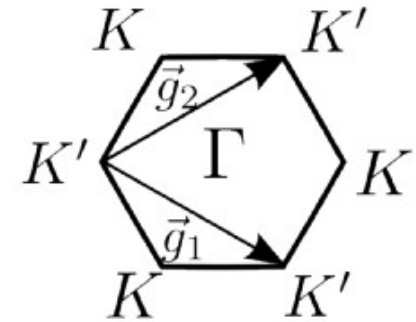
(a) Honeycomb lattice



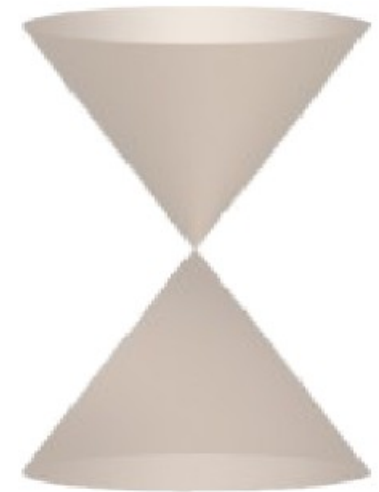
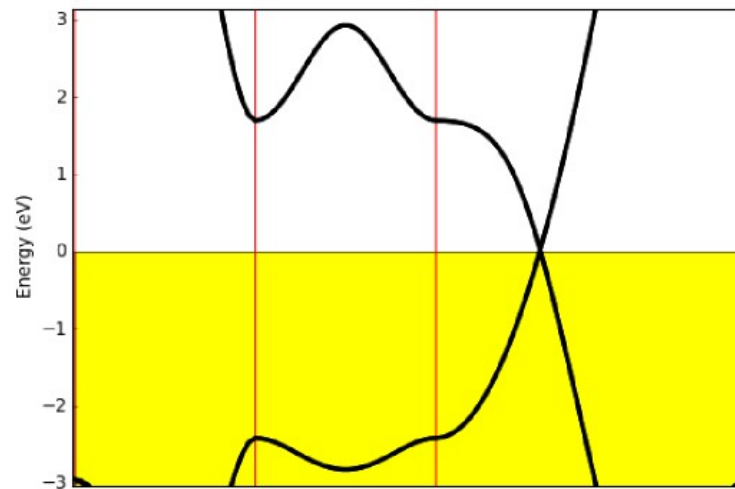
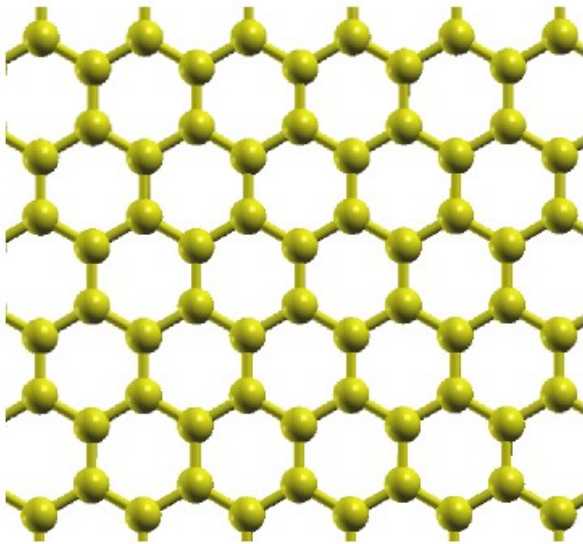
(b) Triangular lattice



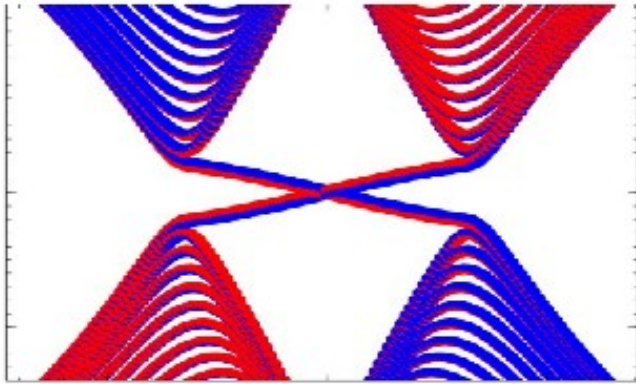
(c) First Brillouin zone



Low energy described by Dirac equation



Quantum Spin Hall in graphene



$\mathbb{Z}_2$ topological insulator

Theory

Kane Mele PRL (2005)



Spin Chern insulator

Theory

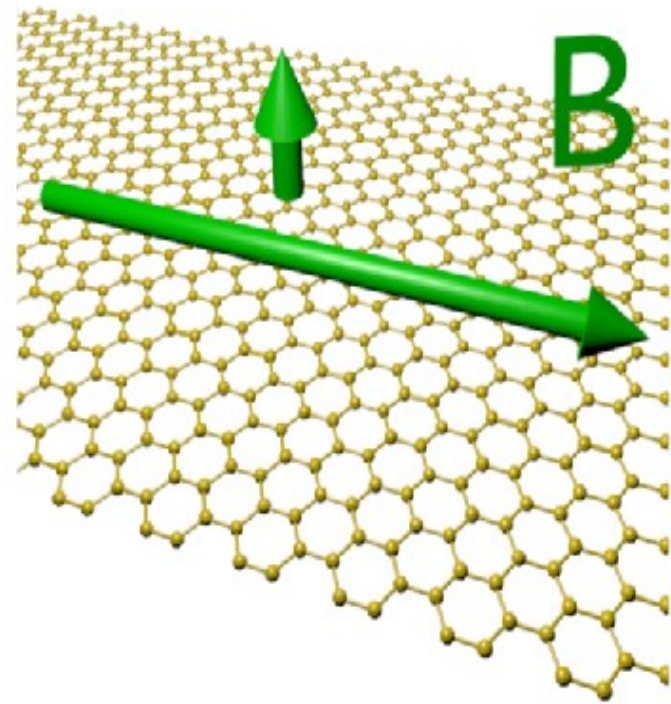
Abanin et al, PRL (2006)

Experiment

A. F. Young et al, Nature (2014)

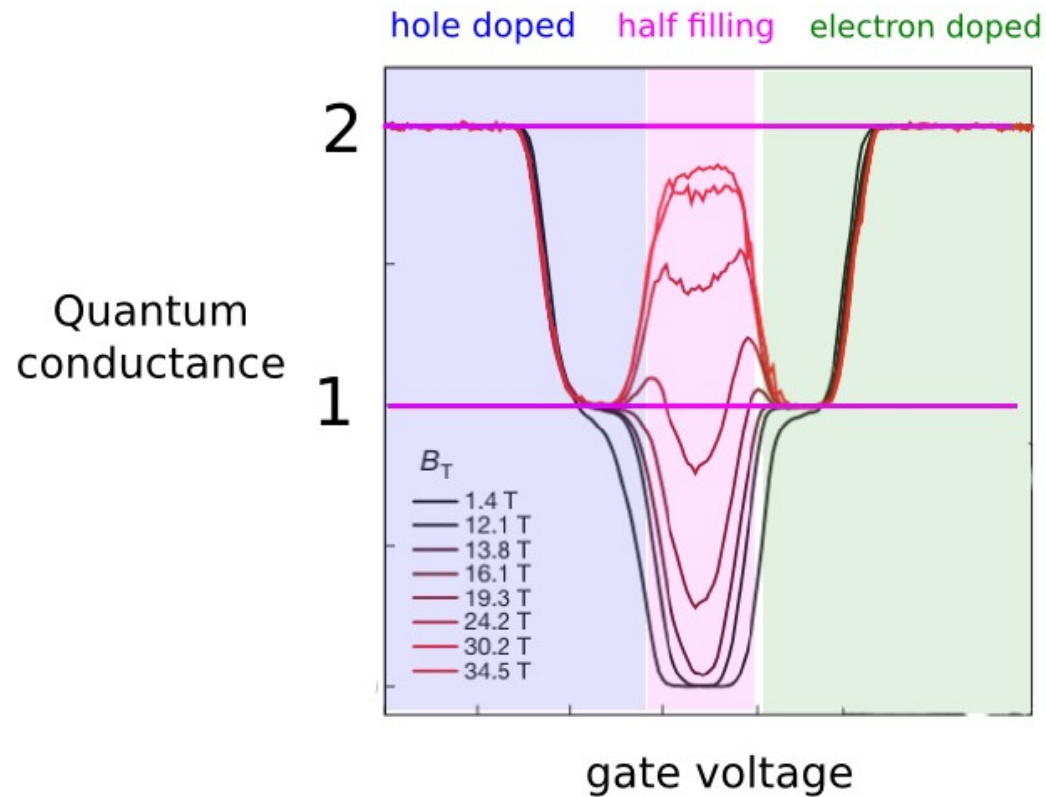
The experimental set-up

- High quality graphene
- Off-plane magnetic field
- In-plane magnetic field
- Variable doping via gate

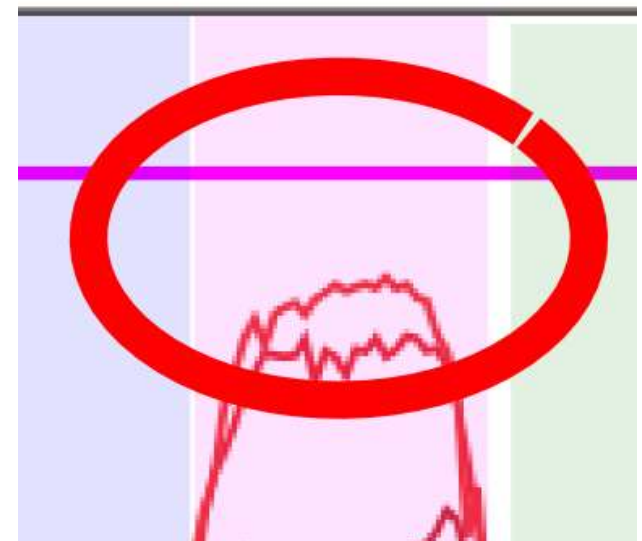


A. F. Young et al, Nature (2014)

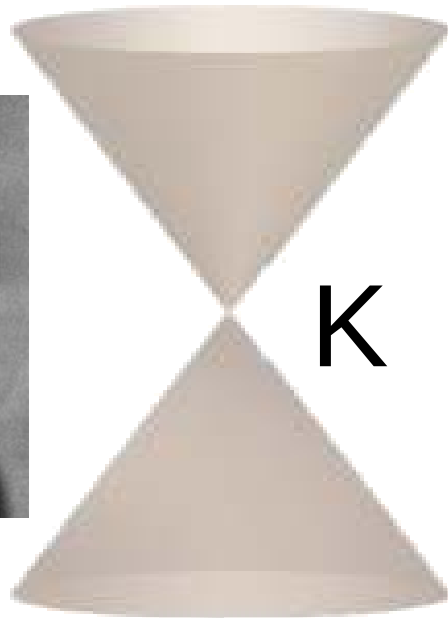
Not perfect conductance???



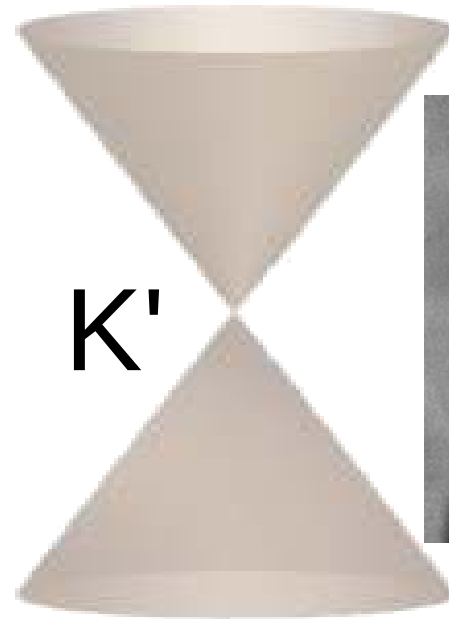
half filling



Low energy Hamiltonian



K



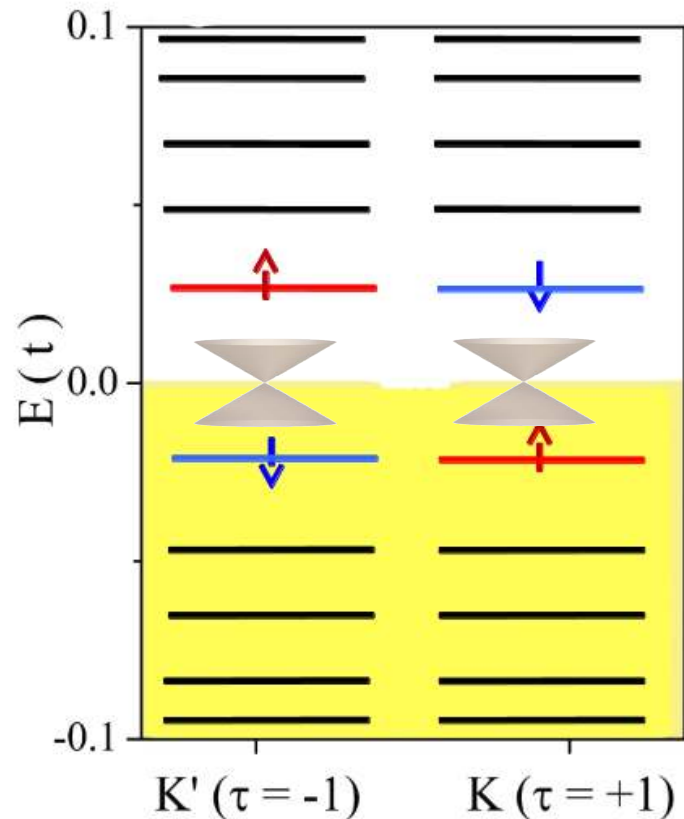
K'



Two Dirac equations

The Landau level picture

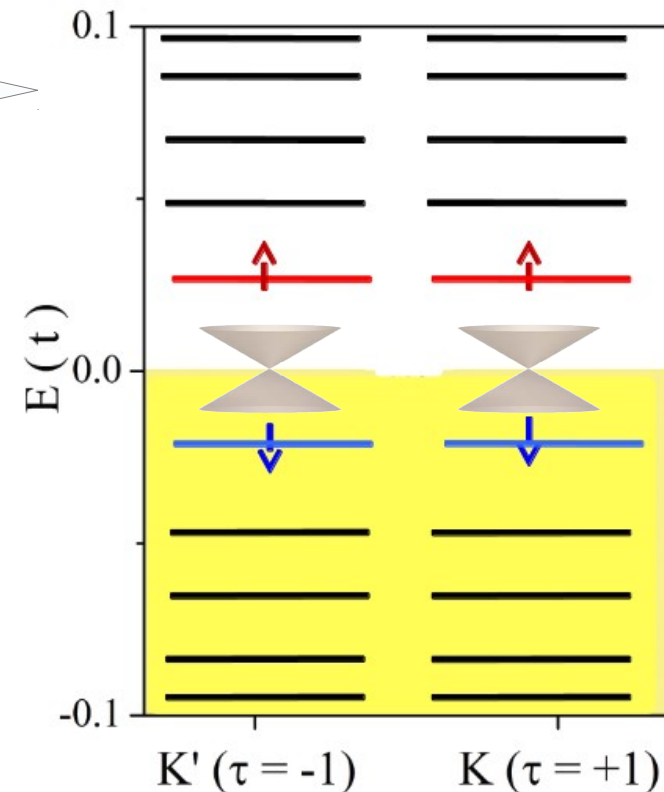
Antiferromagnetic



Spin Chern number = 0

Ferromagnetic

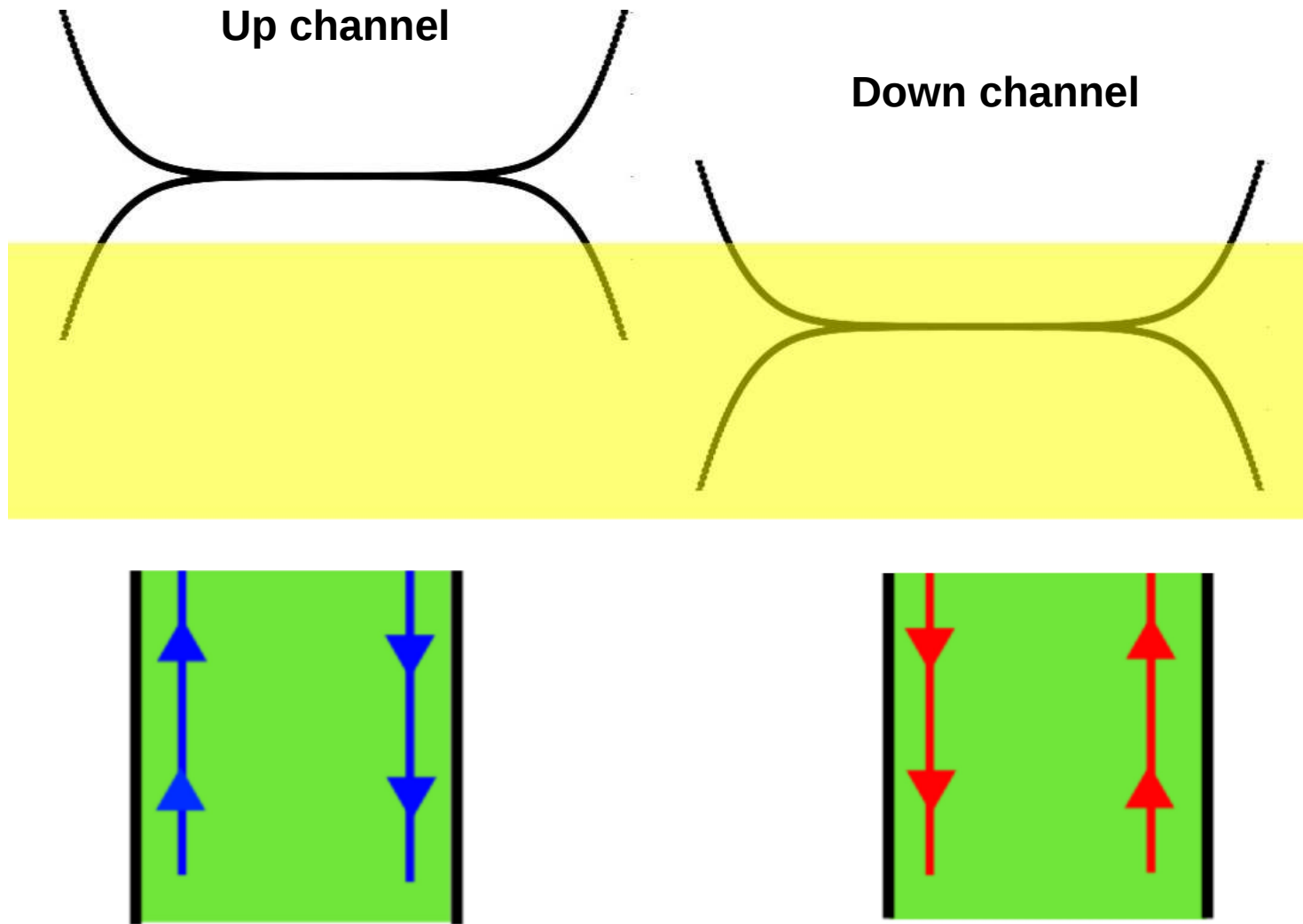
Increasing
Zeeman



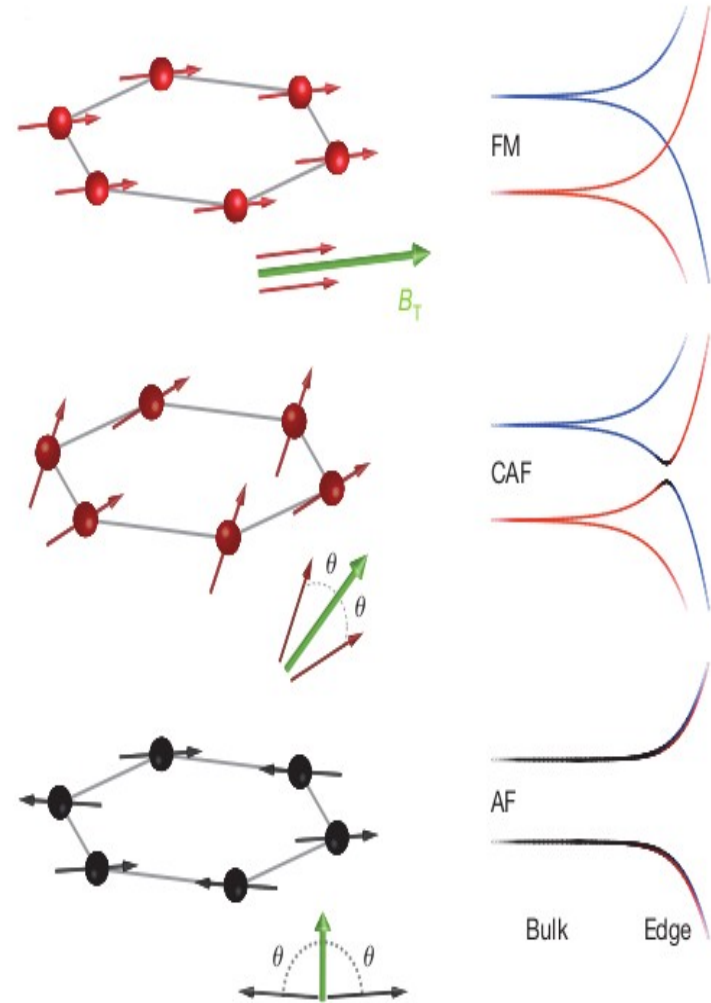
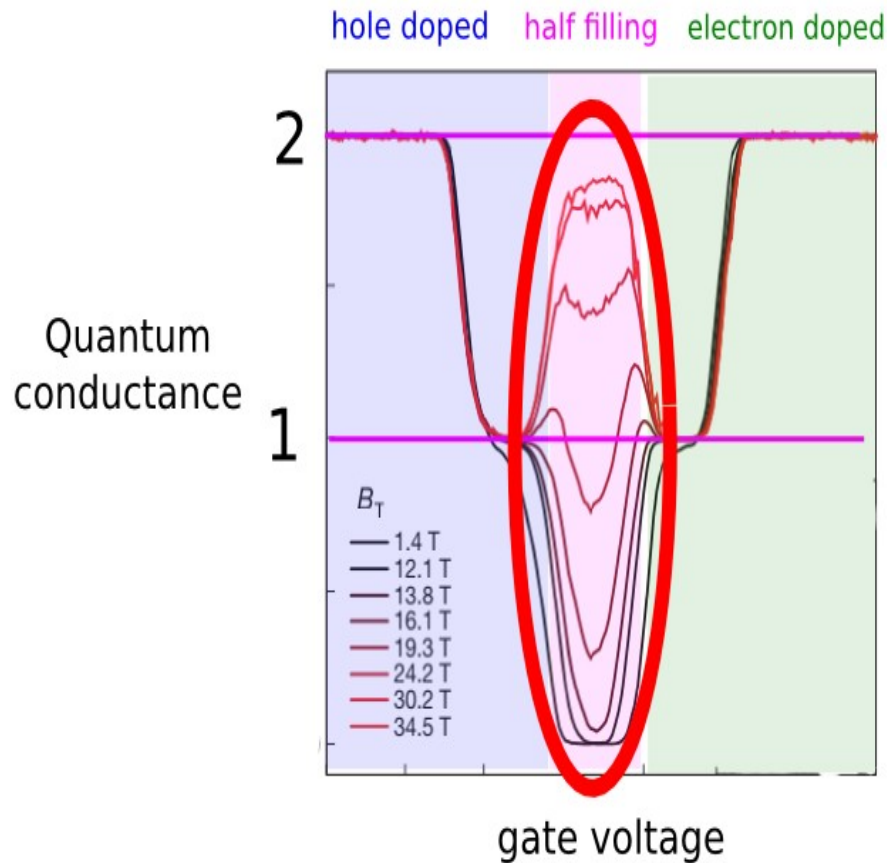
Spin Chern number = 2

Abanin et al, PRL (2006)

QSH in QH graphene



Canted magnetism



A. F. Young et al, Nature (2014)

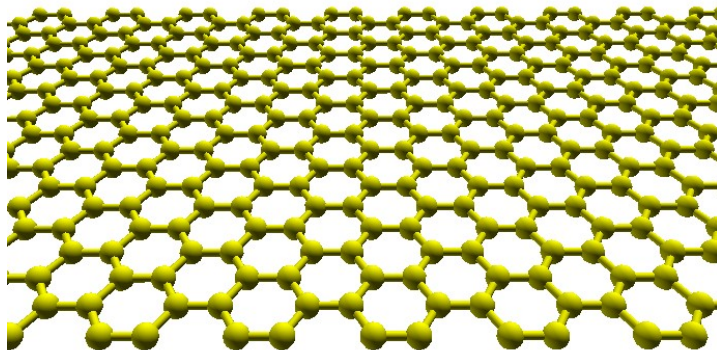
The model

$$\mathcal{H} = \mathcal{H}_0(B_z) + g\mu_B \vec{B} \cdot \vec{S} + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

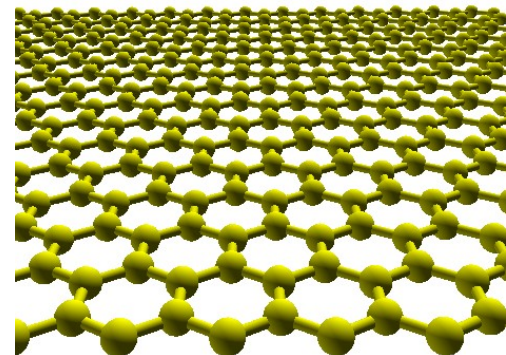
spin Zeeman field

nearest neighbor hopping
(with Peierls phase)

local e-e interaction
(Hubbard like)



ribbon geometry



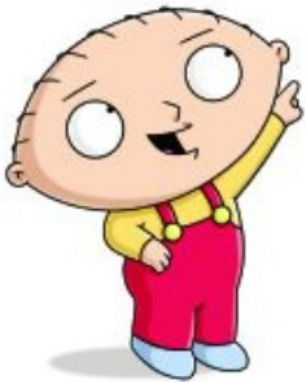
Solving the model

$$\mathcal{H} = \mathcal{H}_0(B_z) + g\mu_B \vec{B} \cdot \vec{S} + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

Mean field ansatz

$$H^{MF} = H_0 + g\mu_B \vec{B} \cdot \vec{S} + H_H + H_F + E_{DC}$$

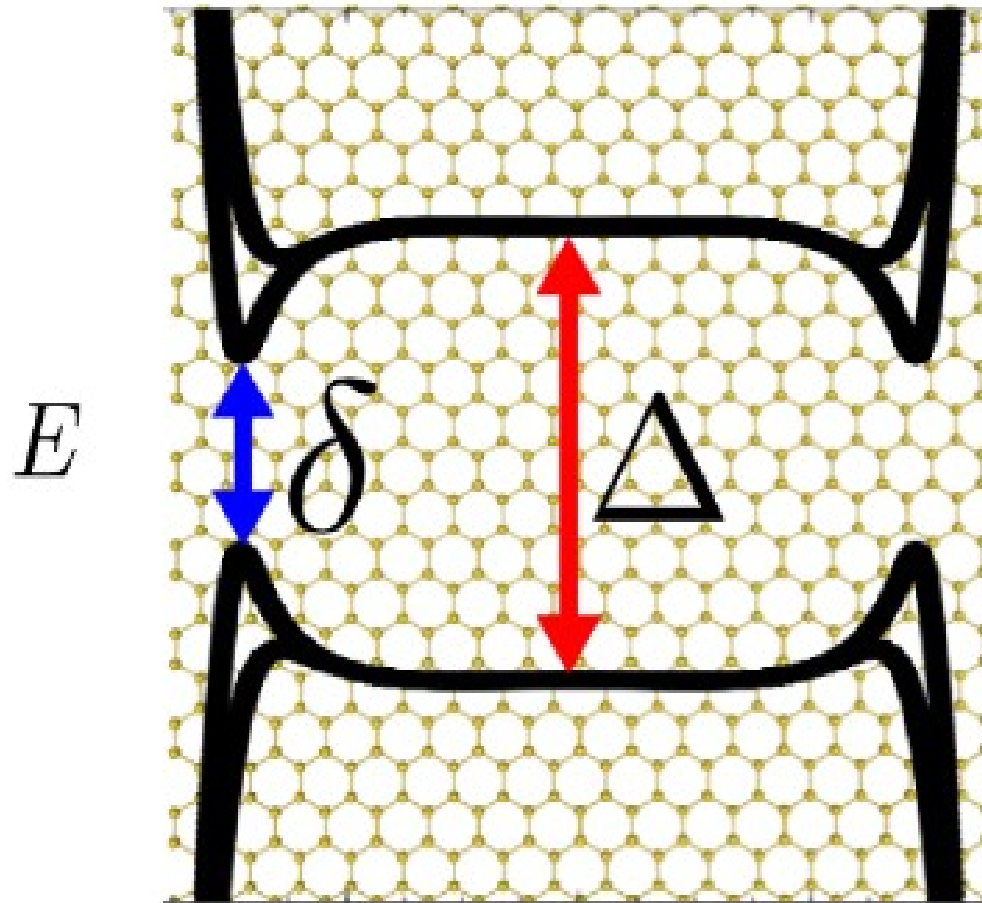
Collinear (Hartree) Exchange (Fock) Double counting



Variational wave function

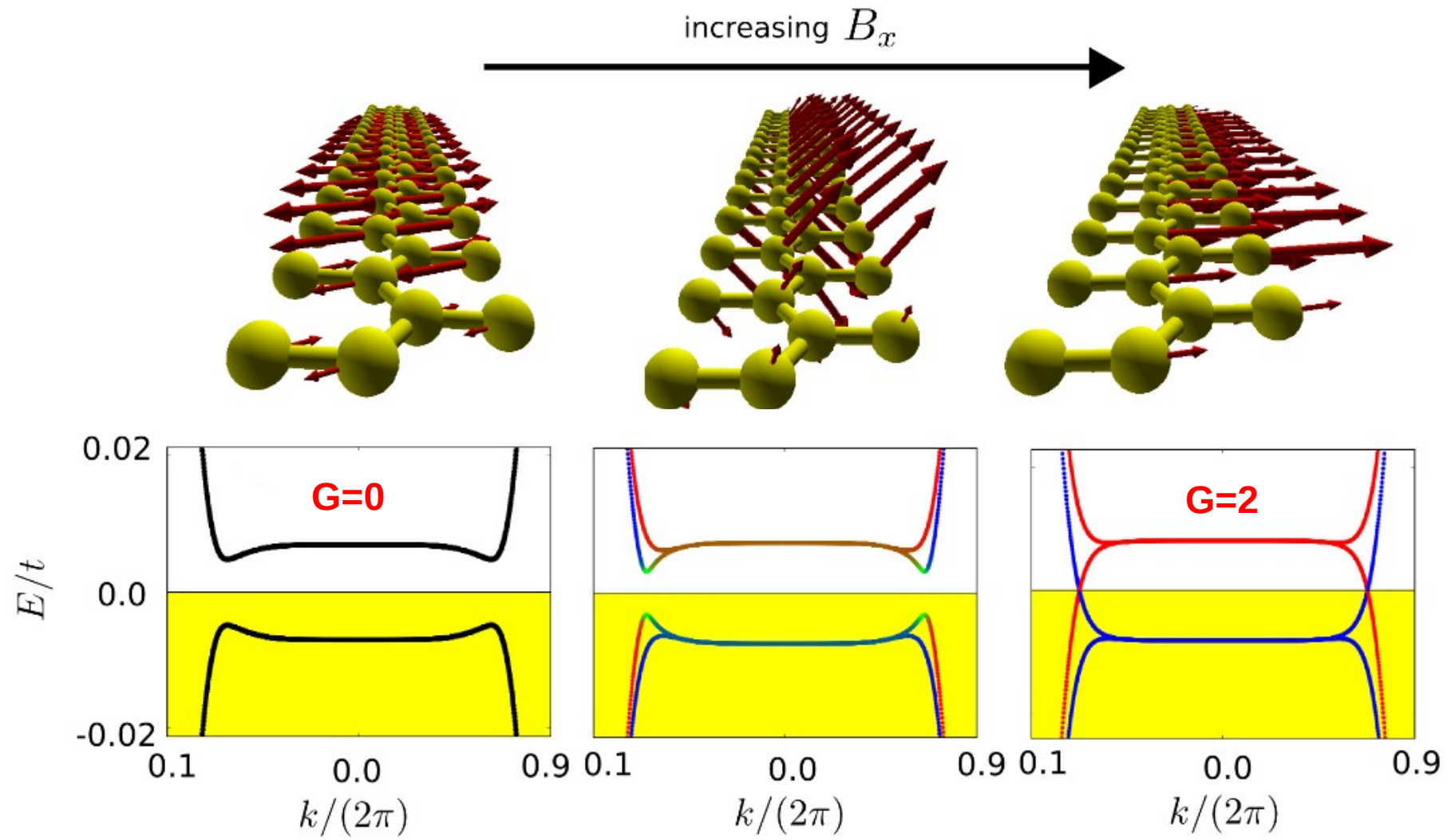
$$|\Omega\rangle = \prod_i \left(u_i c_{i\uparrow}^\dagger + v_i c_{i\downarrow}^\dagger \right) |0\rangle$$

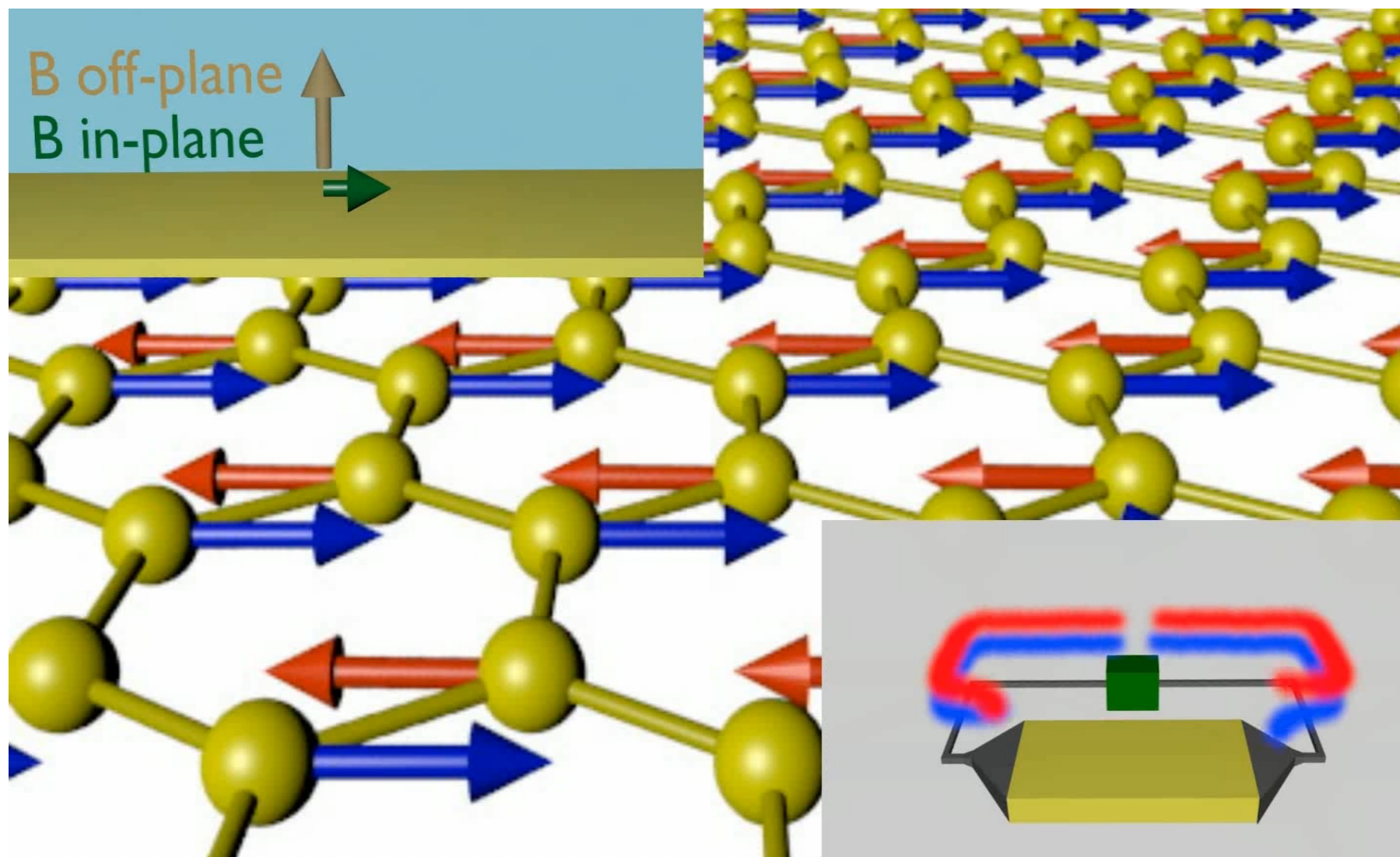
Bands of the system



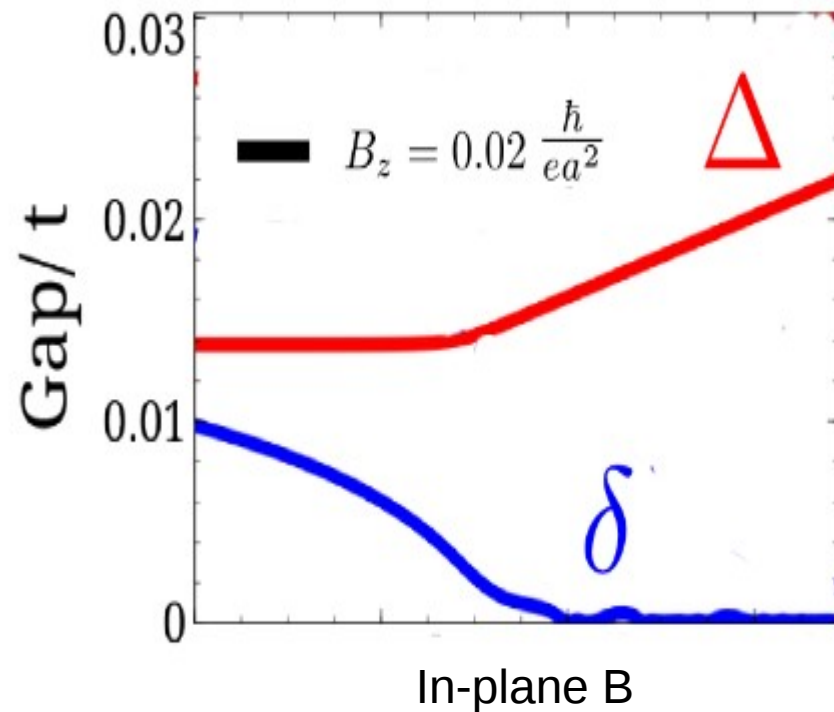
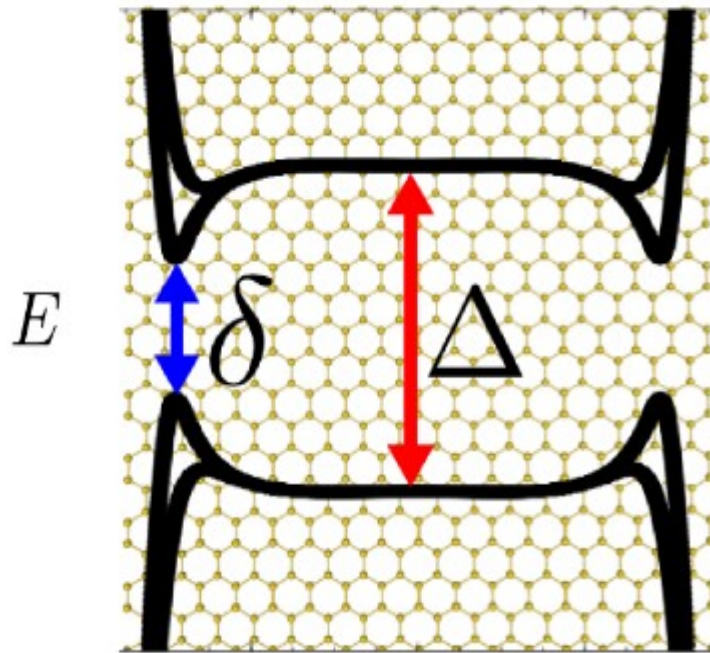
K-vector (and position across the ribbon)

Transition at half filling with B_x





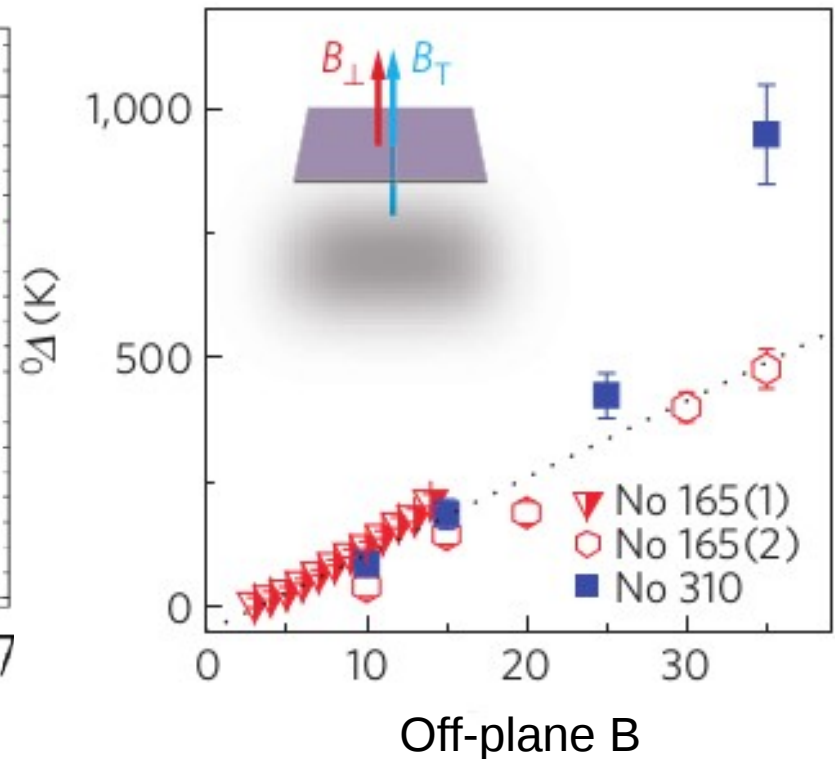
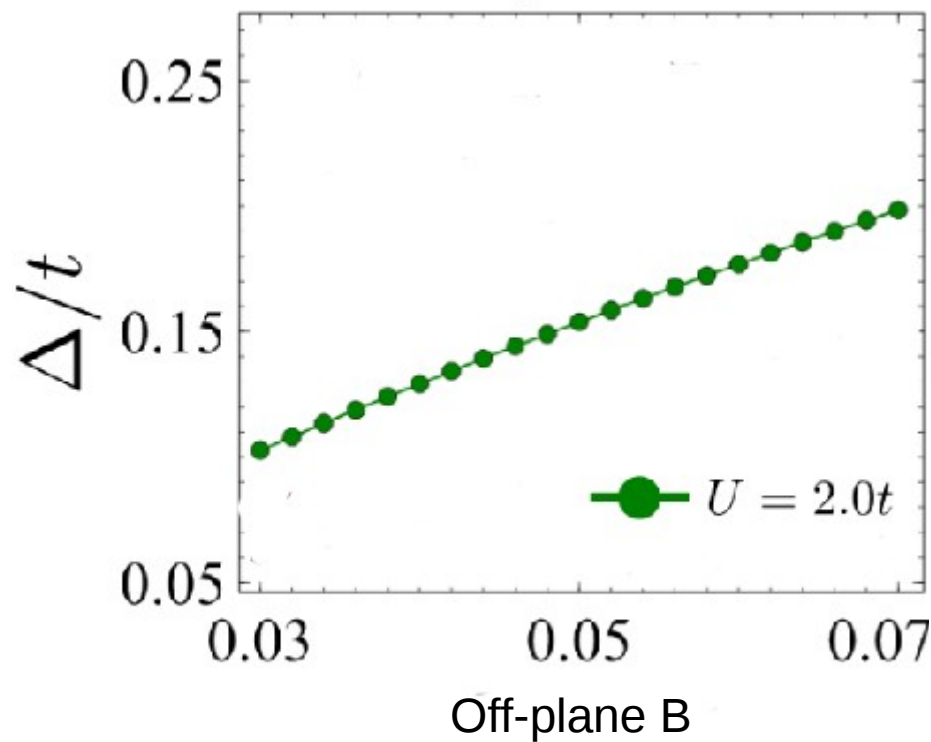
Transition at half filling with B_x



Edge gap closing without bulk gap closing

Gap increases at half filling

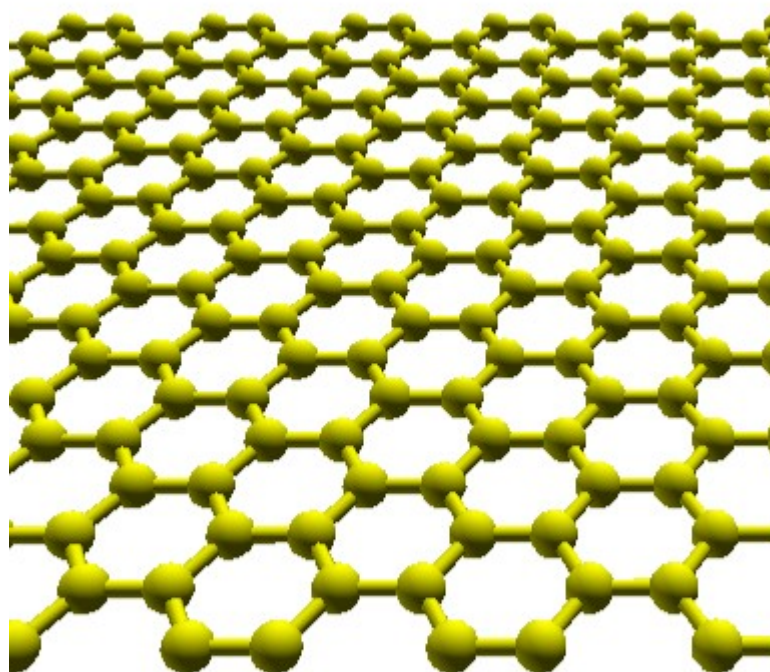
Gap in the bulk



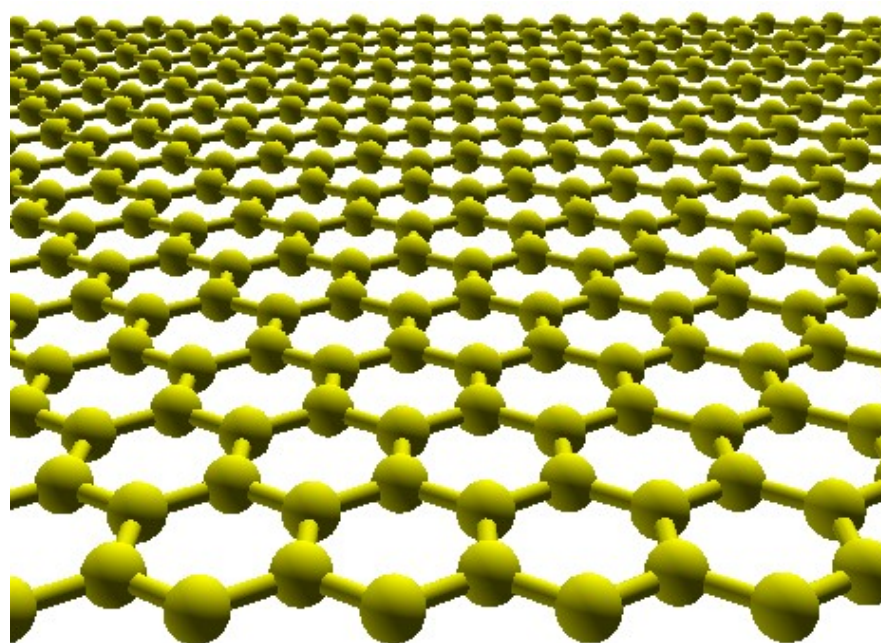
Gap increases linearly with off-plane field

Different edges

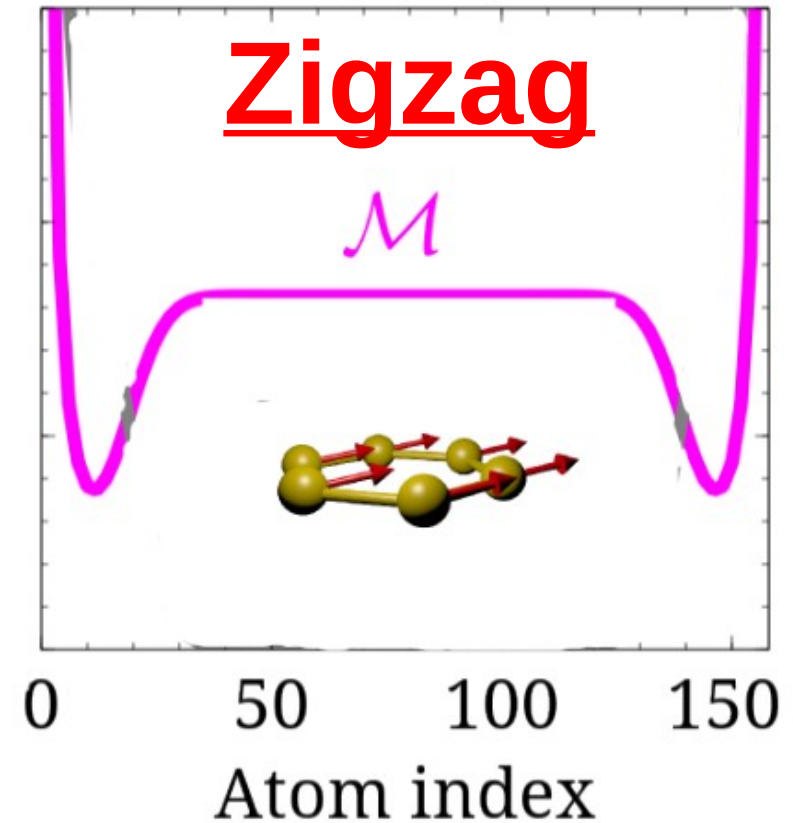
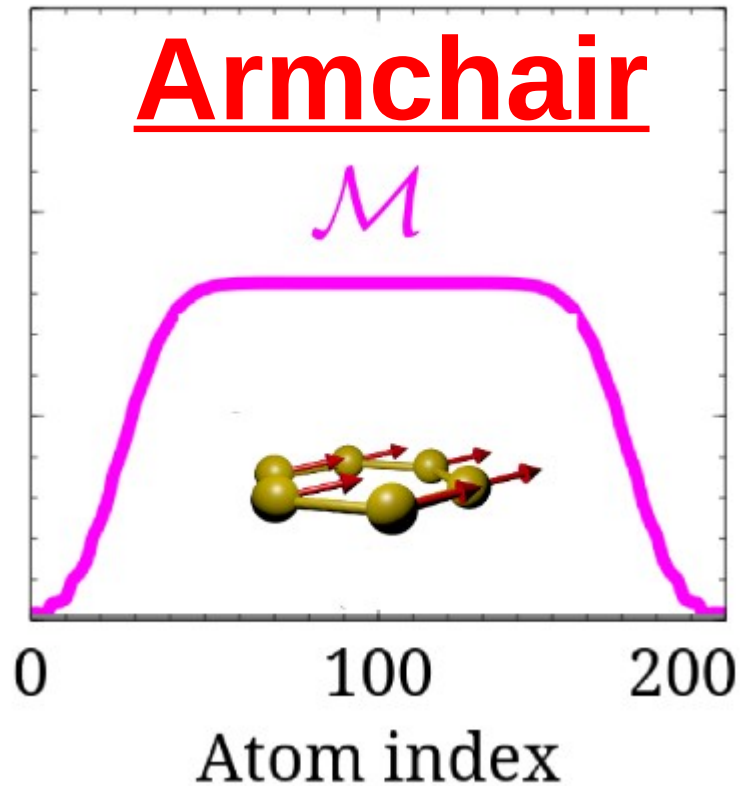
Armchair



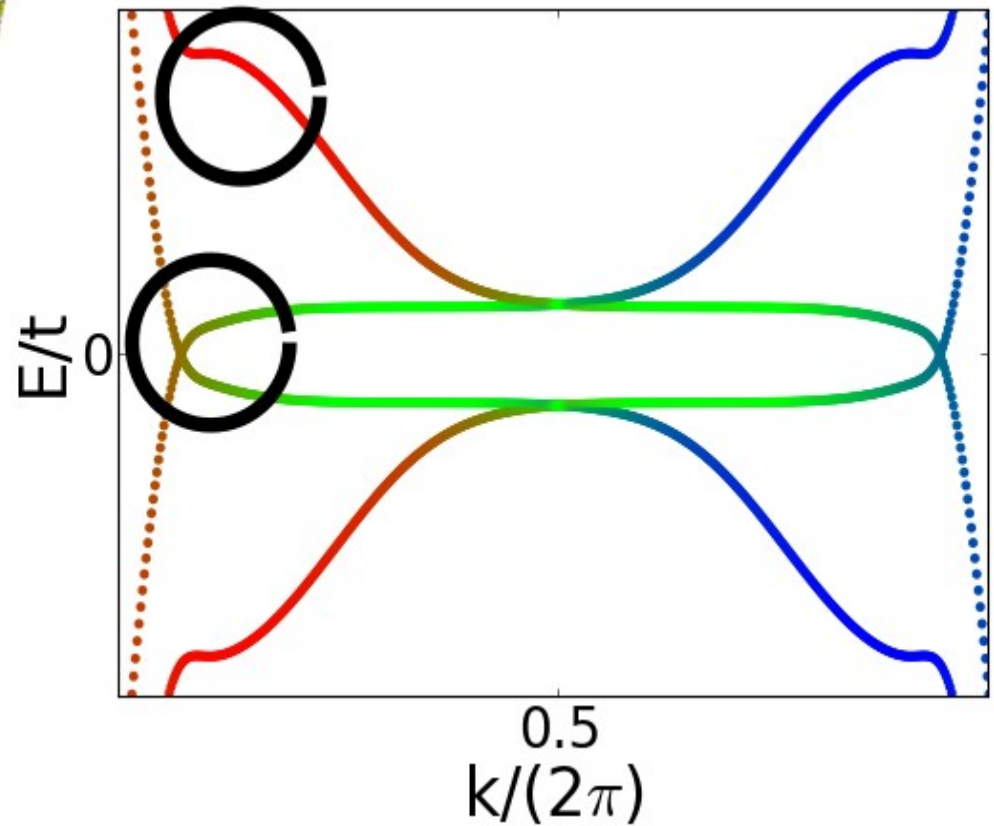
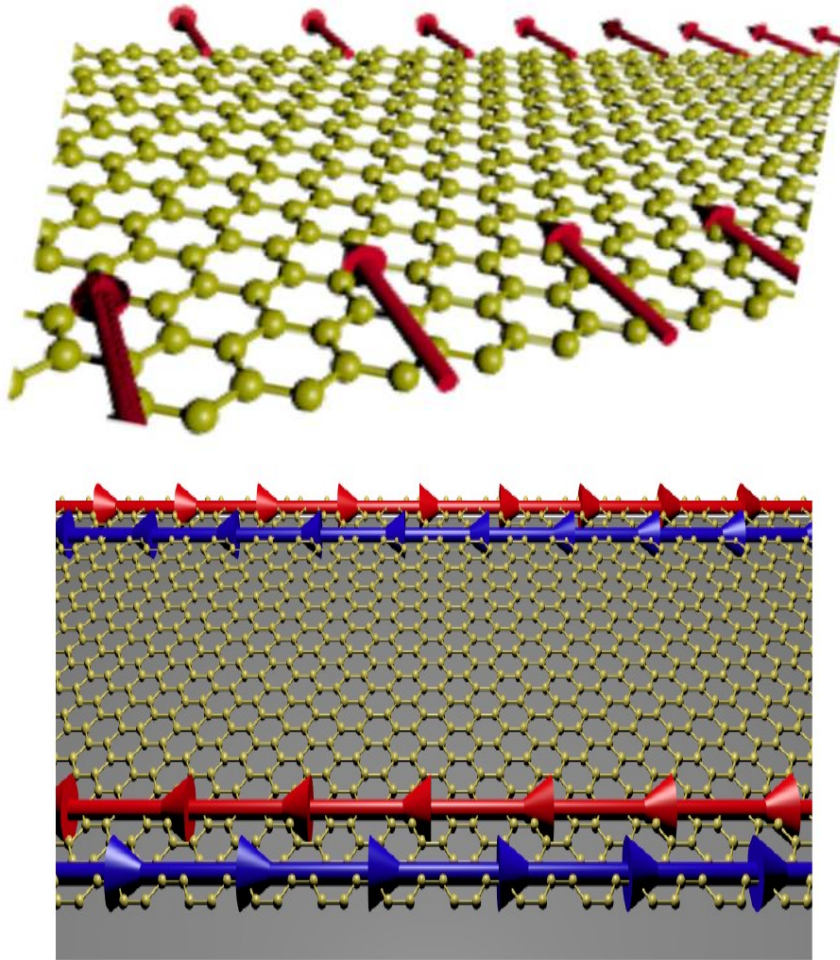
Zigzag



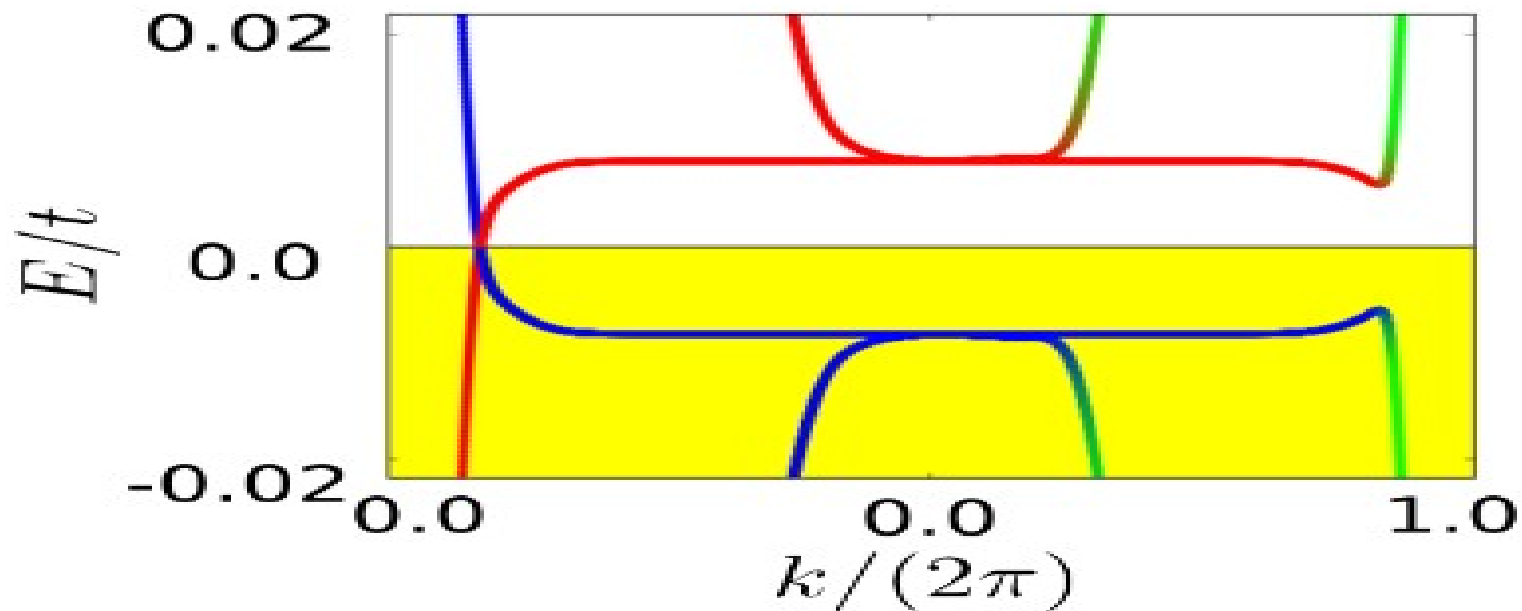
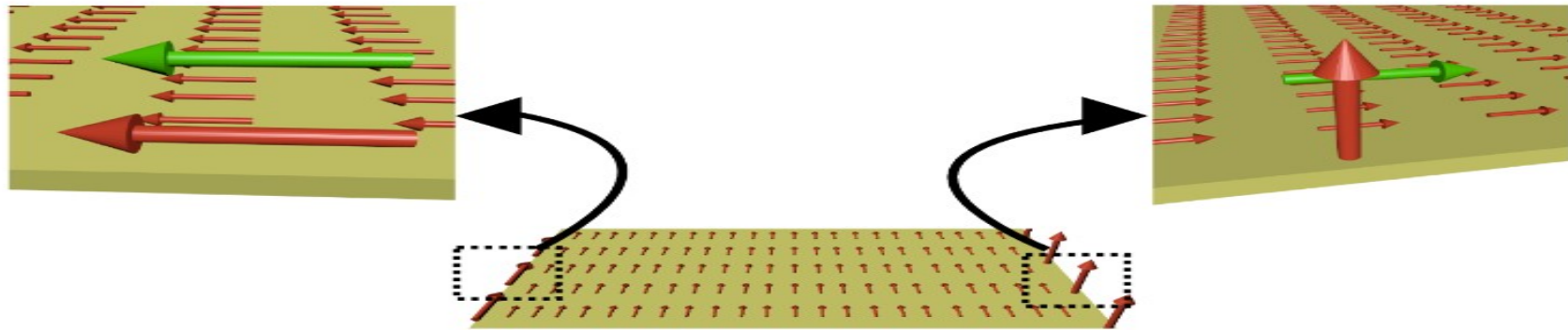
Edge dependent magnetism



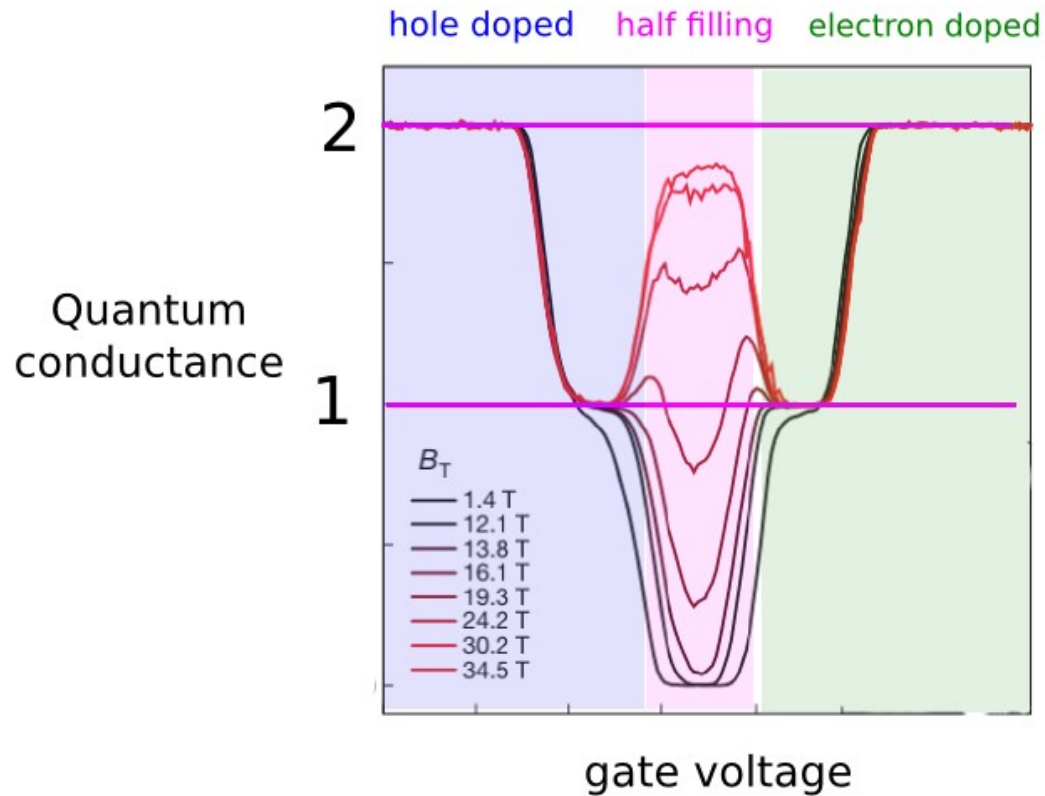
Two different edge states



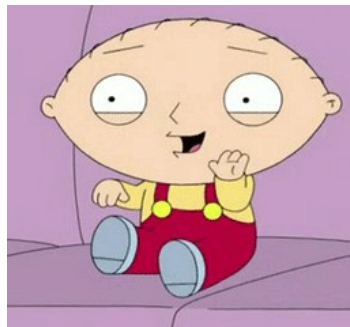
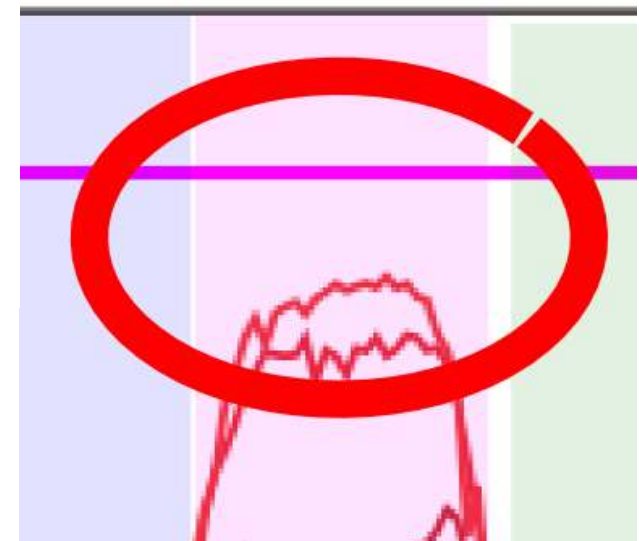
Electron-magnon coupling

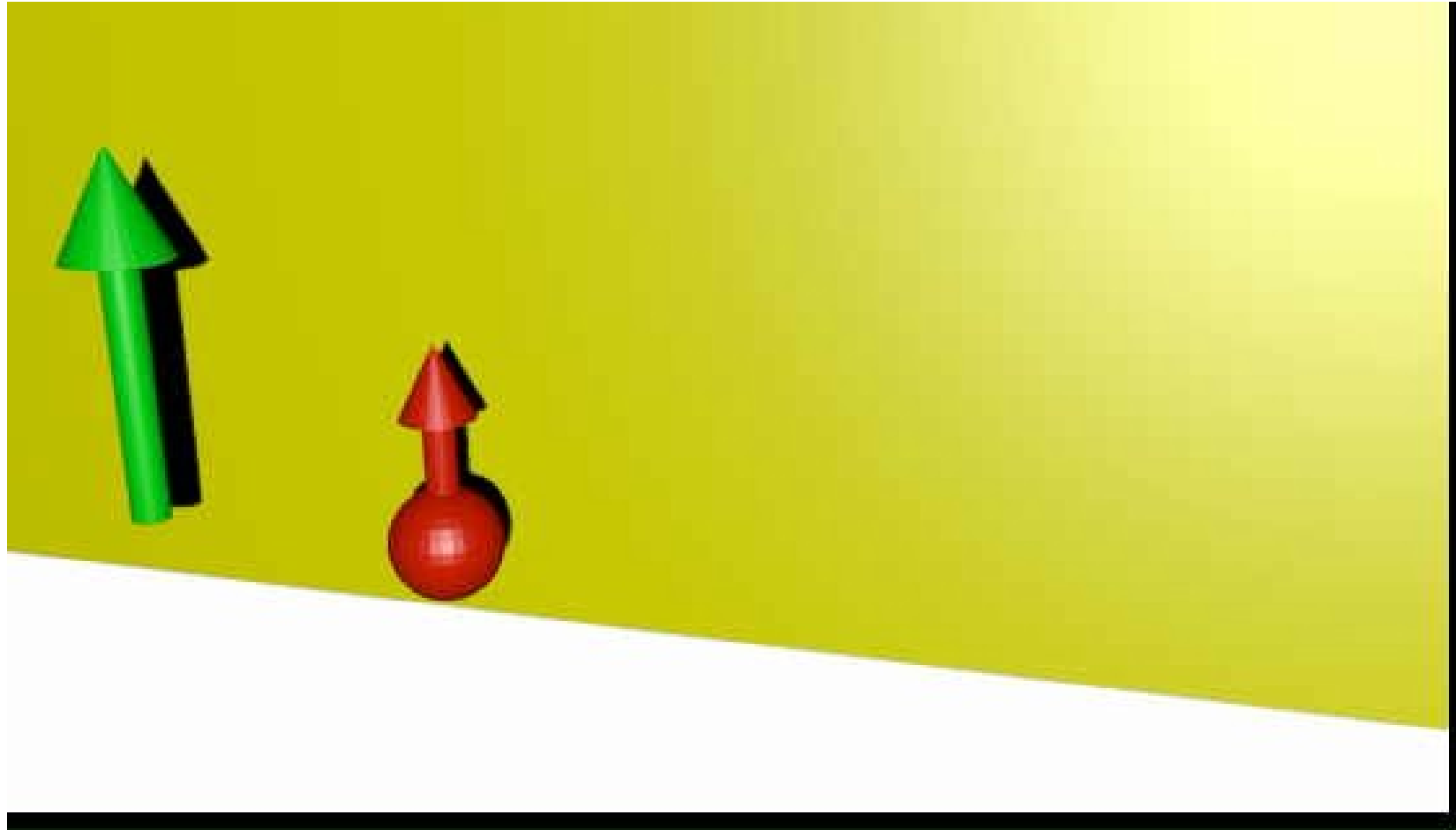


Back scattering channel

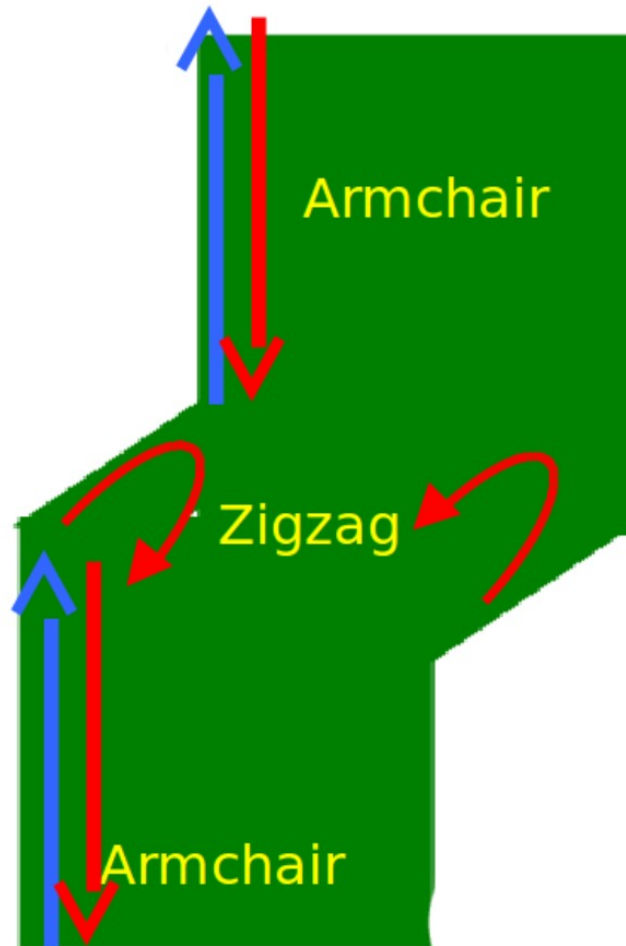
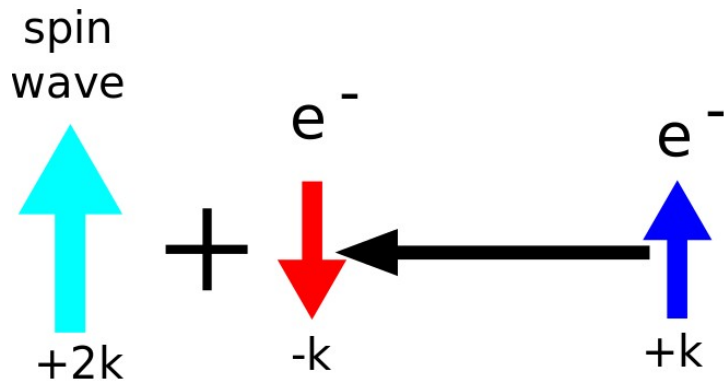


half filling





Mixed edges



Backscattering channel in zigzag part

Take home

- Interaction of different edge states opens up a channel for backscattering

But before finishing...

Funded by



Thank you!!!