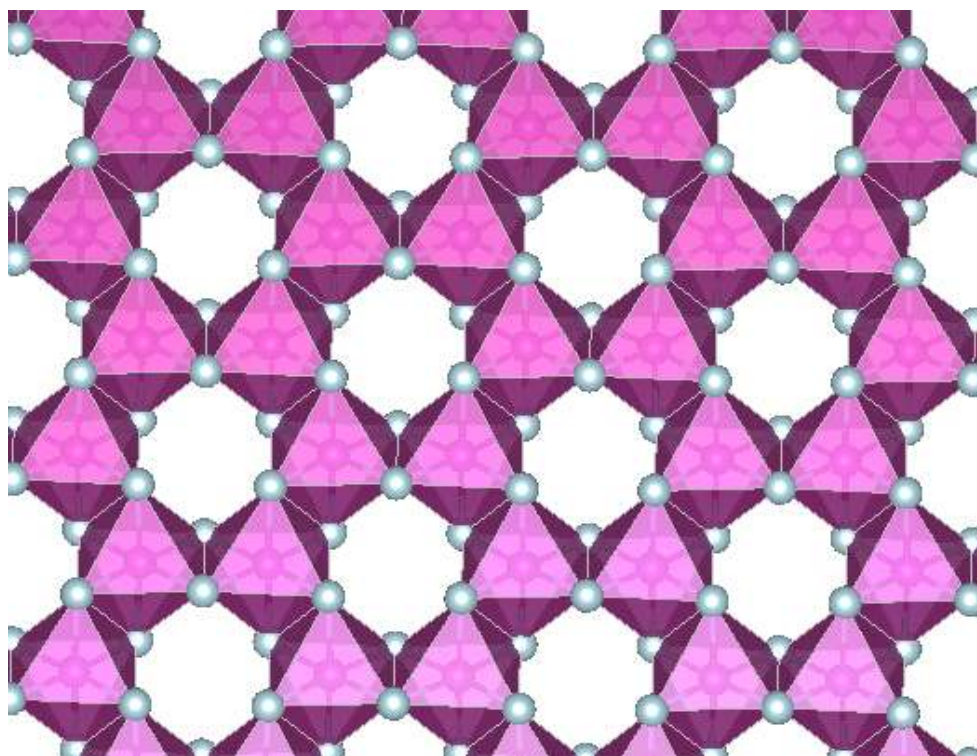


On the origin of magnetic anisotropy in two dimensional CrI_3

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International Iberian Nanotechnology Laboratory, Portugal



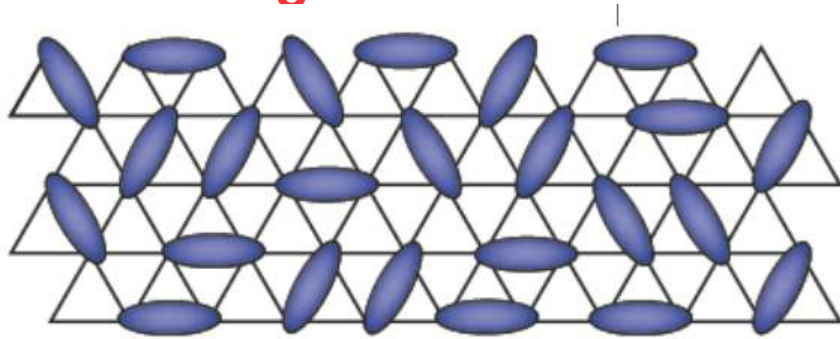
arXiv:1704.03849

Thursday 15th June, CDT Conference, Cambridge 2017

2d magnets join the game

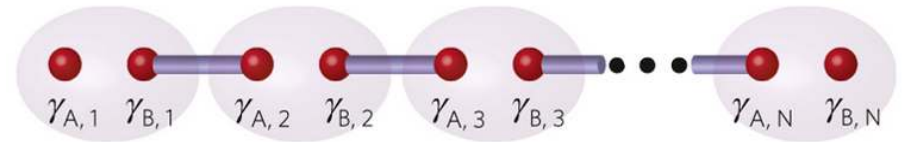
Quantum spin liquid

2d magnet + frustration



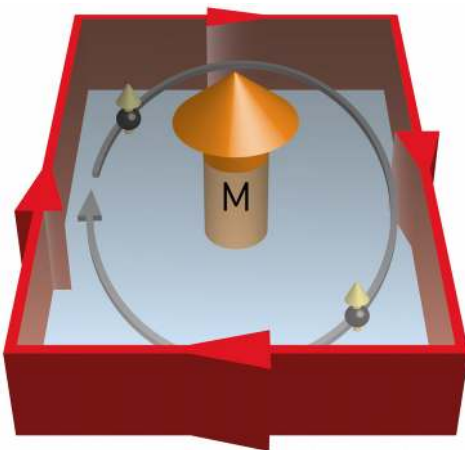
2d topological superconductor

2d magnet + Rashba + superconductor



Quantum Anomalous Hall insulator

2d magnet + 3d topological insulator



Magnetism in two dimensions

A perfect isotropic system does not order at finite temperature in 2d (Mermin-Wagner theorem)

Magnetic anisotropy is a must for 2d magnetism

Two iso-structural compounds

Chromium iodine (CrI_3)

Off-plane ferromagnet

Anisotropy field of 3 T

Chromium chlorine (CrCl_3)

In-plane ferromagnet

Anisotropy field of 0.1 T

Michael A. McGuire Crystals 7.5 (2017): 121.

Too similar, yet too different

What is the ultimate origin of the magnetic anisotropy in CrI_3 ?

Ferromagnetism in 2d CrI_3

nature

Layer-dependent ferromagnetism in a van der Waals crystal down to the monolayer limit

B. Huang et al, Nature 2017

Two simple open questions

Off-plane easy axis, are they Ising spins?

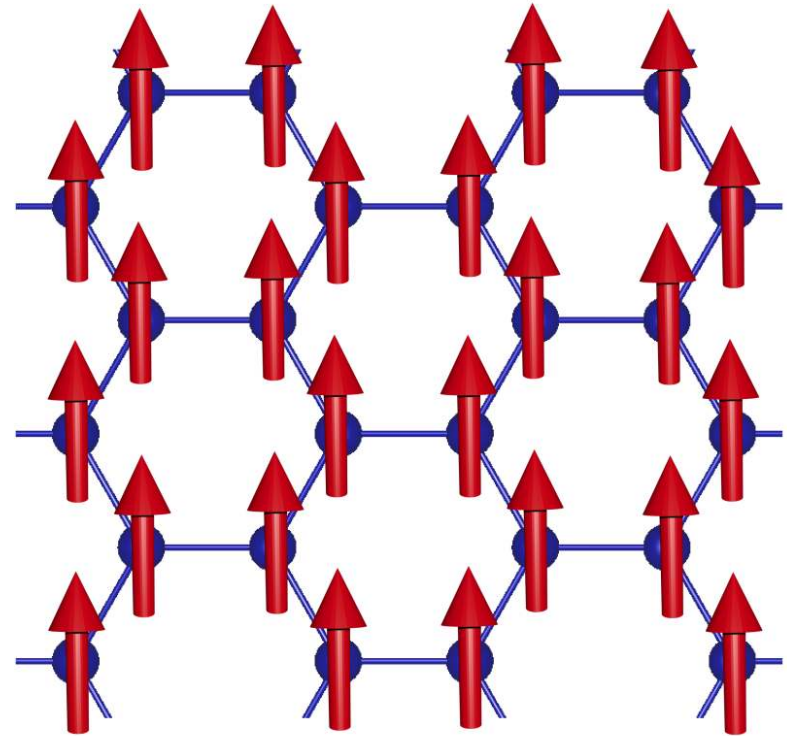
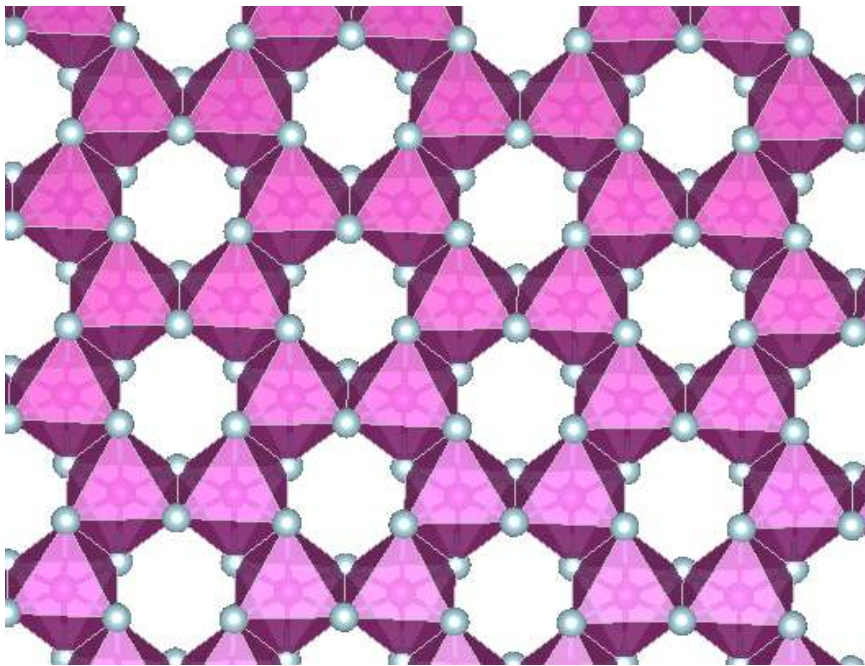
Critical temperature of 45 K, but what controls this?

Is CrI_3 a two dimensional Ising model?

If not, what is the right model for CrI_3 ?

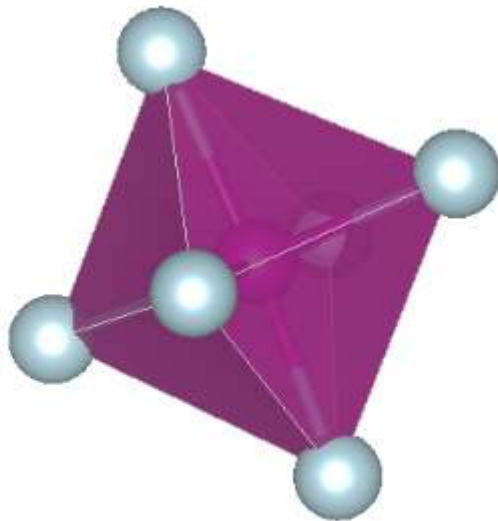


CrI_3 , ferromagnetic honeycomb lattice

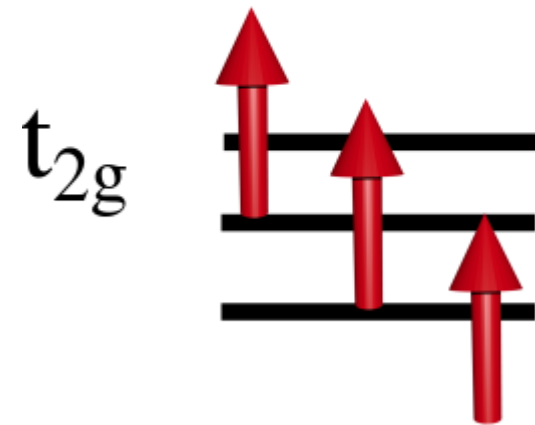


Each chromium atom contributes with $S=3/2$

Electronic filling and anisotropy



Octahedral iodine environment

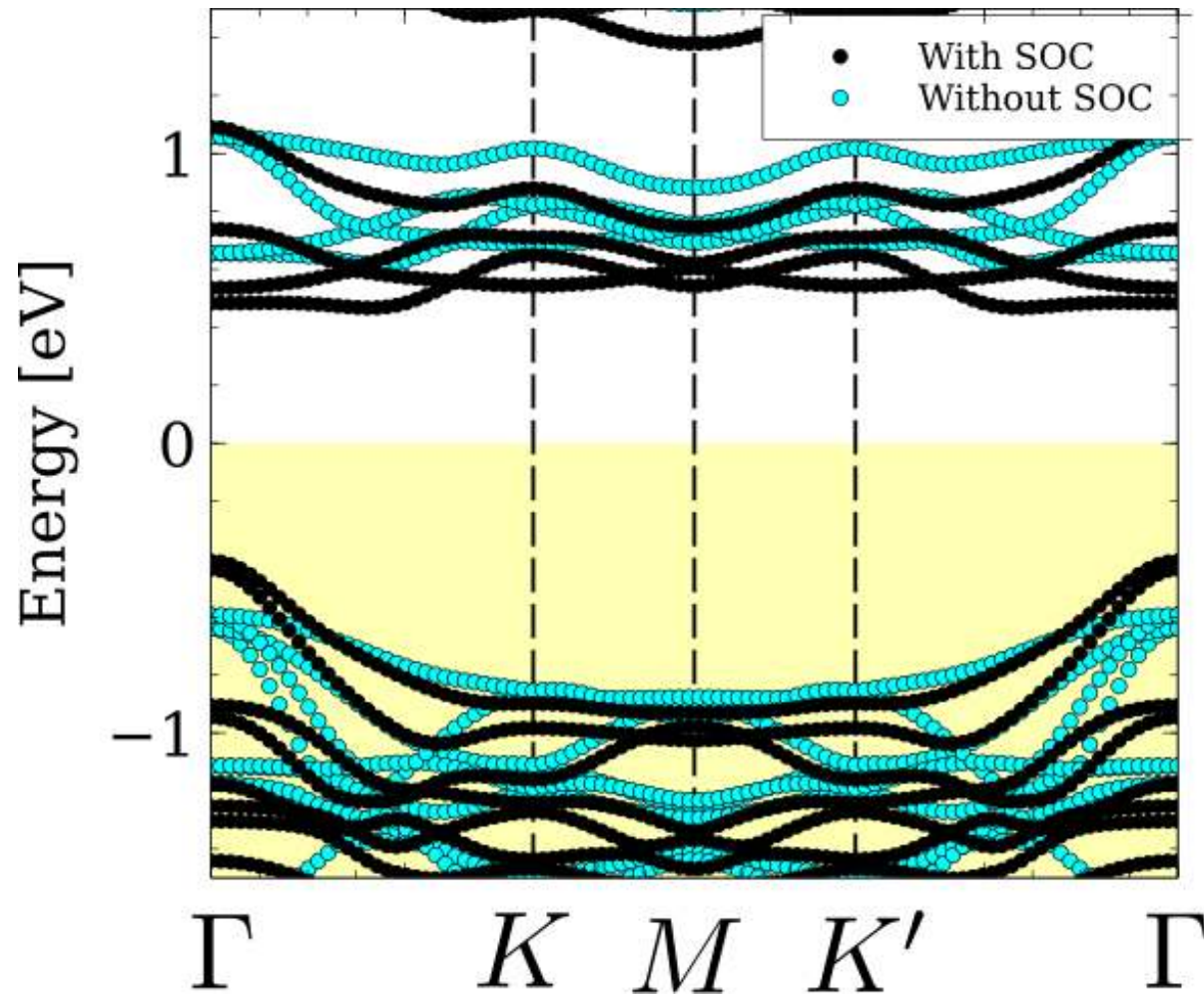


Quenched angular momenta

With quenched angular momenta, single ion anisotropy is extremely small

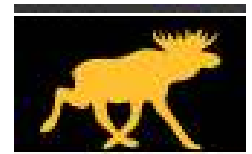
How is there a sizable magnetic anisotropy in CrI_3 ?

Effect of spin orbit coupling in CrI_3

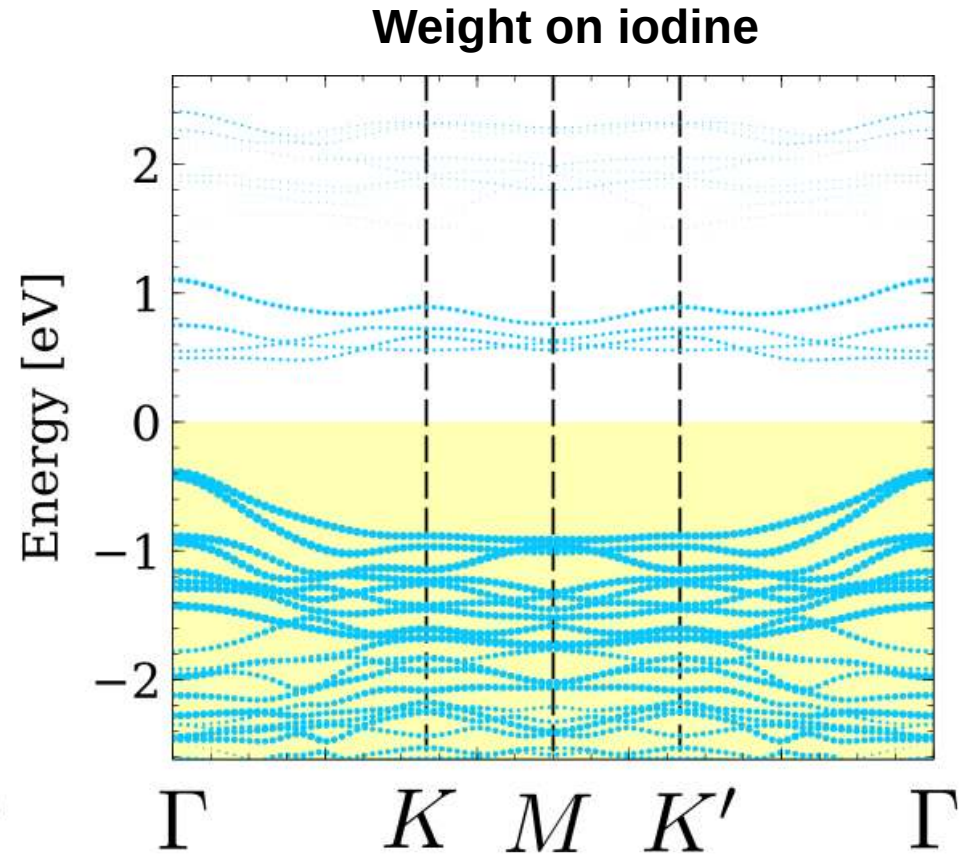
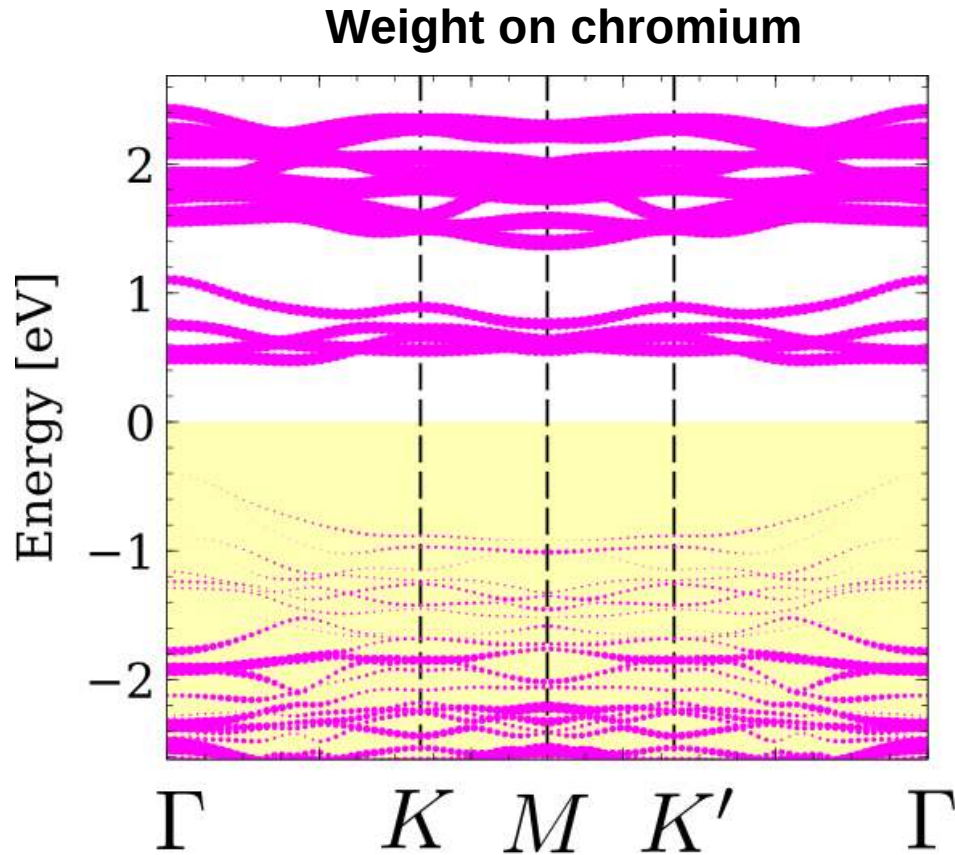


Elk LAPW code

Effect of SOC in the band structure is extremely large



Weights over the atoms in CrI_3



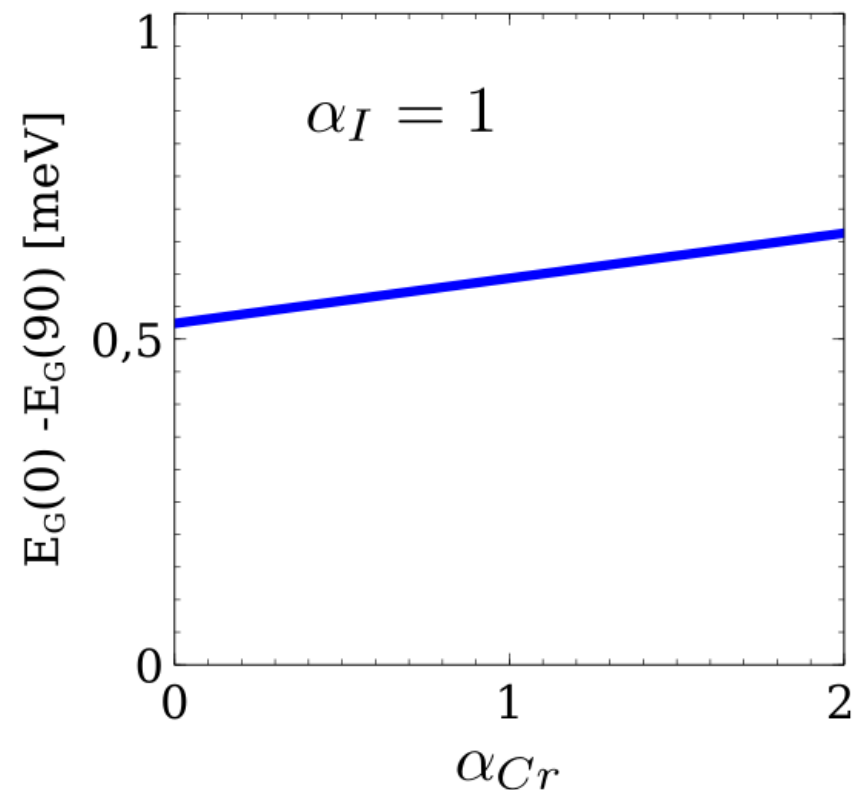
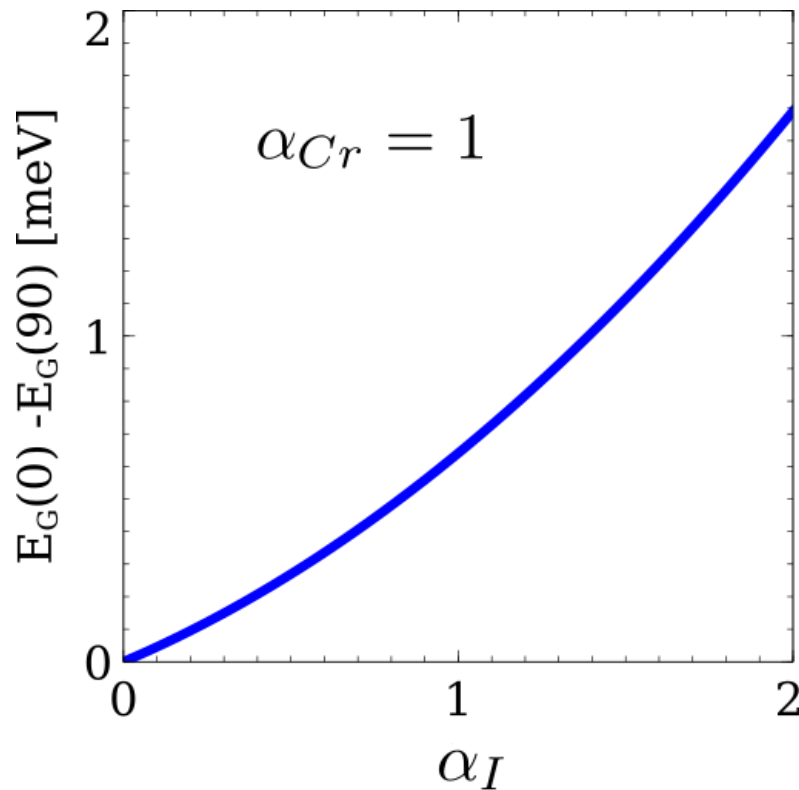
There is a sizable contribution of iodine above and below the gap

Is iodine a key to determine the magnetic properties?

Magnetic anisotropy in CrI_3

We can separately see the effect of the two atoms to magnetic anisotropy

$$\mathcal{H}_{\text{DFT}}(\alpha_I, \alpha_{Cr}) = \mathcal{H}_0 + \alpha_I \mathcal{H}_I^{\text{SOC}} + \alpha_{Cr} \mathcal{H}_{Cr}^{\text{SOC}}$$



The major contribution to magnetic anisotropy comes from iodine

Disentangling the spin Hamiltonian

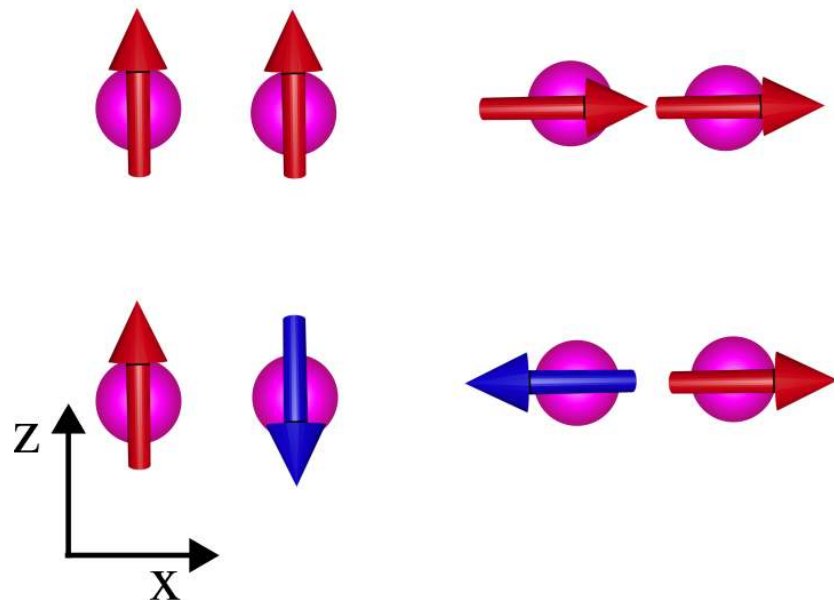
Minimal Spin Hamiltonian

$$\mathcal{H} = - \left(\sum_i D(S_i^z)^2 + \frac{J}{2} \sum_{i,i'} \vec{S}_i \cdot \vec{S}_{i'} + \frac{\lambda}{2} \sum_{i,i'} S_i^z S_{i'}^z \right)$$

Single ion
anisotropy

Isotropic
Heisenberg

Anisotropic
interaction



$$\mathcal{E}_{FM,z} = -2S^2 D - 3S^2 (J + \lambda)$$

$$\mathcal{E}_{AF,z} = -2S^2 D + 3S^2 (J + \lambda)$$

$$\mathcal{E}_{FM,x} = -3S^2 J$$

$$\mathcal{E}_{AF,x} = +3S^2 J$$

The true spin model for CrI_3

Hamiltonian

$$\mathcal{H} = - \left(\sum_i D (S_i^z)^2 + \frac{J}{2} \sum_{i,i'} \vec{S}_i \cdot \vec{S}_{i'} + \frac{\lambda}{2} \sum_{i,i'} S_i^z S_{i'}^z \right)$$

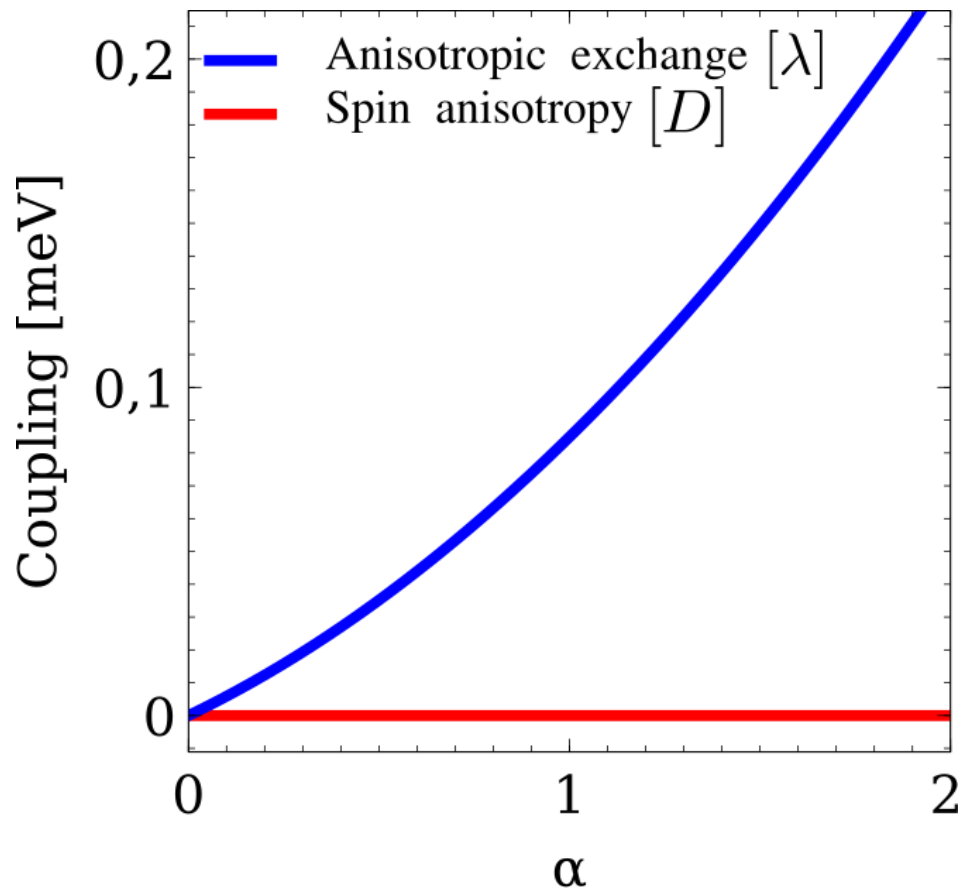
Energetics

$$\mathcal{E}_{FM,z} = -2S^2 D - 3S^2 (J + \lambda)$$

$$\mathcal{E}_{AF,z} = -2S^2 D + 3S^2 (J + \lambda)$$

$$\mathcal{E}_{FM,x} = -3S^2 J$$

$$\mathcal{E}_{AF,x} = +3S^2 J$$



Single ion anisotropy is negligible $\rightarrow$ **not a 2d Ising system**

Sizable anisotropic exchange $\rightarrow$ **CrI_3 is a 2d realization of the XXZ model**

Magnetic fluctuations: spin waves

Holstein-Primakoff

Spin Hamiltonian

$$\mathcal{H} = - \left(\sum_i D(S_i^z)^2 + \frac{J}{2} \sum_{i,i'} \vec{S}_i \cdot \vec{S}_{i'} + \frac{\lambda}{2} \sum_{i,i'} S_i^z S_{i'}^z \right)$$

$$S_i^+ = \sqrt{2S} \sqrt{1 - \frac{b_i^\dagger b_i}{2S}} b_i$$

$$S_i^- = \sqrt{2S} b_i^\dagger \sqrt{1 - \frac{b_i^\dagger b_i}{2S}}$$

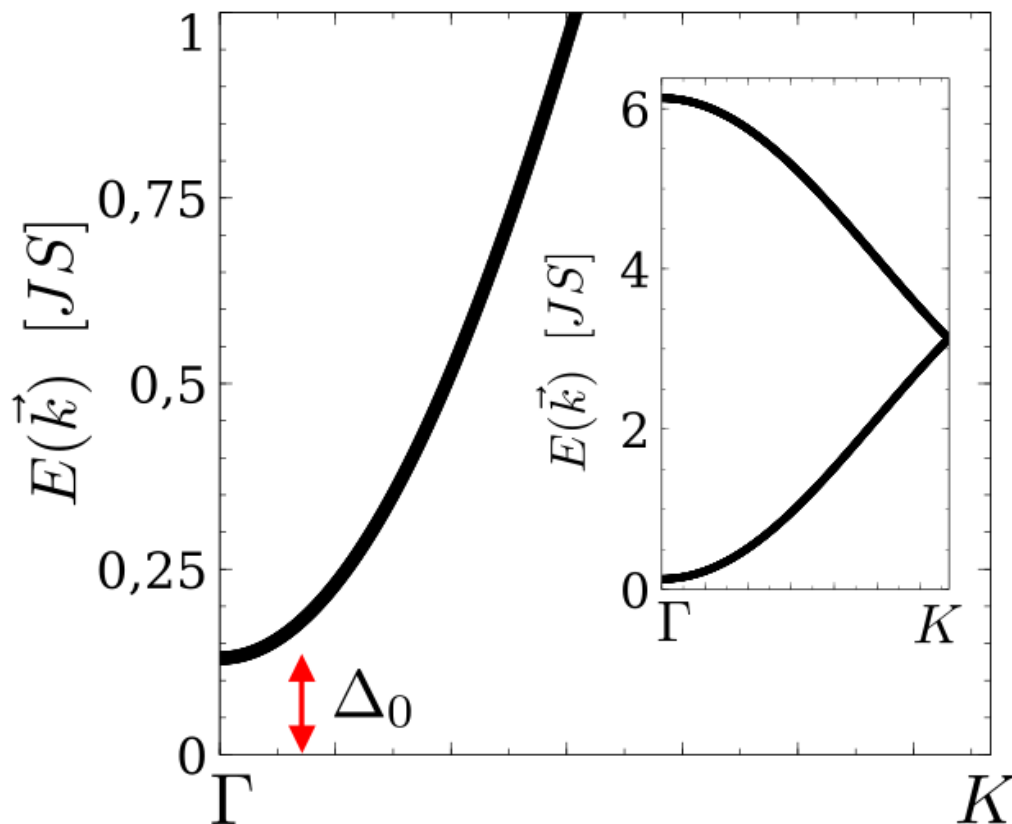
$$S_i^z = S - b_i^\dagger b_i$$

Spin wave Hamiltonian (bosonic Hamiltonian)

$$\mathcal{H}_{\text{spin waves}} = \sum_i (2DS + 3S(J + \lambda)) b_i^\dagger b_i - JS \sum_{\langle ij \rangle} b_i^\dagger b_j$$

Spin wave spectra

Spin wave spectra



$$\Delta_0 = 2DS + 3S\lambda$$

Spin Wave Bloch Hamiltonian

$$\mathcal{H}_{SW}(\vec{k}) = \begin{pmatrix} \epsilon_0 & -JSf(\vec{k}) \\ -JSf^*(\vec{k}) & \epsilon_0 \end{pmatrix}$$

**Correction to magnetization
without magnetic anisotropy**

$$\delta M = \frac{1}{4\pi} \int_0^{k_c} \frac{k dk}{\beta \rho k^2} \rightarrow \infty$$

Spin wave Hamiltonian at finite temperature

First order expansion

$$S_i^+ \approx \sqrt{2S} \left(1 - \frac{b_i^\dagger b_i}{4S} \right) b_i$$

$$S_i^- \approx \sqrt{2S} b_i^\dagger \left(1 - \frac{b_i^\dagger b_i}{4S} \right)$$

$$S_i^z = S - b_i^\dagger b_i$$

Mean field approximation

$$b_i^\dagger b_i b_j^\dagger b_j \approx \langle b_i^\dagger b_i \rangle b_j^\dagger b_j + b_i^\dagger b_i \langle b_j^\dagger b_j \rangle + \mathcal{C}$$

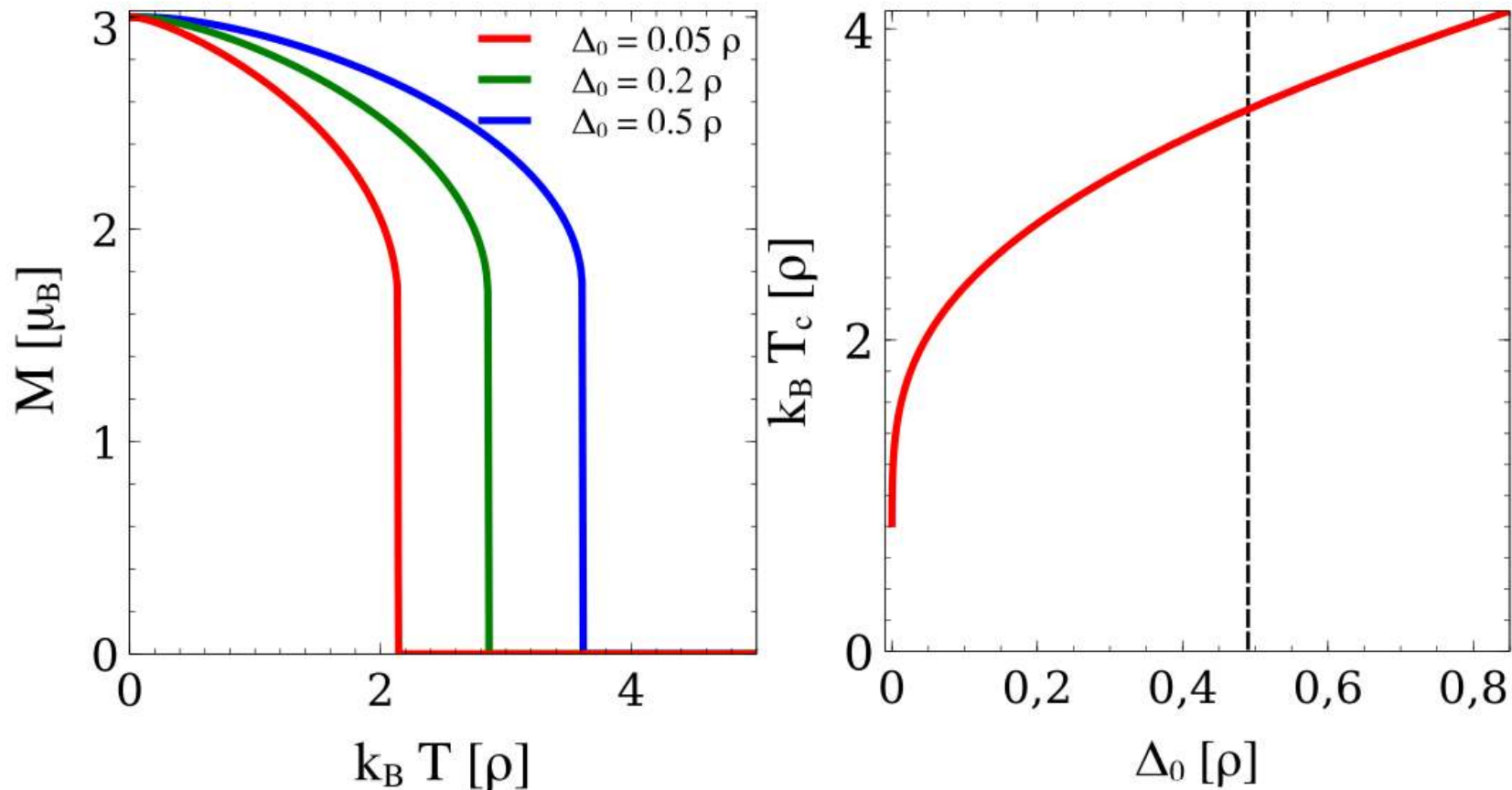
$$JS \rightarrow J(S - \langle b^\dagger b \rangle) = JM(T)$$

$$\lambda S \rightarrow \lambda(S - \langle b^\dagger b \rangle) = \lambda M(T)$$

Selfconsistent magnetization at finite temperature

$$M = S - \frac{1}{2(2\pi)^2} \int_{BZ} \frac{d^2 \vec{k}}{e^{\beta M E(\vec{k})/S} - 1}$$

The critical temperature



**The value of the spin wave gap
determines the critical temperature**

Take home

- CrI_3 is a 2d realization of the XXZ model
- Iodine is the ultimate source of magnetic anisotropy in this material

Thank you!!!