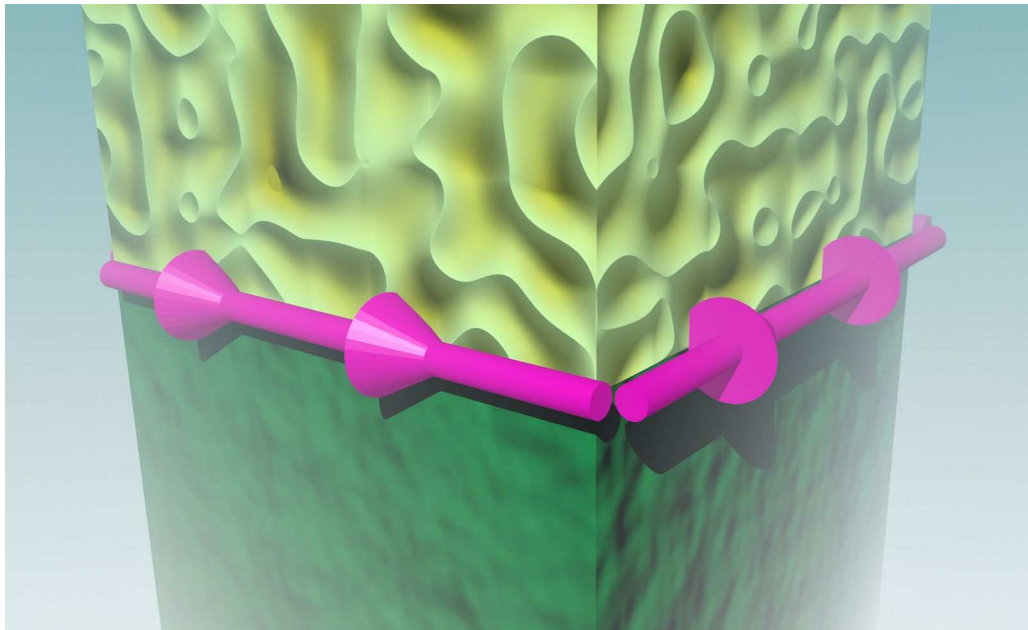


Engineering and detecting unconventional superconductivity with disguised Dirac points

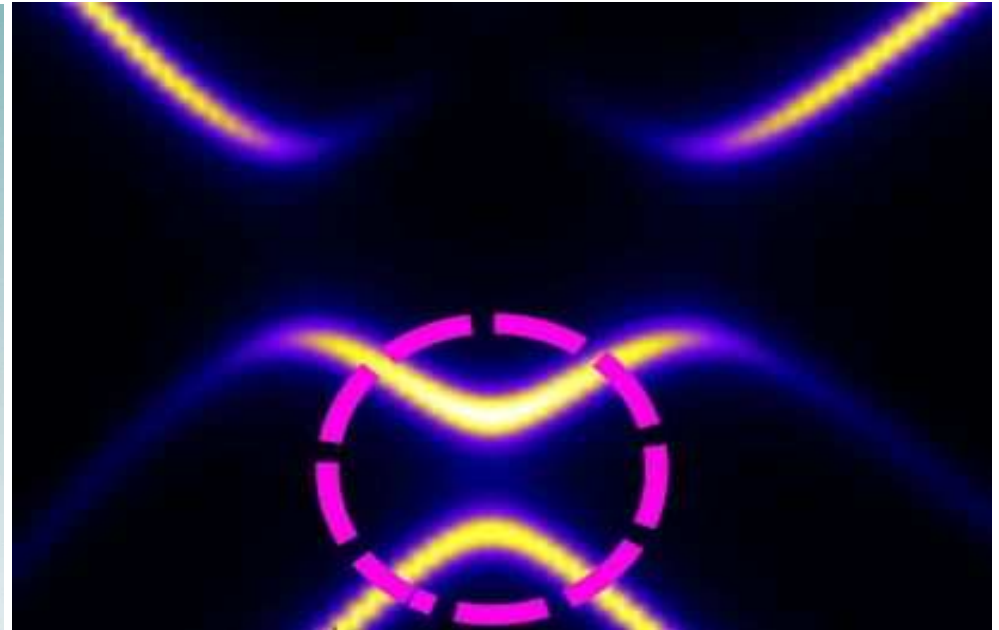
Jose Lado¹ and Manfred Sigrist²

¹*Department of Applied Physics, Aalto University, Finland*

²*Institute for Theoretical Physics, ETH Zurich, Switzerland*



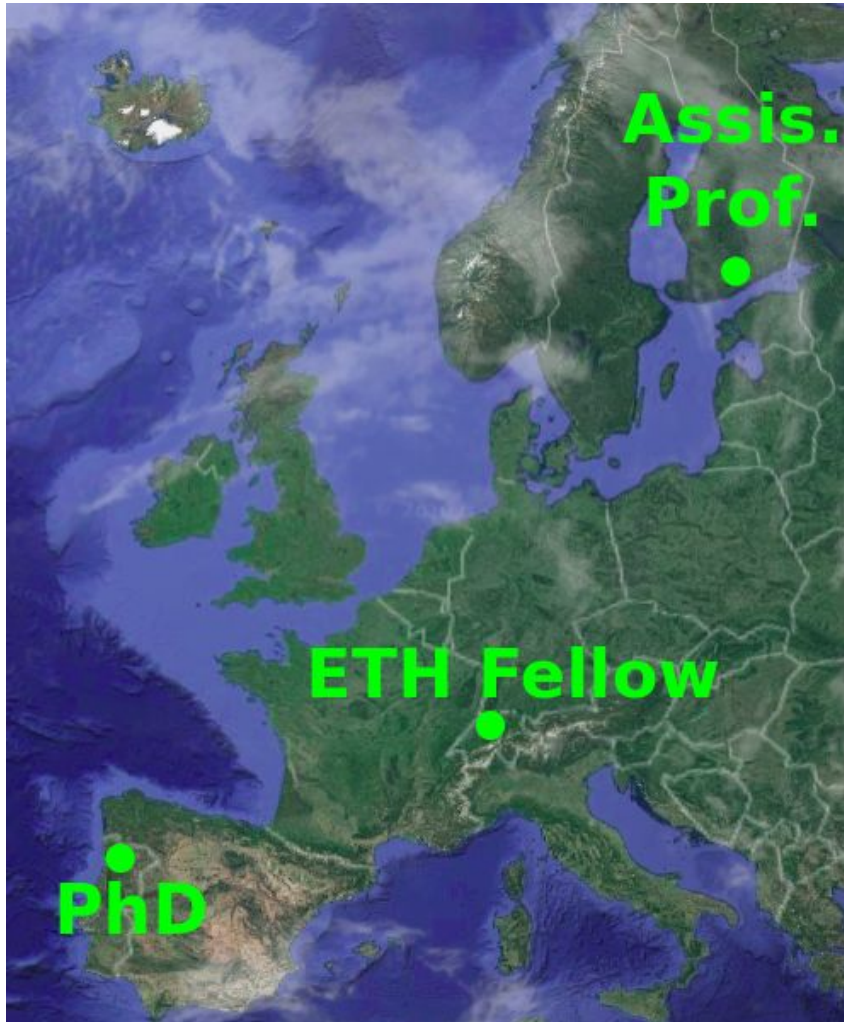
Phys. Rev. Lett. 121, 037002 (2018)



Phys. Rev. Research 1, 033107 (2019)

Radboud University, November 21st 2019

About me



Ph.D. 2013-2016, INL Braga (Portugal)

Prof. Joaquin Fernandez-Rossier



ETH Fellow 2017-2019, ETH Zurich (Switzerland)

Prof. Manfred Sigrist

Prof. Oded Zilberberg

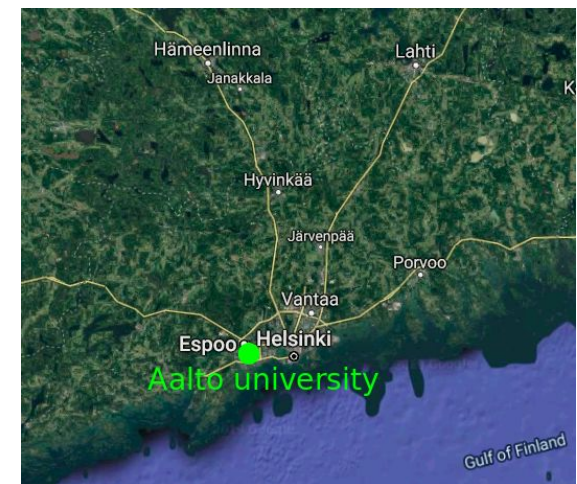


Assistant Professor, Aalto University (Sep. 2019)

Aalto University

About Aalto University, Finland

- Located in Espoo, 10km from Helsinki
- Established in 2010
- Science, business and arts
- Second largest University in Finland
- Around 17000 students and 4000 staff
- Around 20 research groups in the Department of Applied Physics



The quantum materials age

Electronics



Imaging techniques in medicine



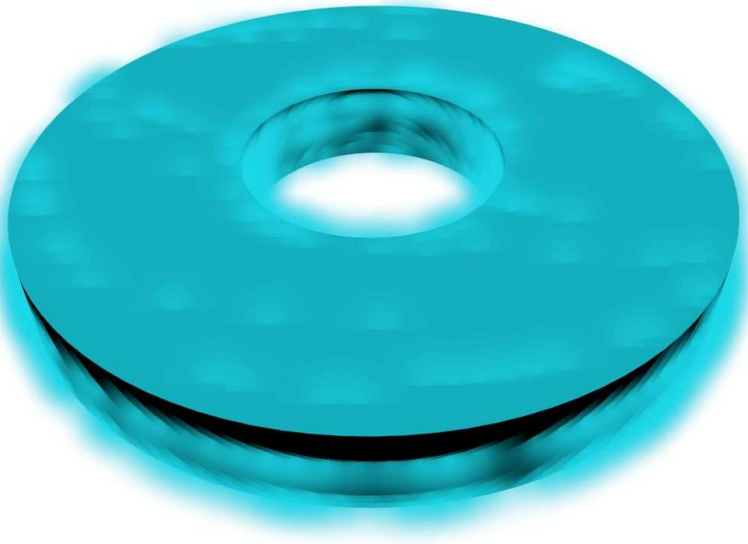
Renewable energy



Materials with exotic properties allow to tackle key technological challenges

Macroscopic materials, emergence and fundamental physics

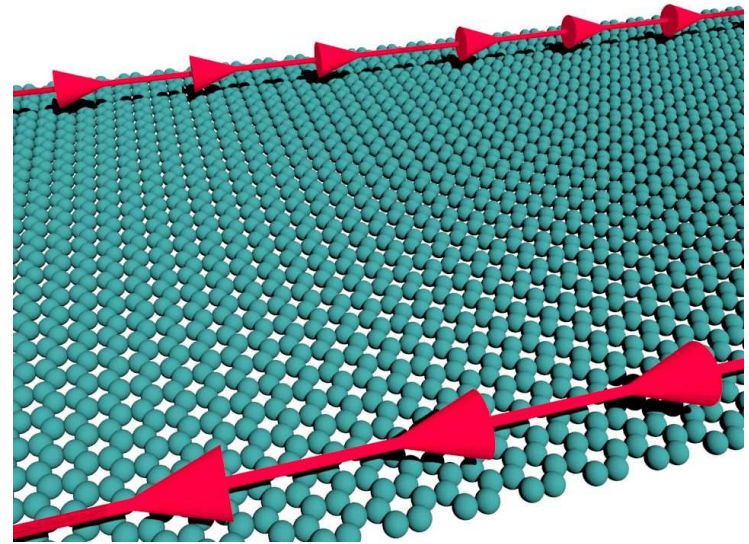
Superconductivity



$$\Phi = \frac{h}{2e} \quad \text{Many-body state}$$

Quantization of flux

Quantum Hall effect



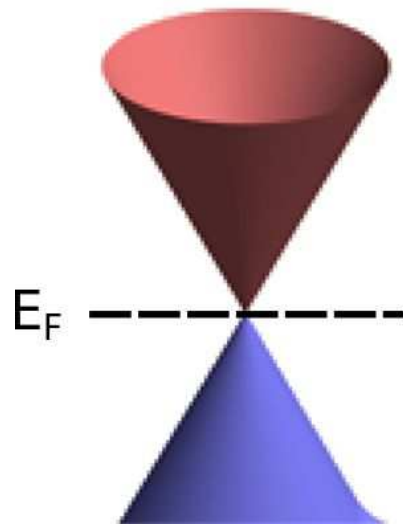
$$\sigma_{xy} = \nu \frac{e^2}{h} \quad \text{Single-particle state}$$

Quantization of conductance

Macroscopic quantum phenomena determine fundamental constants with the highest precision

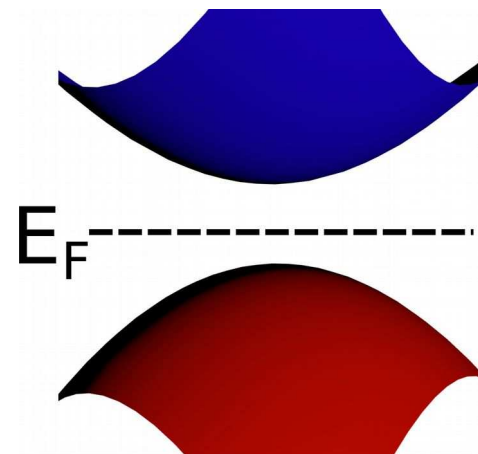
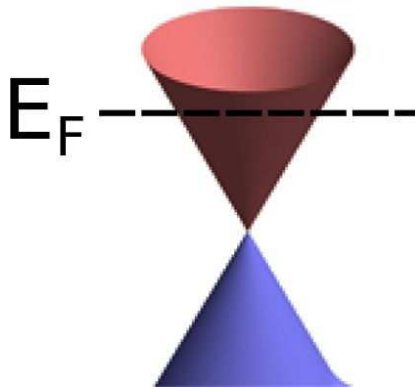
On Dirac crossings

A platform to engineer exotic physics in solid state materials



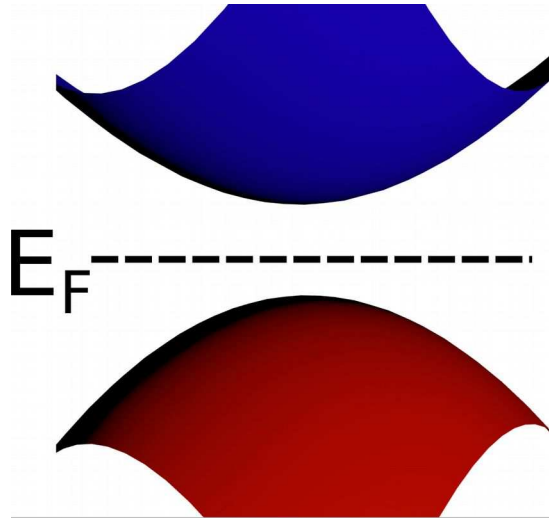
WINDING NUMBER
TOPOLOGY **NODAL LINES**
TWISTED BILAYER GRAPHENE
DIRAC
BERRY FLUX
CHERN NUMBER
BERRY PHASE
Weyl POINTS
NODAL CHAINS
SPINOR
ANOMALOUS HALL

But what if the Dirac points are not observable at the chemical potential?



Exploiting disguised Dirac points

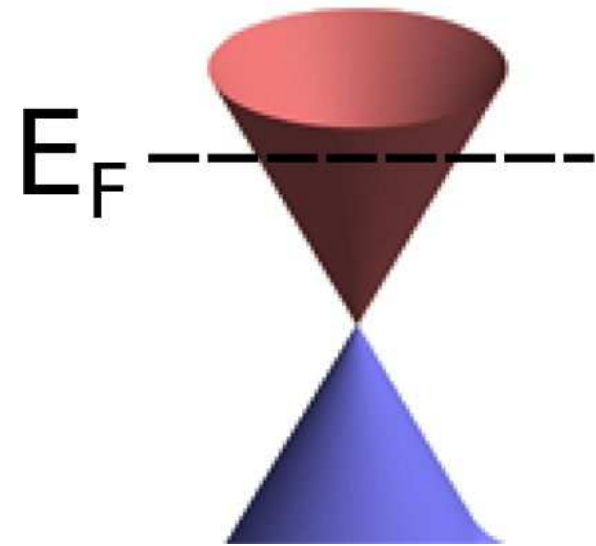
**Platform for an artificial
topological superconductor**



Gaped out by antiferromagnetism

Phys. Rev. Lett. 121, 037002 (2018)

**Platform to detect
non-unitary superconductivity**



Away from the Fermi energy

Phys. Rev. Research 1, 033107 (2019)



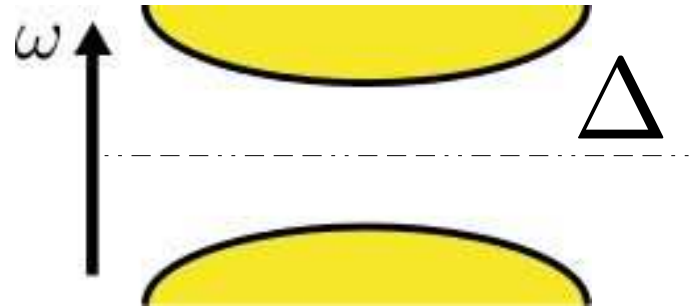
Manfred
Sigrist

**Creating a topological
superconductor with
antiferromagnetically
gaped Dirac points**

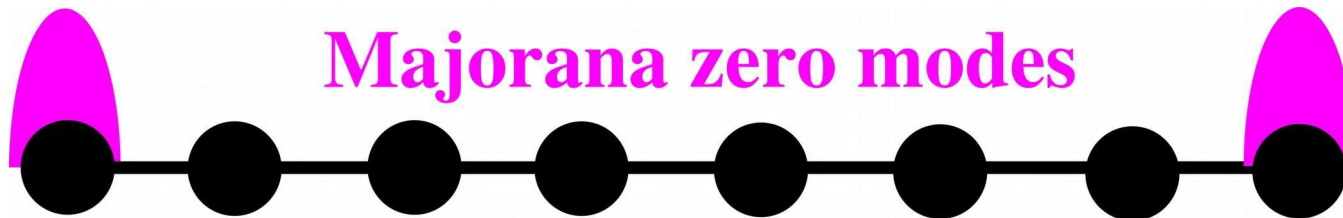
Topological superconductors

Topological superconductors with broken time reversal symmetry

Gapped bulk excitations

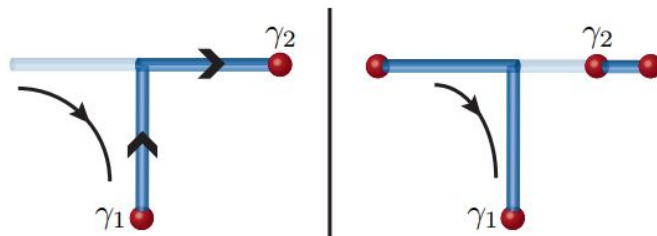


Gapless surface modes, single Majorana mode per edge



A. Y. Kitaev.
Physics-Uspekhi,
44:131, 2001

Anyonic statistics, suitable for quantum computing



Nature Physics 7, 412-417 (2011)

Our objective

Build a 2D topological superconductor with

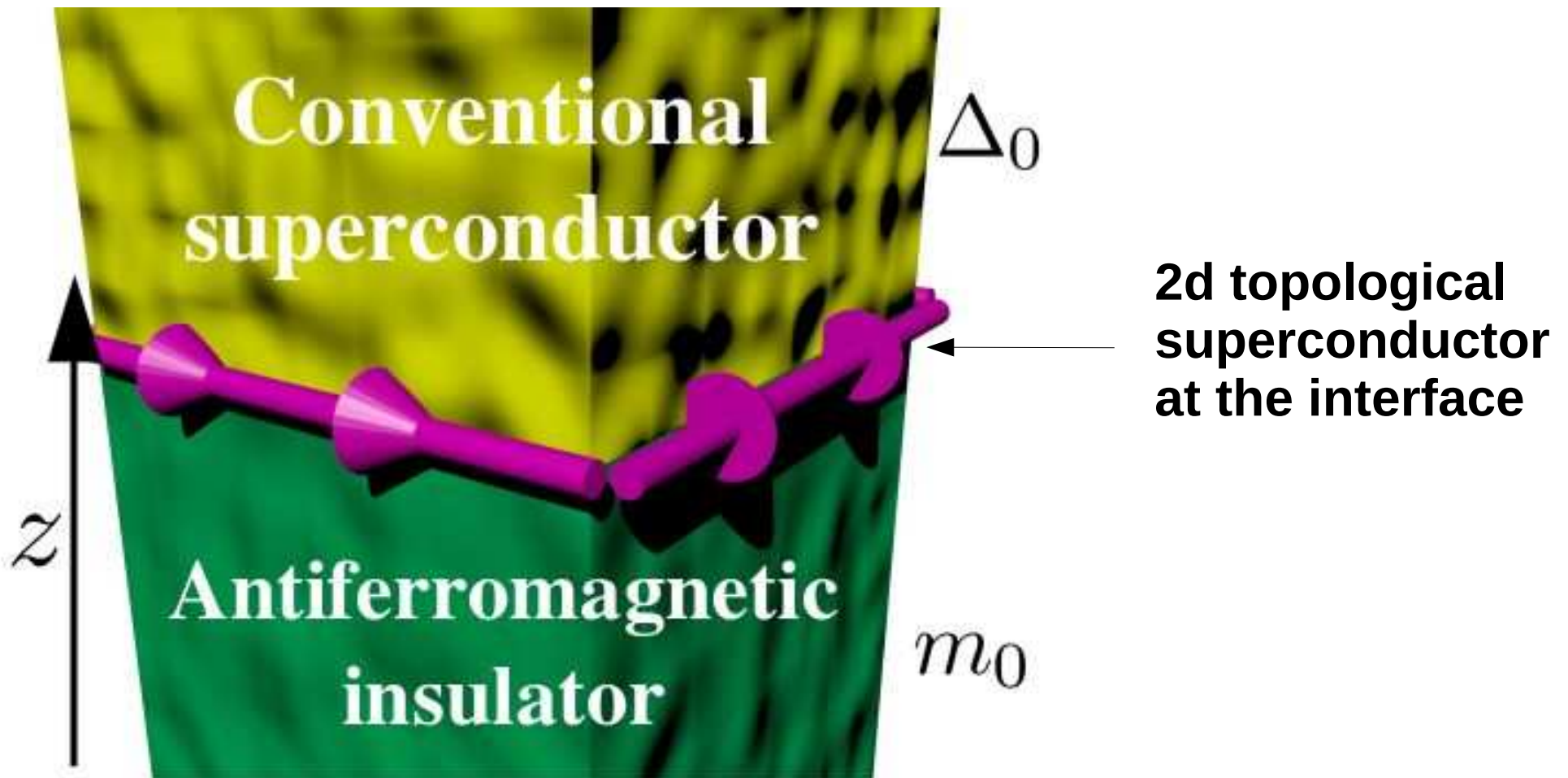
- A conventional (s-wave) 3D superconductor
- A 3D antiferromagnetic insulator

The prize



Bringing a new solid state platform to realize
artificial topological superconductors

The system for TS in an AF insulator



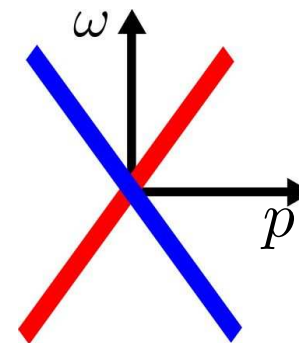
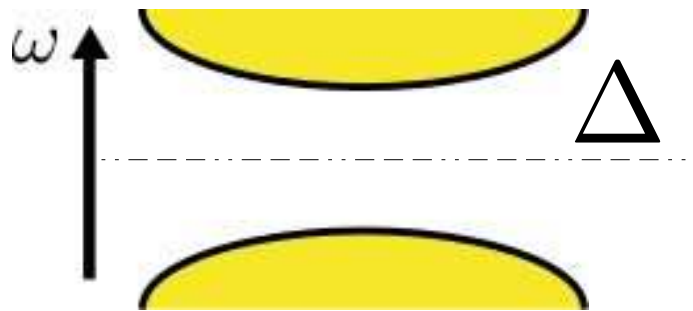
How to do your own topological superconductor



Ingredients

- s-wave pairing
- Helical states

Science 346.6209 (2014)
Science 354.6319 (2016)
Phys. Rev. Lett. 100, 096407 (2008)



Objective: realize a spin-less superconductor

$$H = \sum_n t c_{n+1}^\dagger c_n + \Delta c_n c_{n+1} + c.c.$$

A. Y. Kitaev. *Physics-Uspekhi*, 44:131, 2001

The initial problem

How can we get a topological phase starting from a trivial insulator?



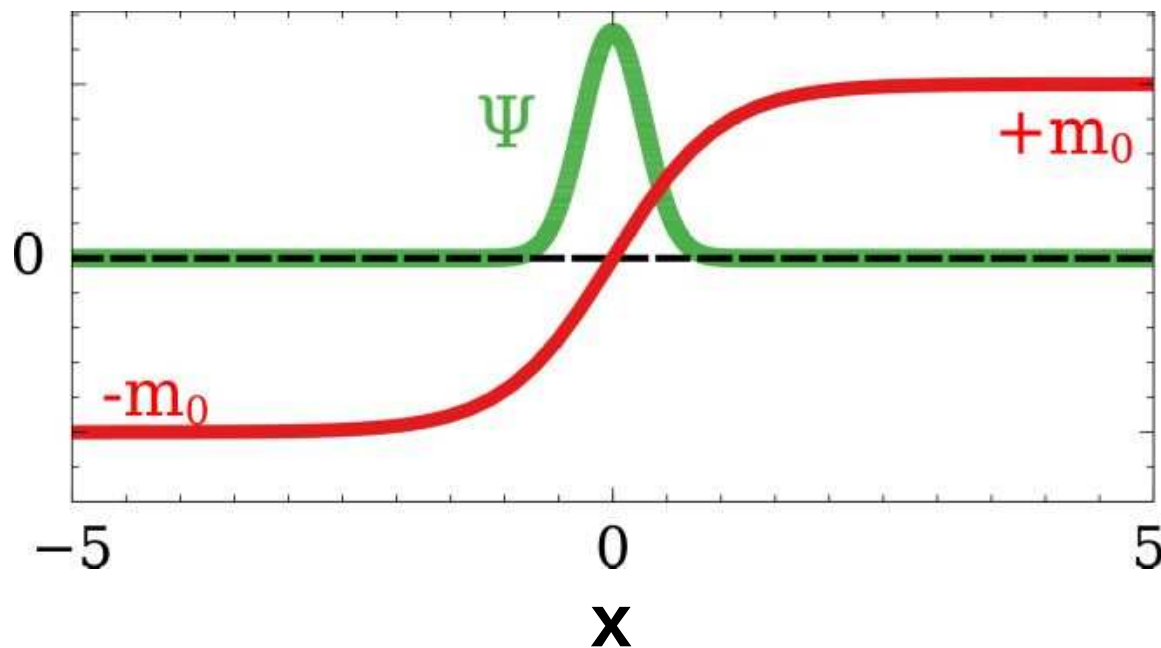
We need to create a “spinless” gapless state out of an insulator

Building an interface gas between two insulators

Lets take a 1d Dirac equation with spatially dependent mass

$$H = \begin{pmatrix} m(x) & p \\ p & -m(x) \end{pmatrix} = \sigma_x p + m(x) \sigma_z$$

$$m(\infty) = -m(-\infty) = m_0$$



**Zero energy solution
bound to the interface**

$$\Psi = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} e^{-\int_0^x m(\lambda) d\lambda}$$

$$H\Psi = 0$$

*R. Jackiw and C. Rebbi
Phys. Rev. D 13, 3398 (1976)*

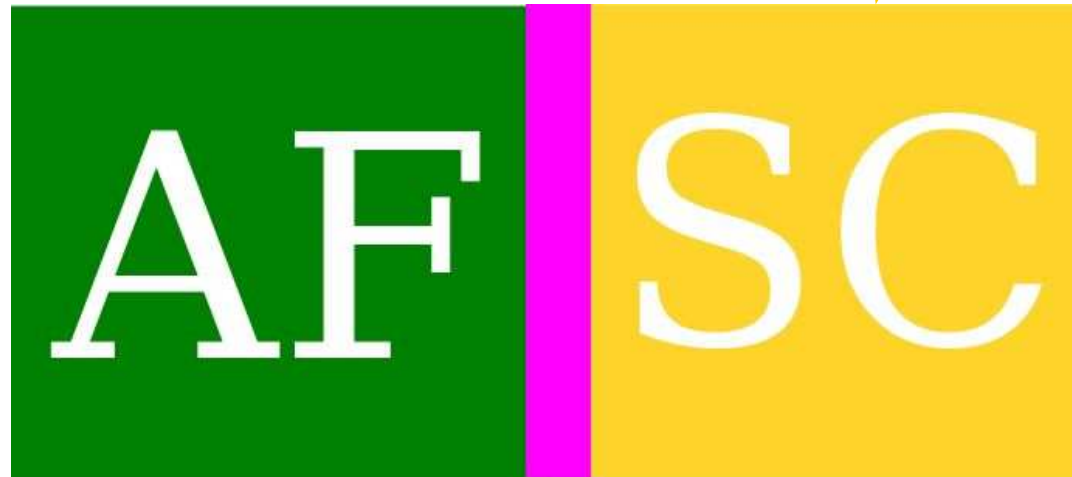
Interface modes between a 1D AF and a 1D SC

Kinetic part

Antiferromagnetic order

Superconducting order

$$H = \hat{H}_{kin} + H_{AF} + H_{SC}$$



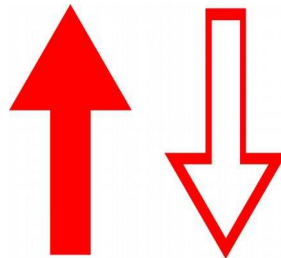
Solitonic modes between Dirac AF and SC

Total Hamiltonian

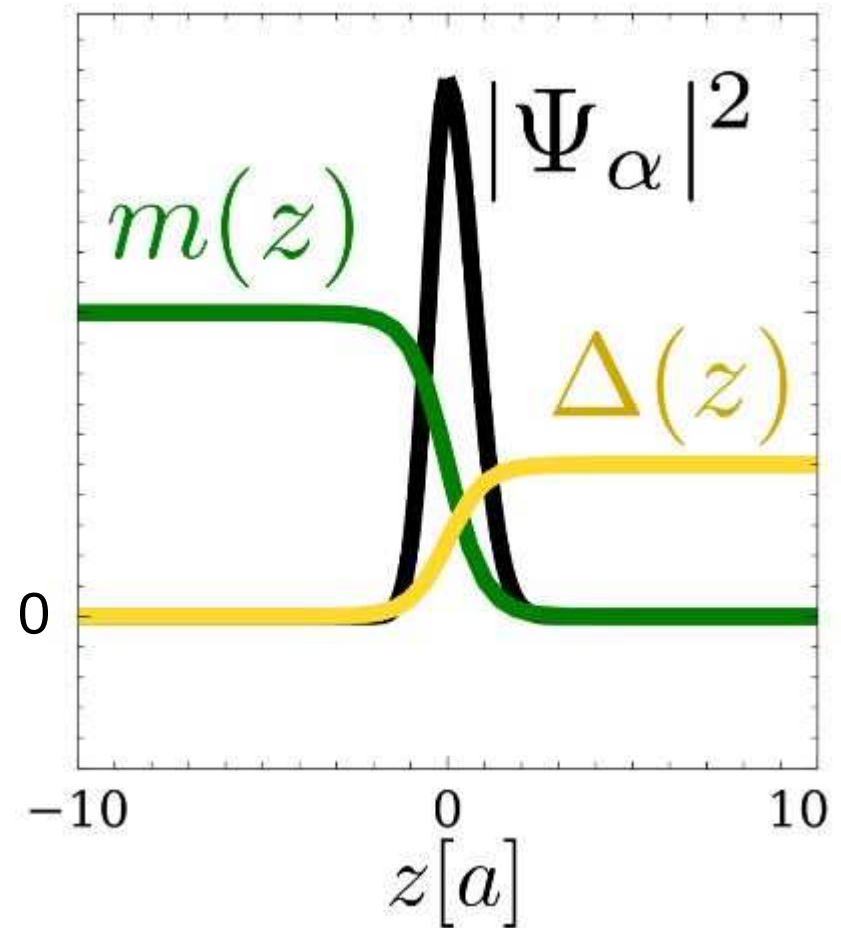
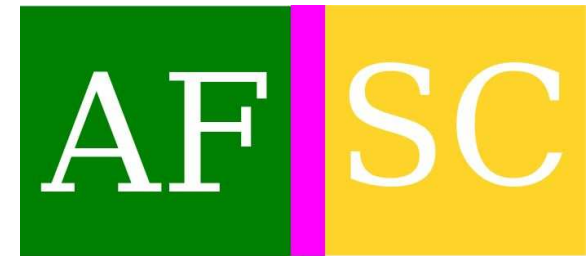
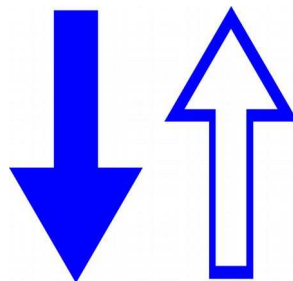
$$H = H_{kin} + H_{AF} + H_{SC}$$

There will be two zero solutions

Sector #1
Up electron, down hole

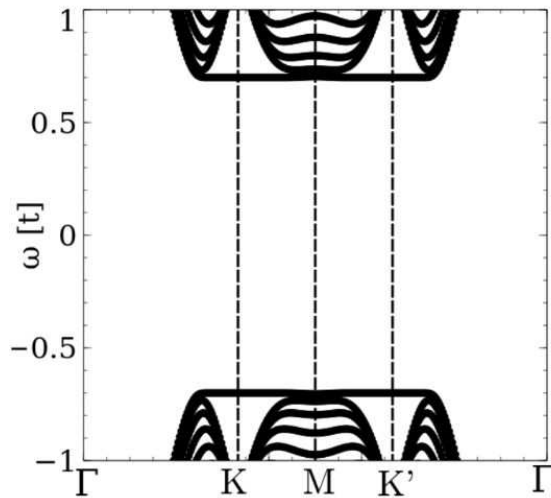


Sector #2
Down electron, up hole

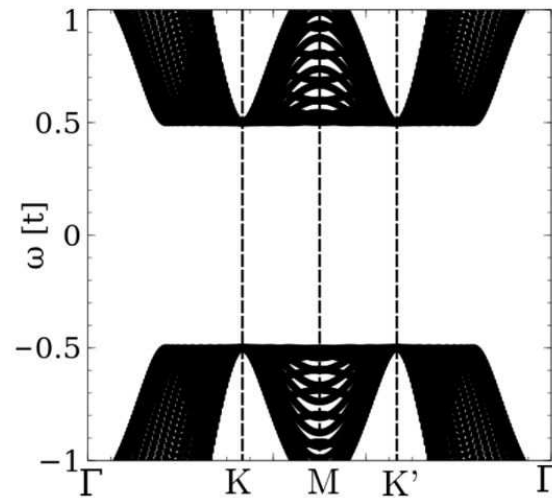


Emergence of interfacial modes, no spin-orbit coupling

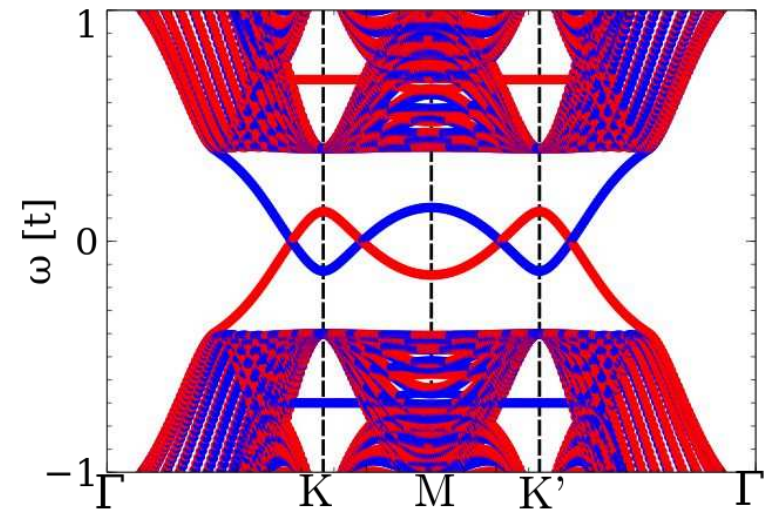
Antiferromagnet



Superconductor



AF/SC heterostructure

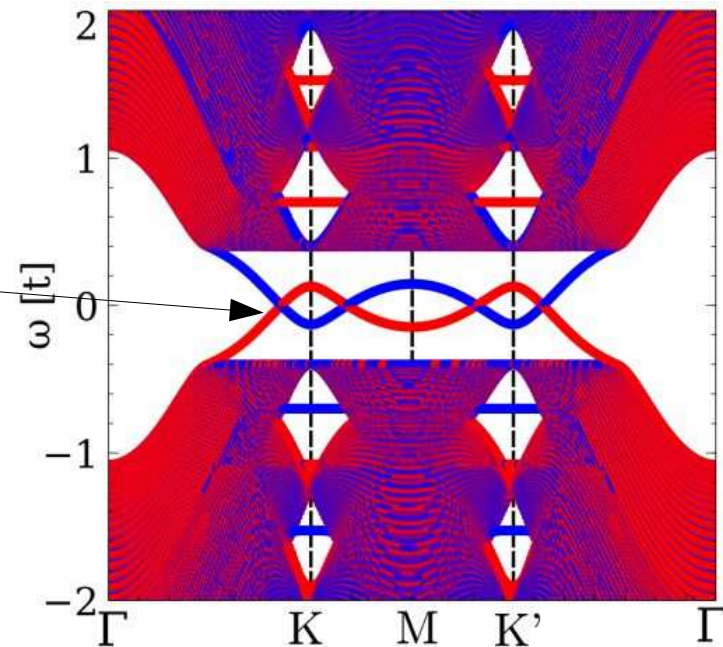
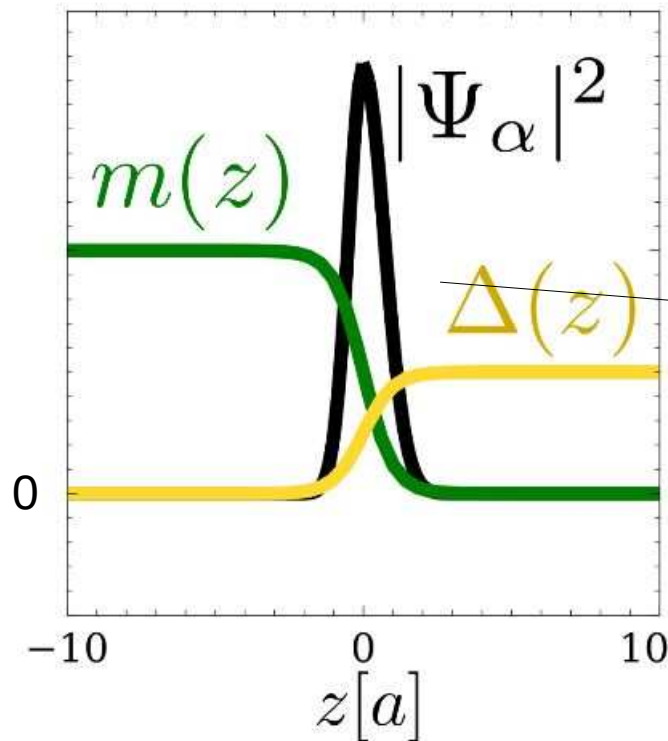
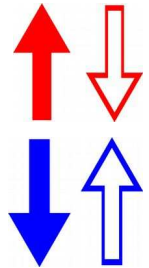


Interface states between AF and SC

Superconducting solitonic excitations

$$\Psi_1(z) = g(z)[c_{A,\vec{k}_{\parallel},\uparrow} - ic_{B,\vec{k}_{\parallel},\uparrow} - c_{A,-\vec{k}_{\parallel},\downarrow}^\dagger - ic_{B,-\vec{k}_{\parallel},\downarrow}^\dagger]$$

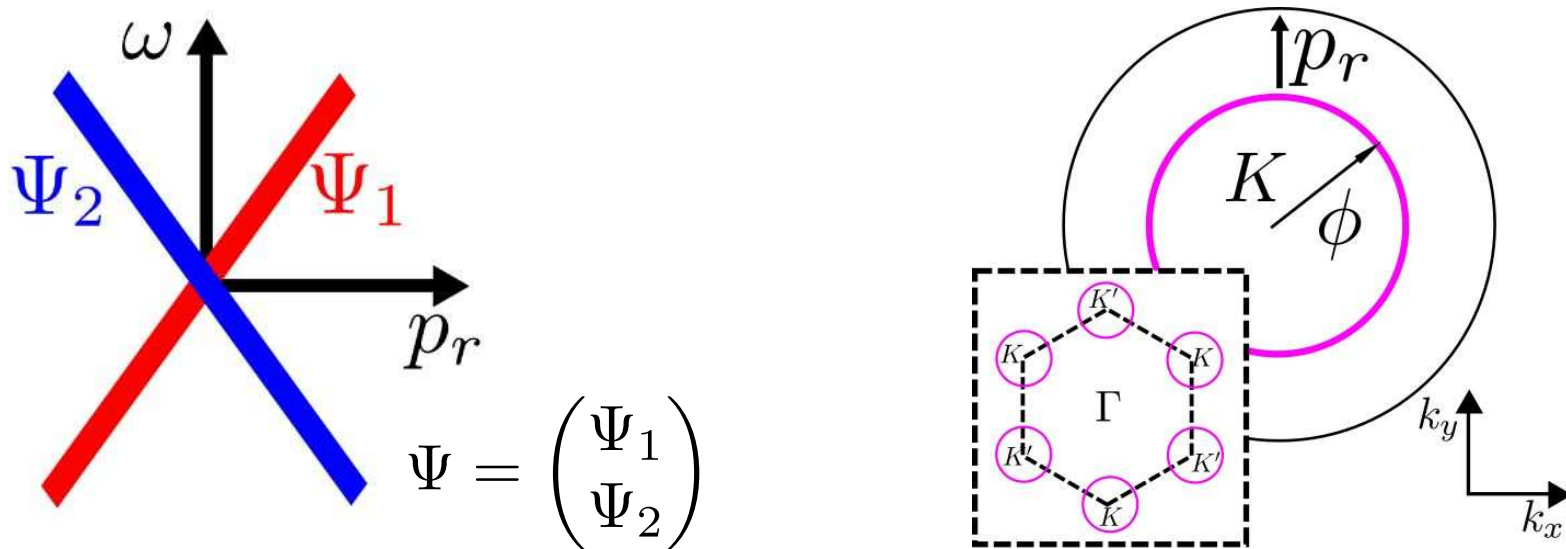
$$\Psi_2(z) = g(z)[c_{A,\vec{k}_{\parallel},\downarrow} + ic_{B,\vec{k}_{\parallel},\downarrow} - c_{A,-\vec{k}_{\parallel},\uparrow}^\dagger + ic_{B,-\vec{k}_{\parallel},\uparrow}^\dagger]$$



Effective model without spin-orbit coupling

Effective Hamiltonian without spin-orbit coupling

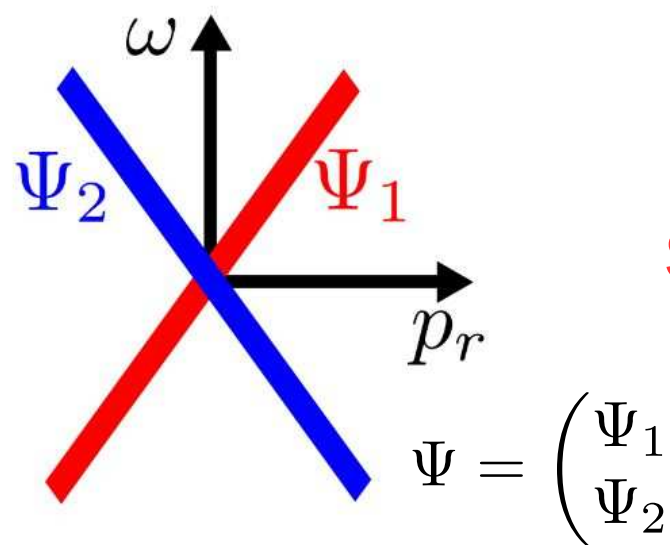
$$\mathcal{H}(p_r, \phi) = \Psi^\dagger \begin{pmatrix} v_r p_r & 0 \\ 0 & -v_r p_r \end{pmatrix} \Psi$$



Every point of the Dirac line creates a zero mode

Effective model with spin-orbit coupling

Effective Hamiltonian with spin-orbit coupling

$$\mathcal{H}(p_r, \phi) = \Psi^\dagger \begin{pmatrix} v_r p_r & -i\lambda e^{i\phi} \\ i\lambda e^{-i\phi} & -v_r p_r \end{pmatrix} \Psi$$


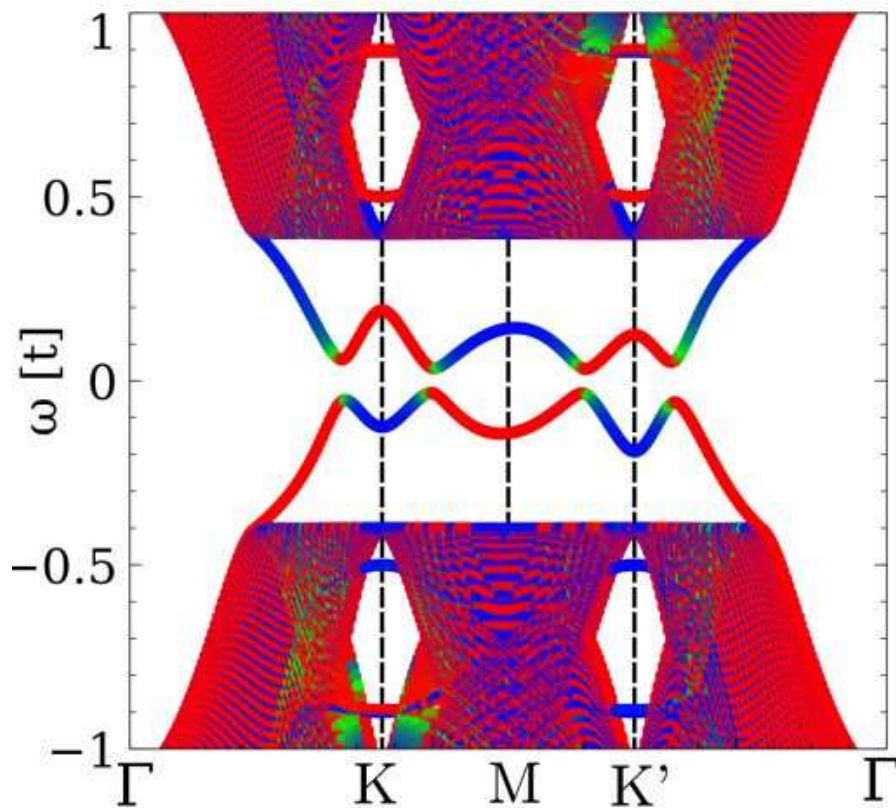
Spin-orbit coupling takes the form of a chiral pairing

$$\Psi = \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix}$$

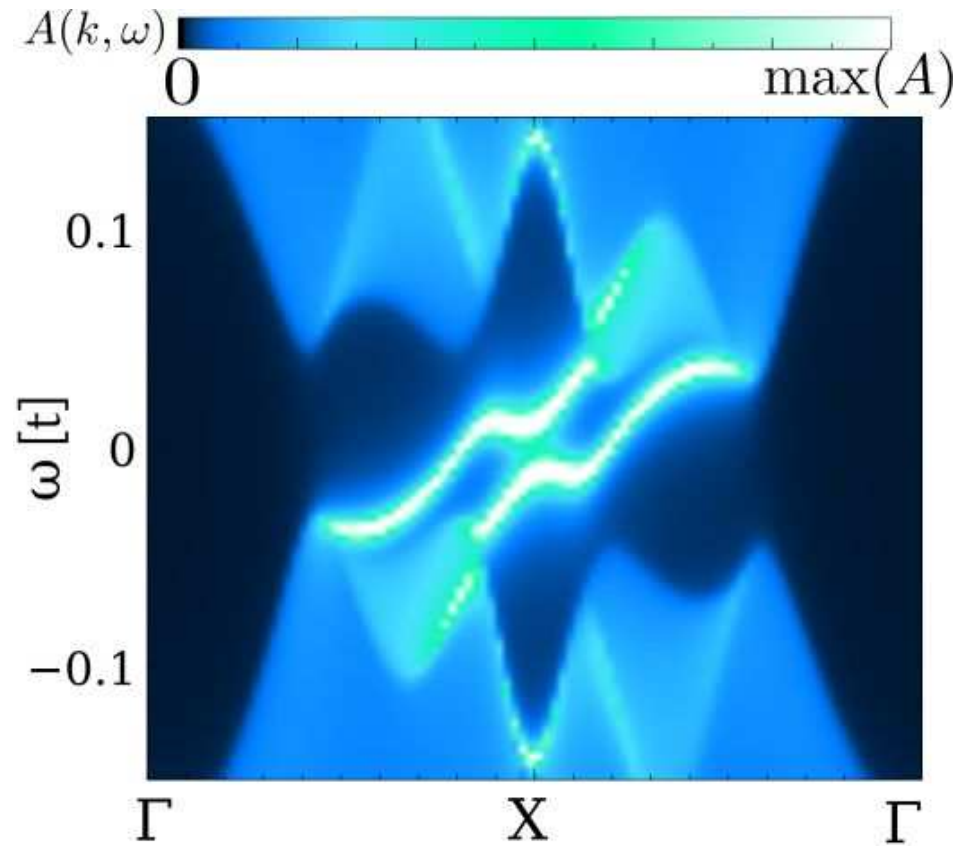
Chern number $\mathcal{C} = \pm 1$ per Dirac line

Adding spin-orbit coupling

Band structure

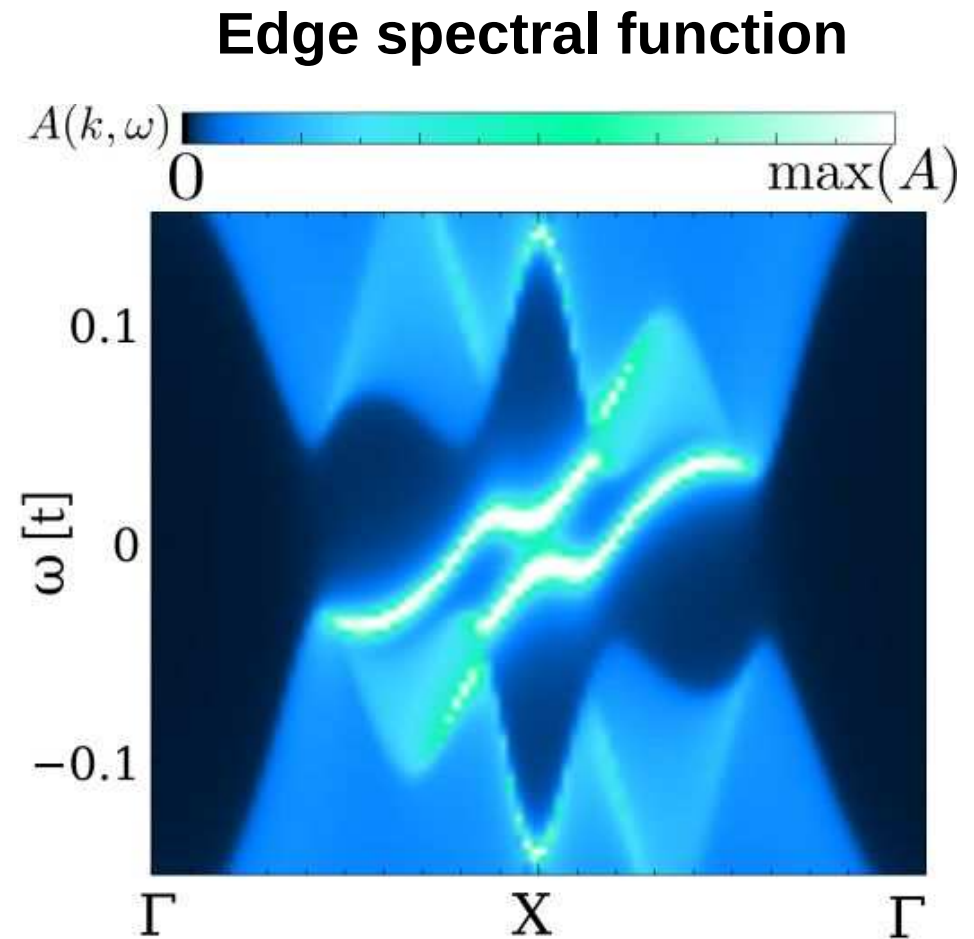
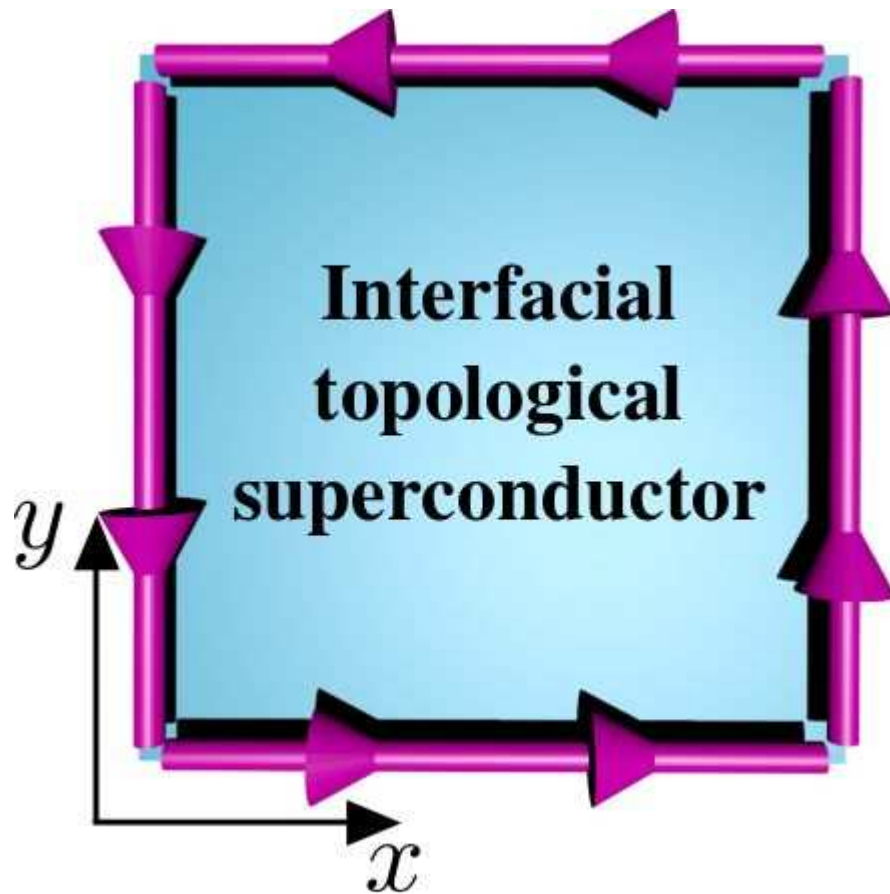


Edge spectral function



Two co-propagating edge modes

An interfacial topological superconductor

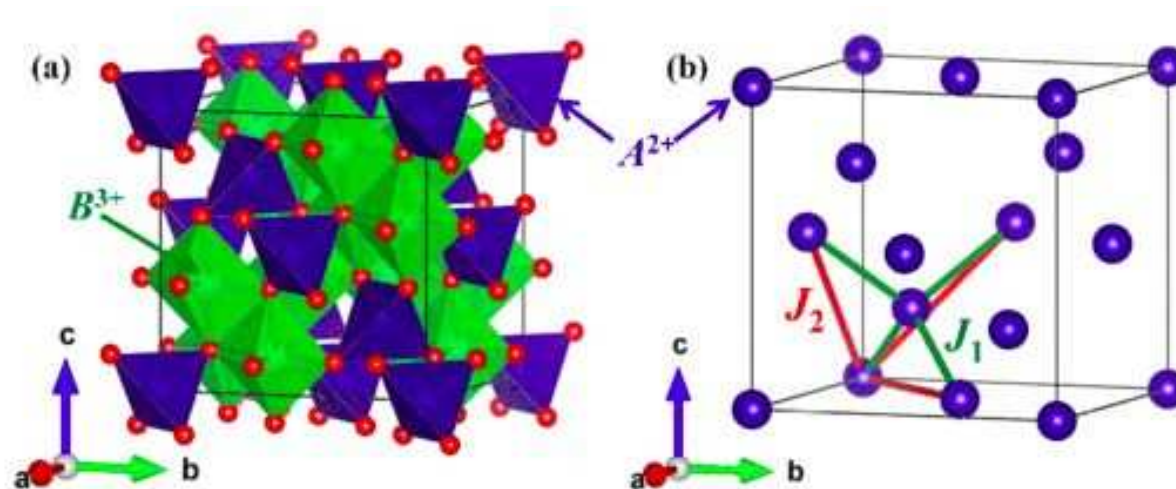


The interface realizes a topological superconductor

Possible materials, spinels

- Antiferromagnet forming a diamond lattice

Antiferromagnetic spinels

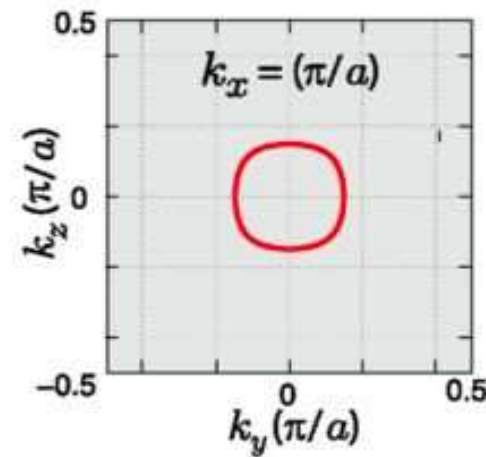
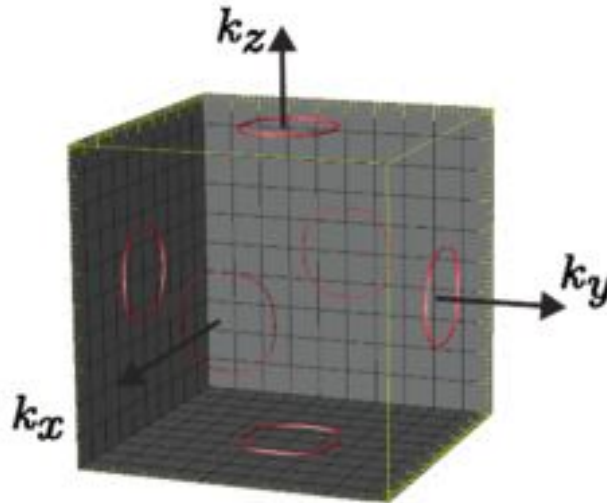


Co atoms form a diamond lattice

Possible materials, Dirac materials in the AF phase

Dirac lines in the absence of spin-orbit coupling and magnetism

Phys. Rev. Lett. 115, 036806 (2015)

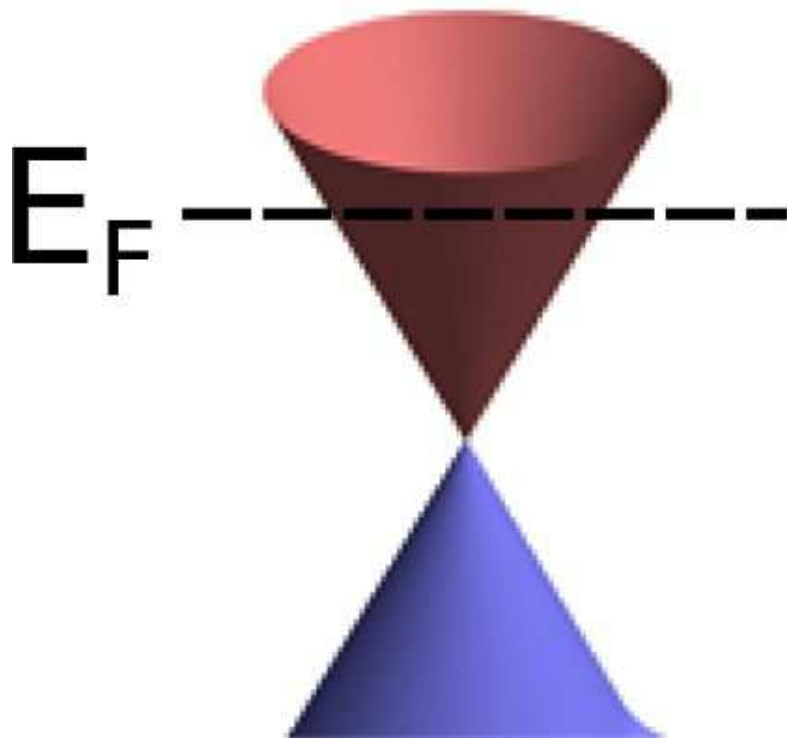


Antiferromagnets whose paramagnetic state hosts Dirac lines

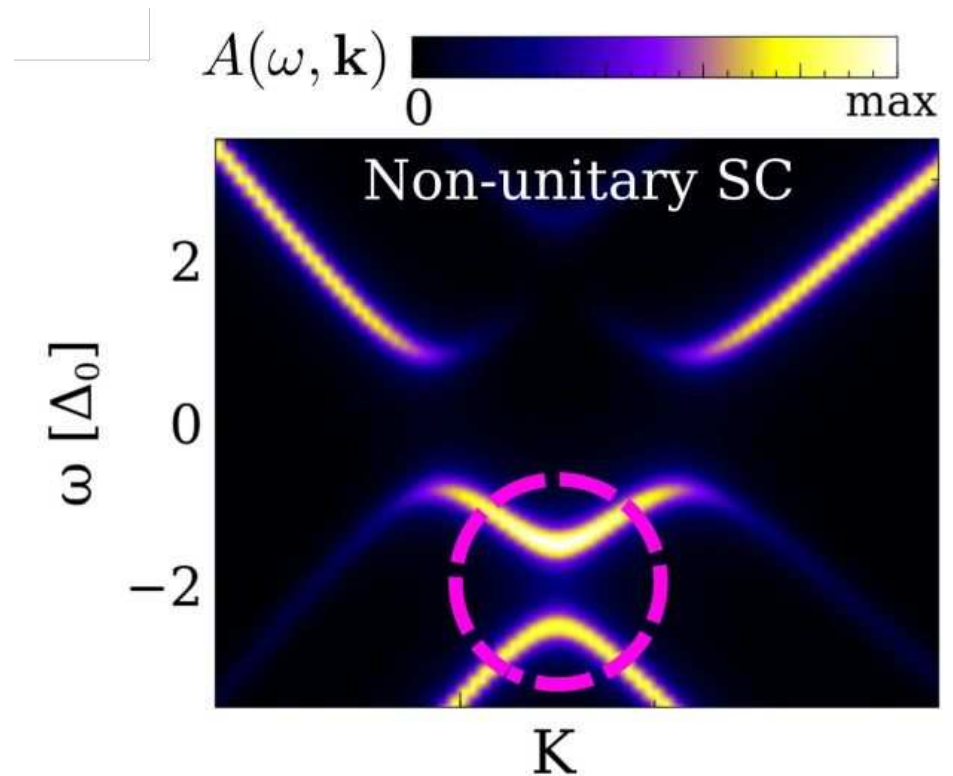
**Detecting non-unitary
superconductivity through
deep Dirac points**

Superconducting symmetry from ARPES experiments

Normal state band structure



ARPES in superconducting state



Can we detect non-unitary superconductivity from an ARPES experiment?

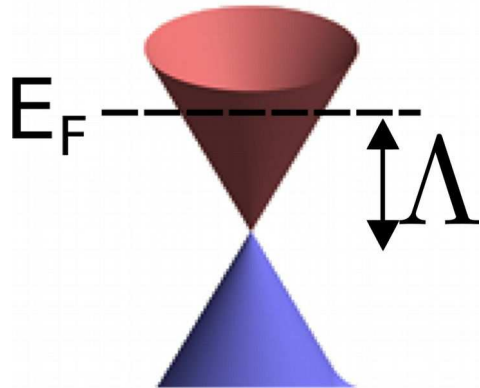
Gap opening by singlet multiorbital superconductivity

Two orbital band structure with a Dirac crossing, singlet pairing

Hamiltonian in normal state

$$\psi_{\mathbf{k}}^\dagger = (c_{\alpha,\mathbf{k}}^\dagger, c_{\beta,\mathbf{k}}^\dagger);$$

$$H_0^{DP}(\mathbf{k}) = -\Lambda\tau_0 + \tau_x k_x + \tau_y k_y$$



Hamiltonian in SC state

$$\Psi_{\mathbf{k}}^\dagger = (c_{\alpha_0,\uparrow,\mathbf{k}}^\dagger, c_{\beta_0,\uparrow,\mathbf{k}}^\dagger, c_{\alpha_0,\downarrow,-\mathbf{k}}, c_{\beta_0,\downarrow,-\mathbf{k}}).$$

$$H_{BdG}(\mathbf{k}) = \begin{pmatrix} H_0^{DP}(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta(\mathbf{k}) & -H_0^{DP}(-\mathbf{k}) \end{pmatrix}$$

Effective Hamiltonian after Schrieffer-Wolff transformation

$$H^{DP} = H_0^{DP} + \frac{\Delta\Delta^\dagger}{2\Lambda}$$

SC gap opening at the Dirac point if $\Upsilon = \text{Tr}[\tau_z \Delta\Delta^\dagger] \neq 0$

Topological invariant of the gap opening

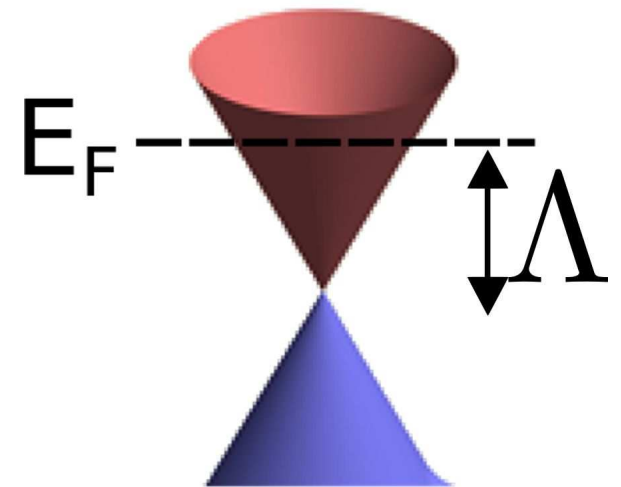
Full Hamiltonian

$$H_{BdG}(\mathbf{k}) = \begin{pmatrix} H_0^{DP}(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta(\mathbf{k}) & -H_0^{DP}(-\mathbf{k}) \end{pmatrix}$$

Effective Hamiltonian

$$H^{DP} = -\bar{\Lambda}\tau_0 + \bar{p}_x\tau_x + \bar{p}_y\tau_y + \gamma_z\tau_z$$

$$\gamma_z = \Upsilon/4\Lambda$$



SC gap opening at the Dirac point if $\Upsilon = \text{Tr}[\tau_z \Delta \Delta^\dagger] \neq 0$

Berry flux (non-integer Chern number) associated to the Dirac crossing

$$C = \frac{1}{2} \text{sign}(\Upsilon) \text{sign}(\Lambda)$$

Conditions for a Dirac gap opening

$$\Upsilon = \text{Tr}(\tau_z \Delta \Delta^\dagger)$$

No gap opening $\Upsilon = 0$

Unitary states

Intraorbital

$$\Delta \sim \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Interorbital

$$\Delta \sim \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

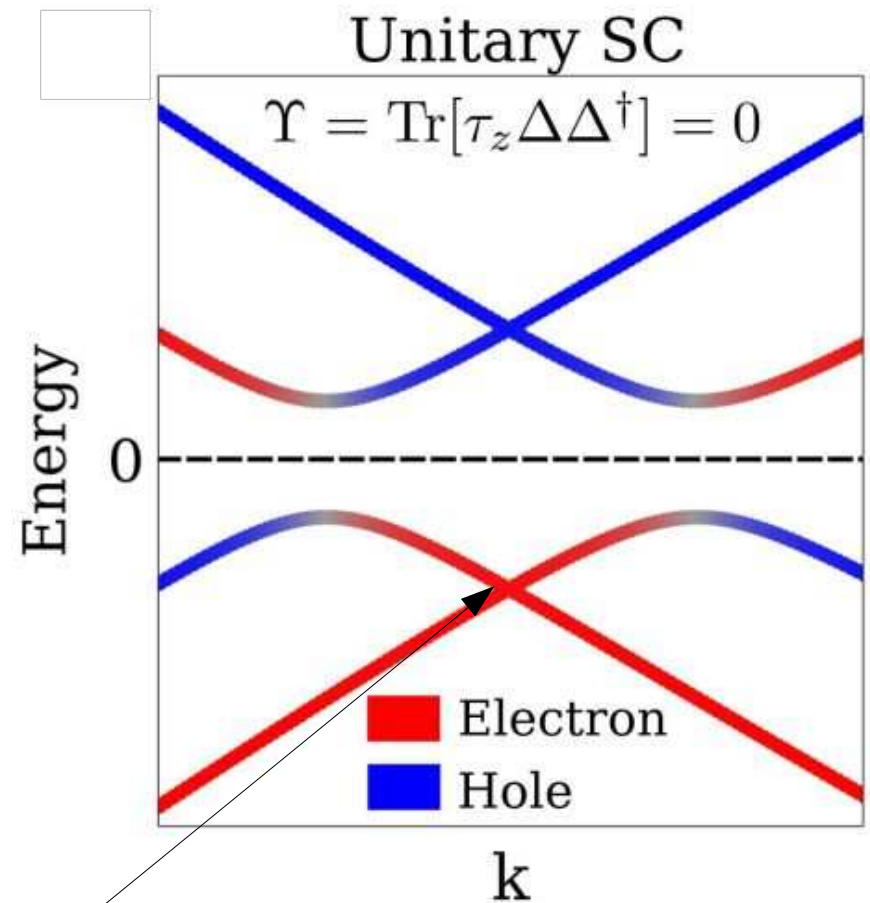
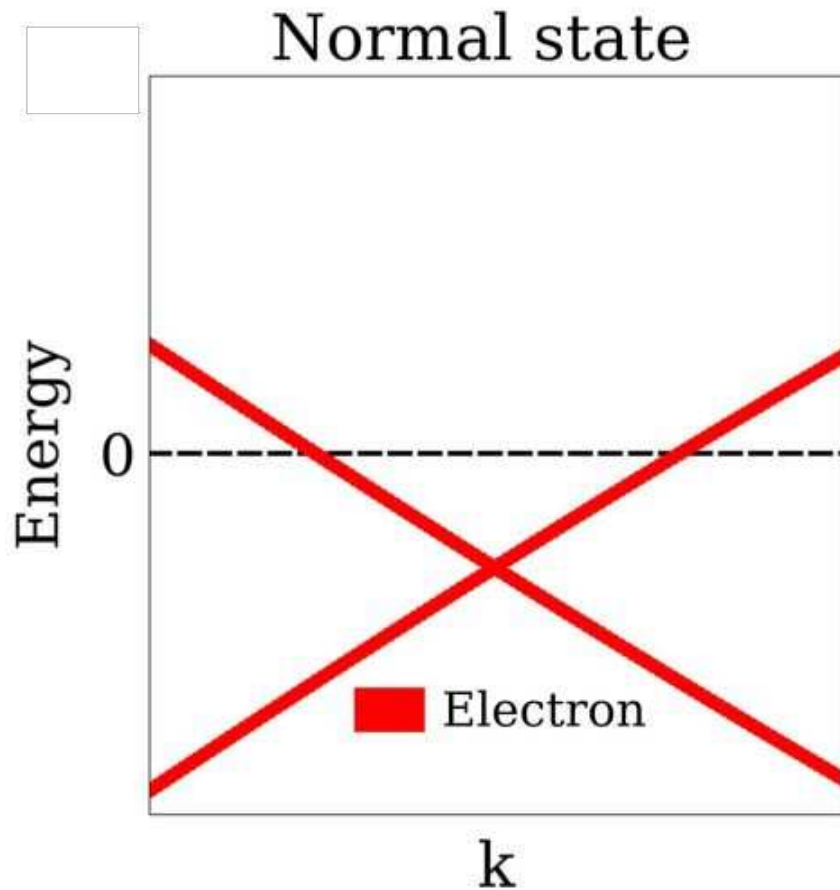
Gap opening $\Upsilon \neq 0$

Non-unitary states

Non-unitary intraorbital

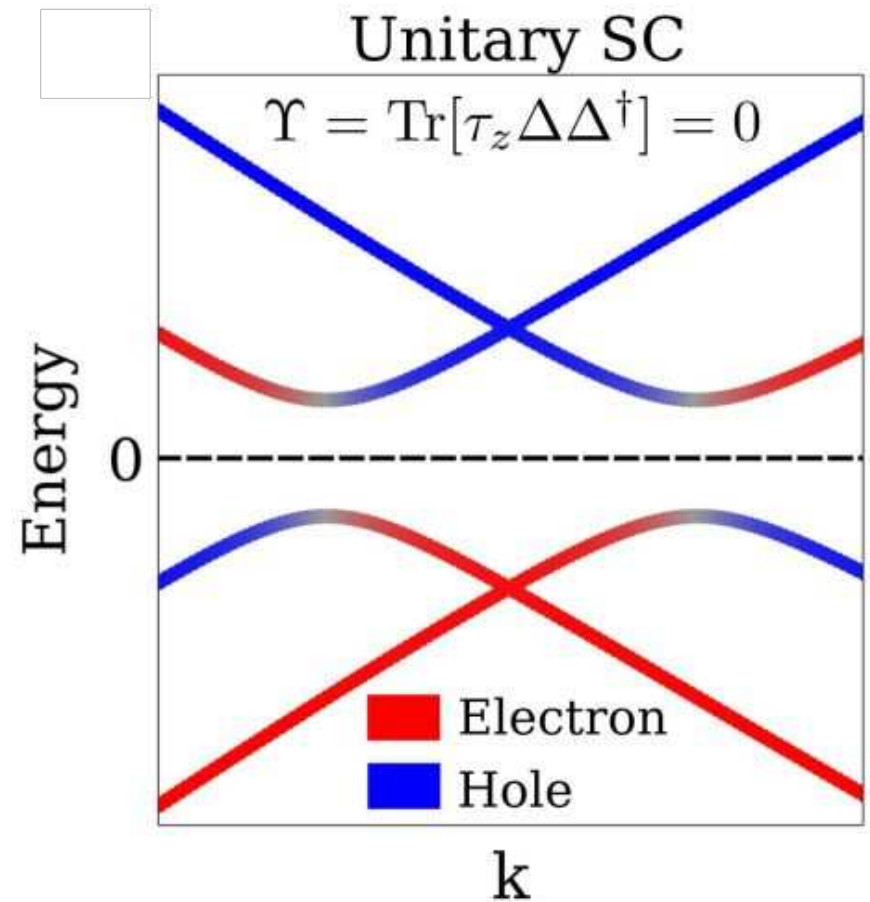
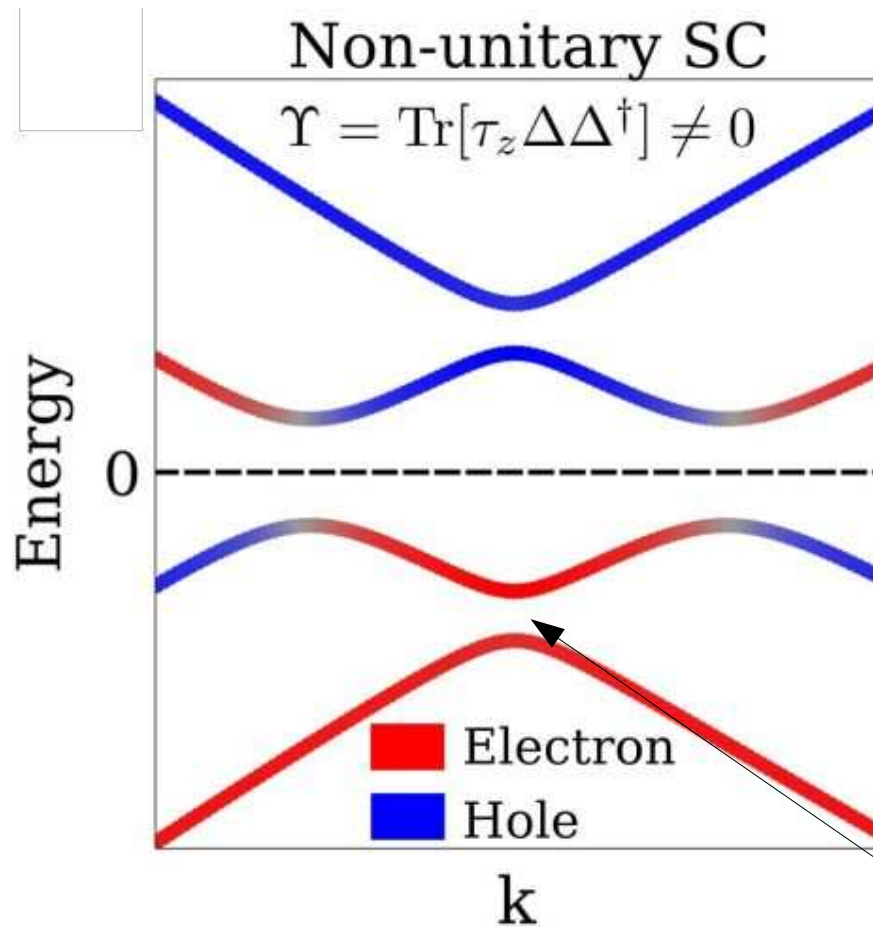
$$\Delta \sim \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Dirac crossing in the superconducting state



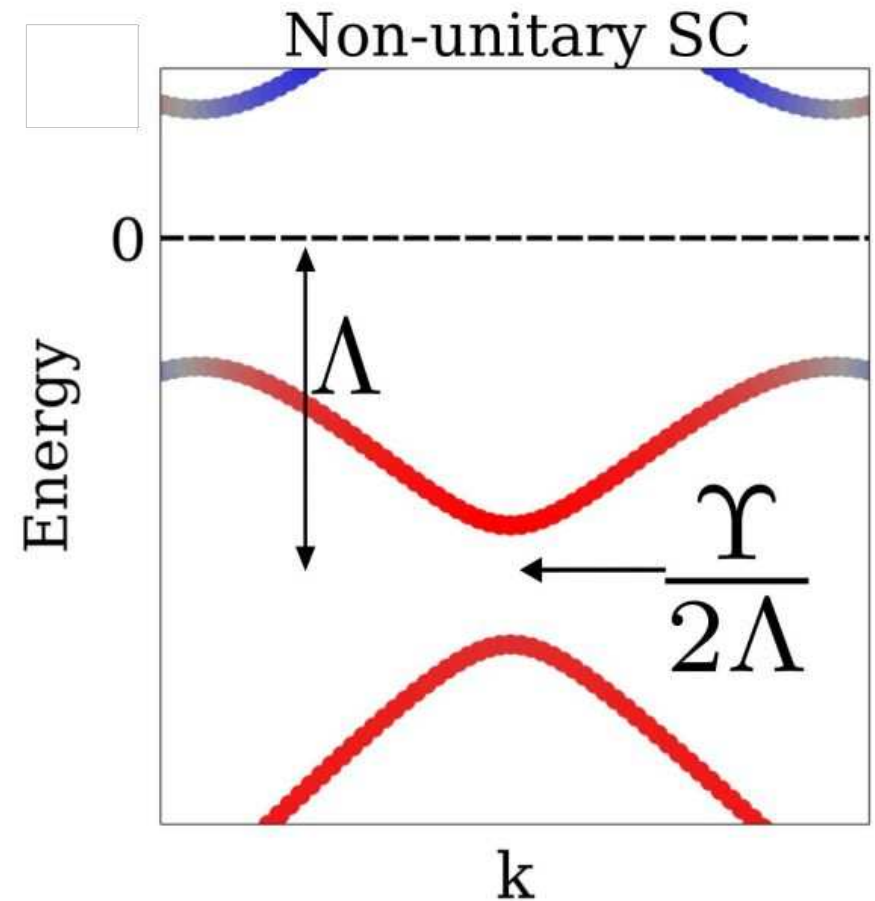
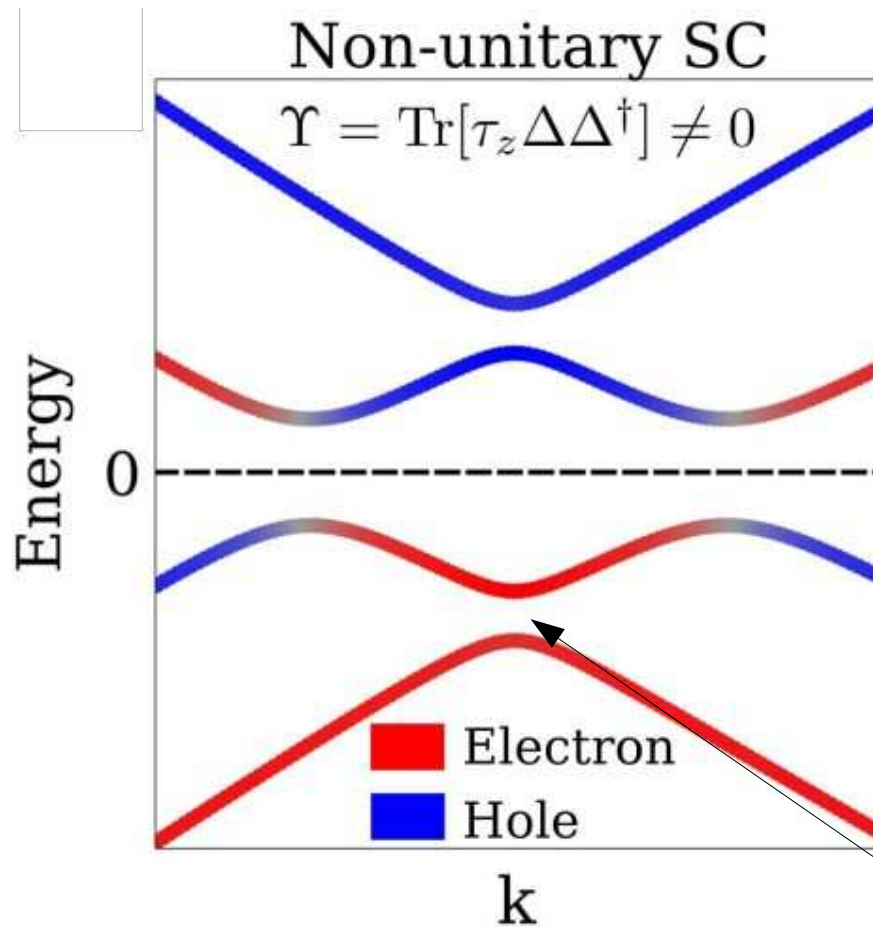
Conventional superconducting state: Dirac point remains closed

Dirac crossing in the superconducting state



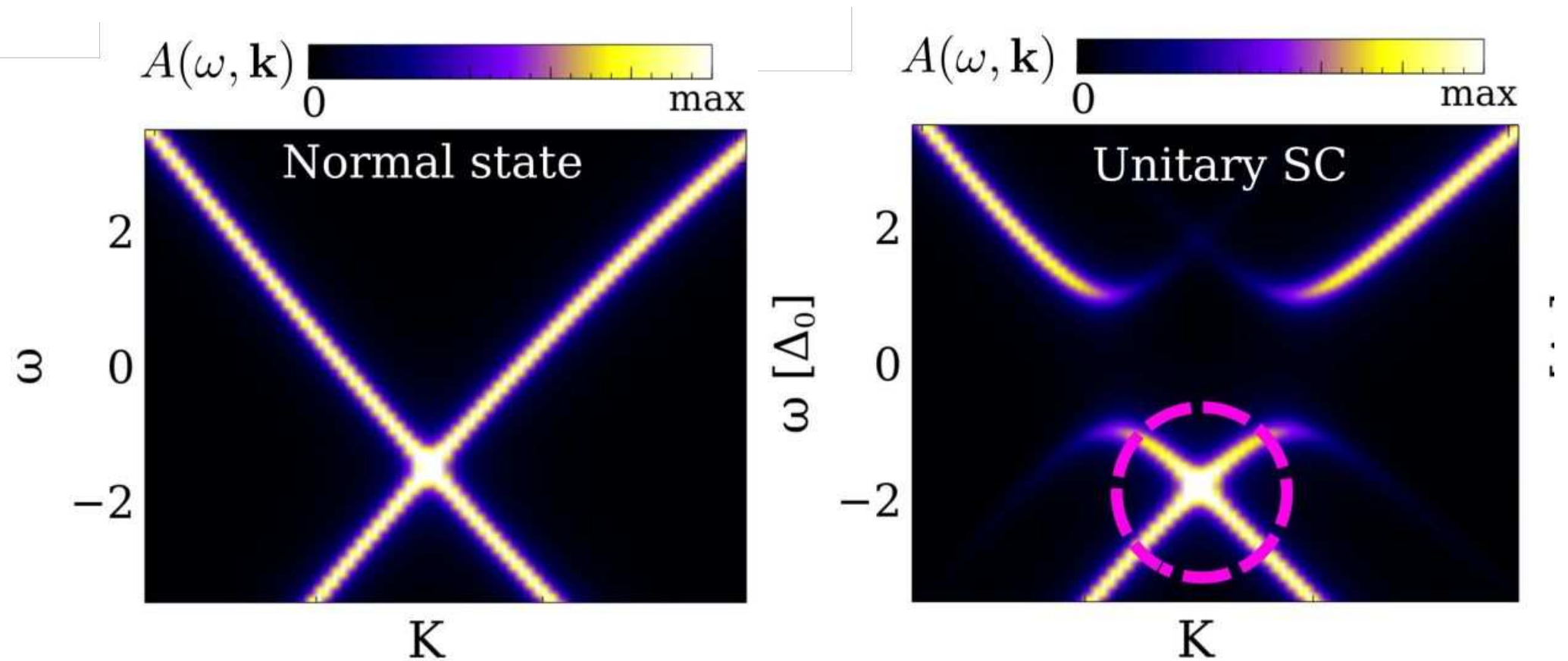
Dirac gap generation is a hallmark of non-unitary superconductivity

Dirac crossing in the superconducting state



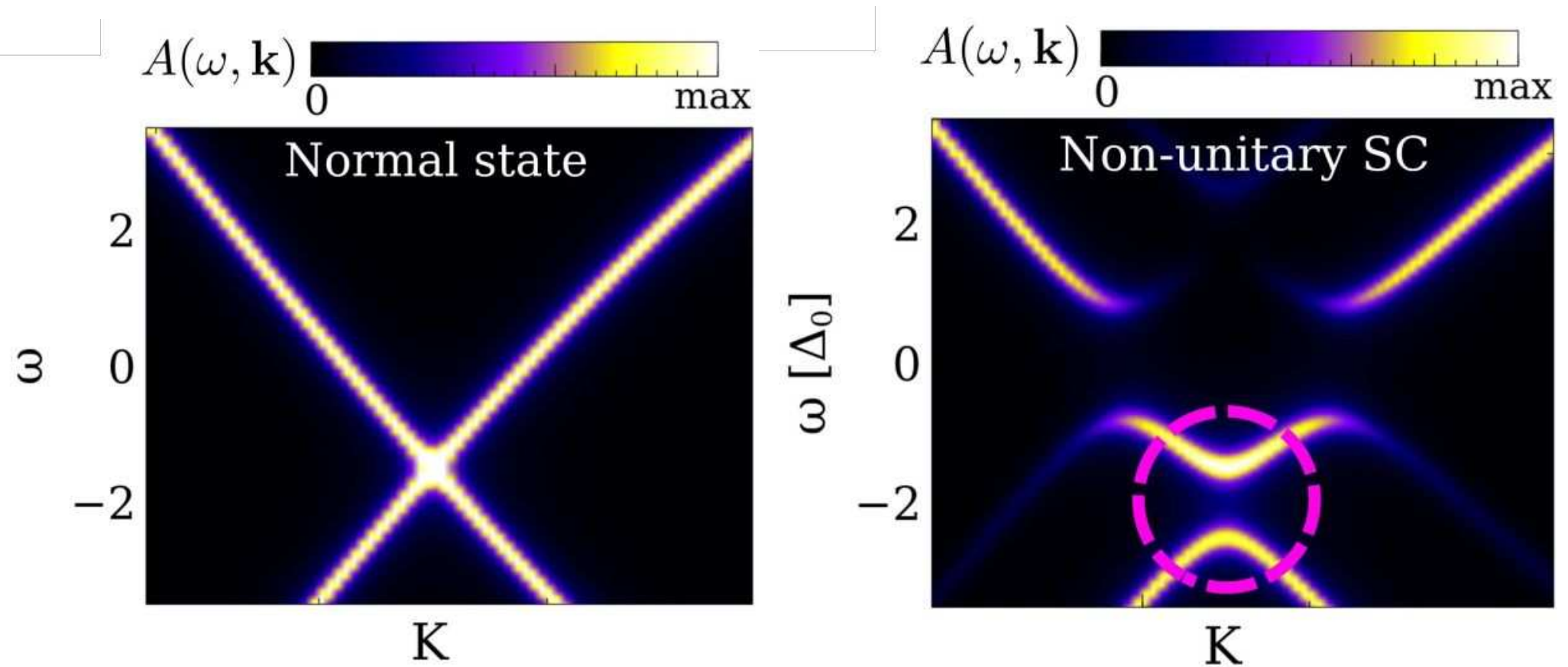
Dirac gap generation is a hallmark of non-unitary superconductivity

ARPES of a unitary SC



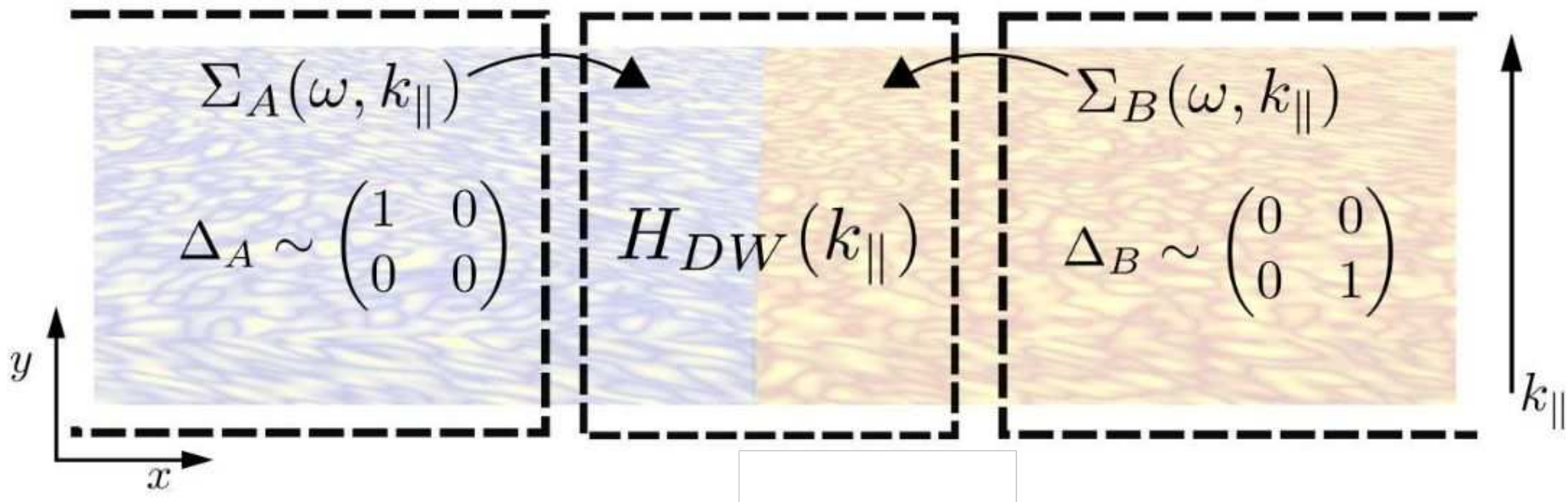
No gap generated at the Dirac point for a unitary state

ARPES of a non-unitary SC



Finite gap generated at the Dirac point for a non-unitary state

Superconducting orbital domain wall

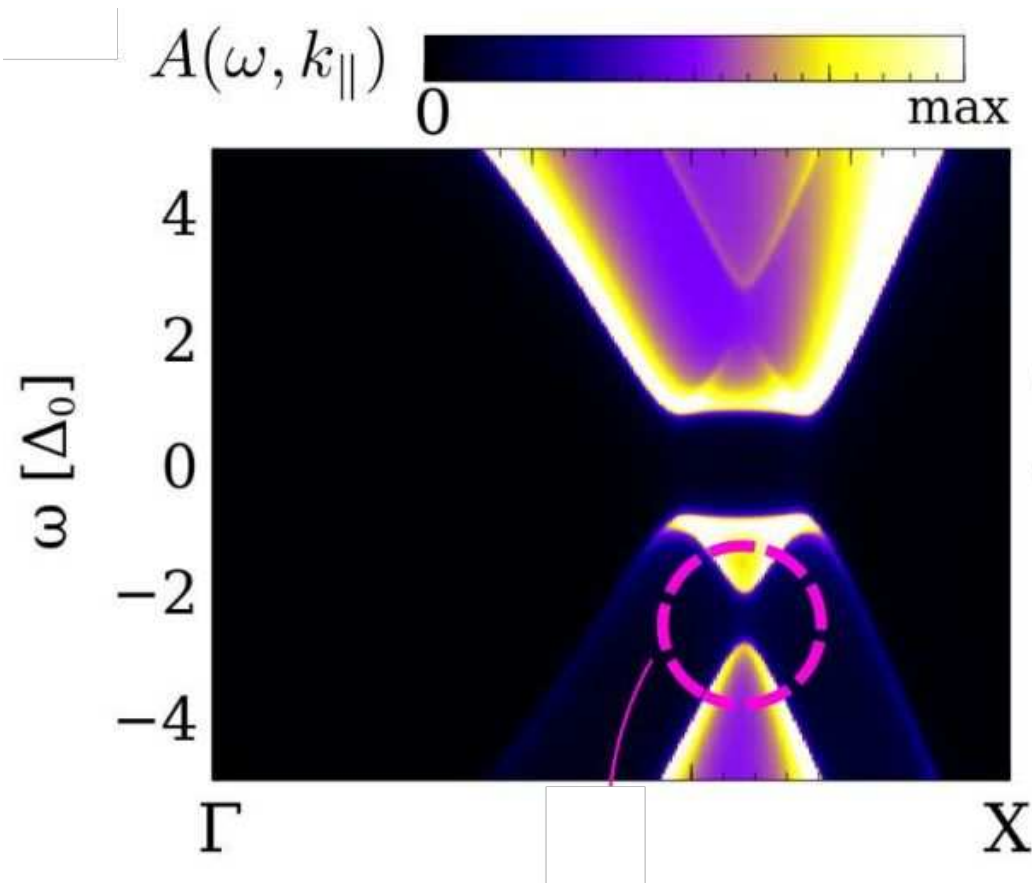


Is a superconducting domain wall visible in a spectral function?

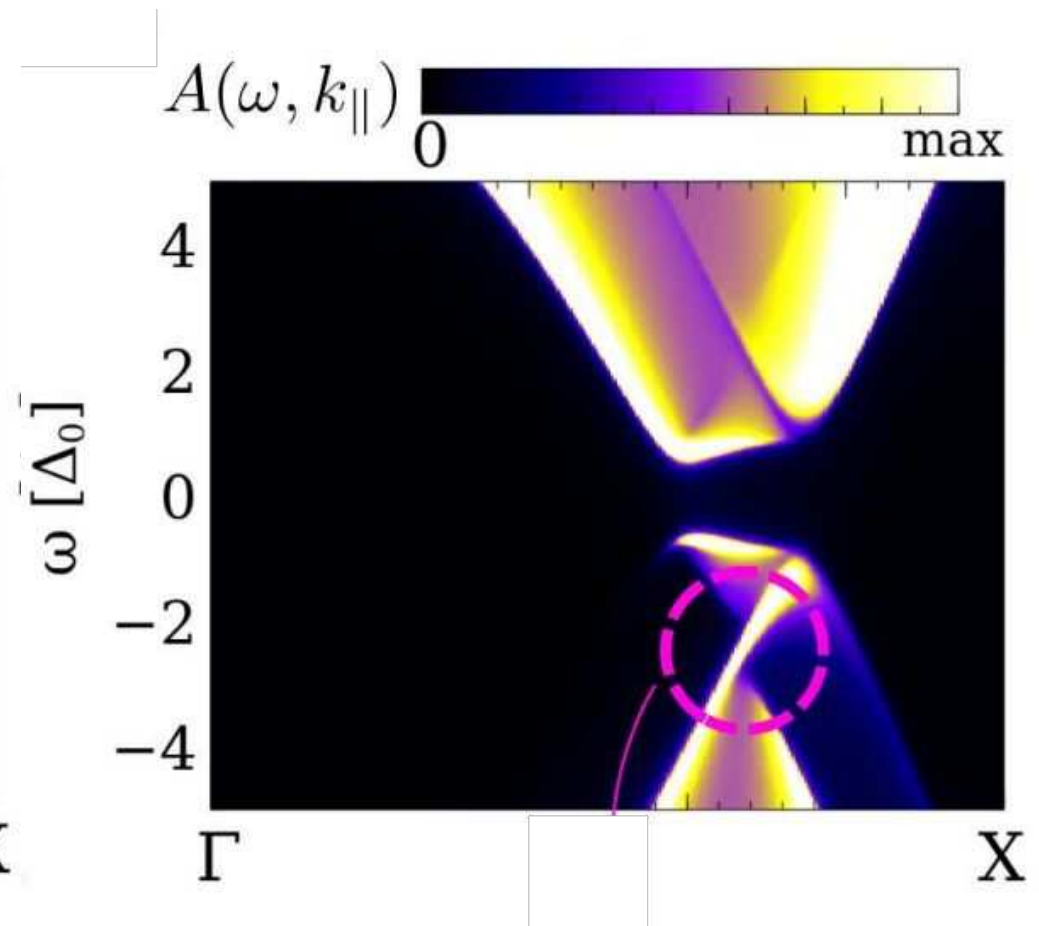
$$A(\omega, \mathbf{k}) = -\frac{1}{\pi} \text{Im}(\text{Tr}\{P_e[\omega - H(\mathbf{k}) + i0^+]^{-1}\})$$

ARPES of the SC with and without SC domain wall

Without domain wall



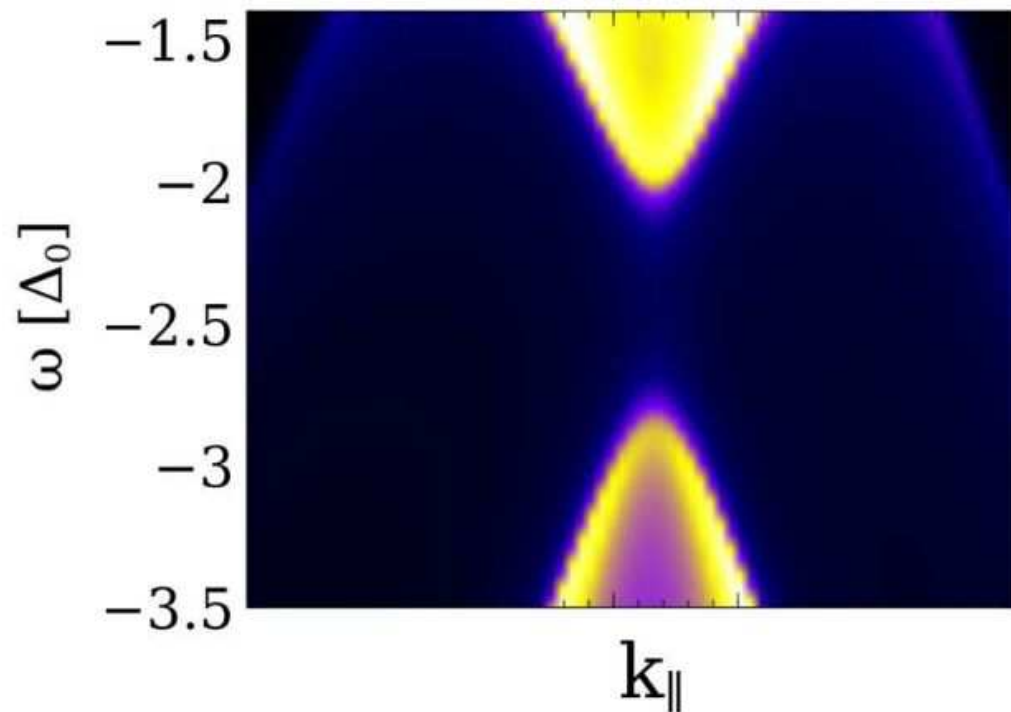
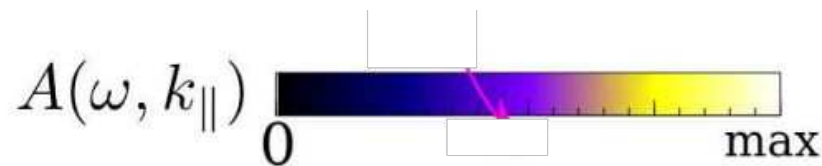
With domain wall



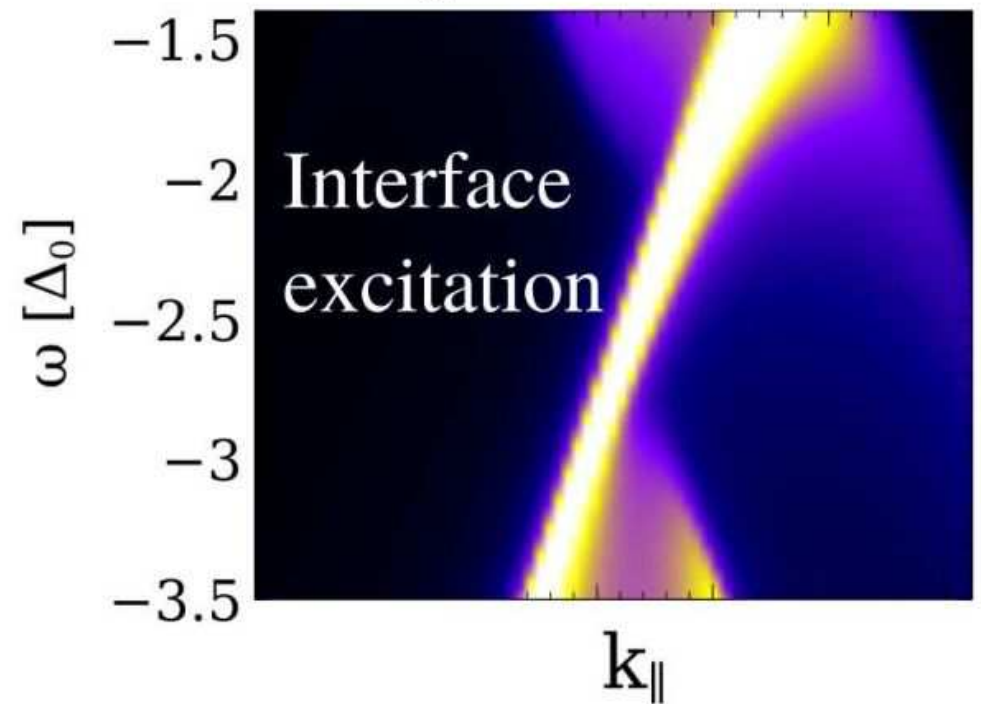
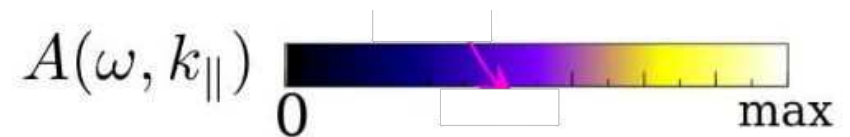
Superconducting domains generate states inside the Dirac gap

ARPES of the SC with and without SC domain wall

Without domain wall



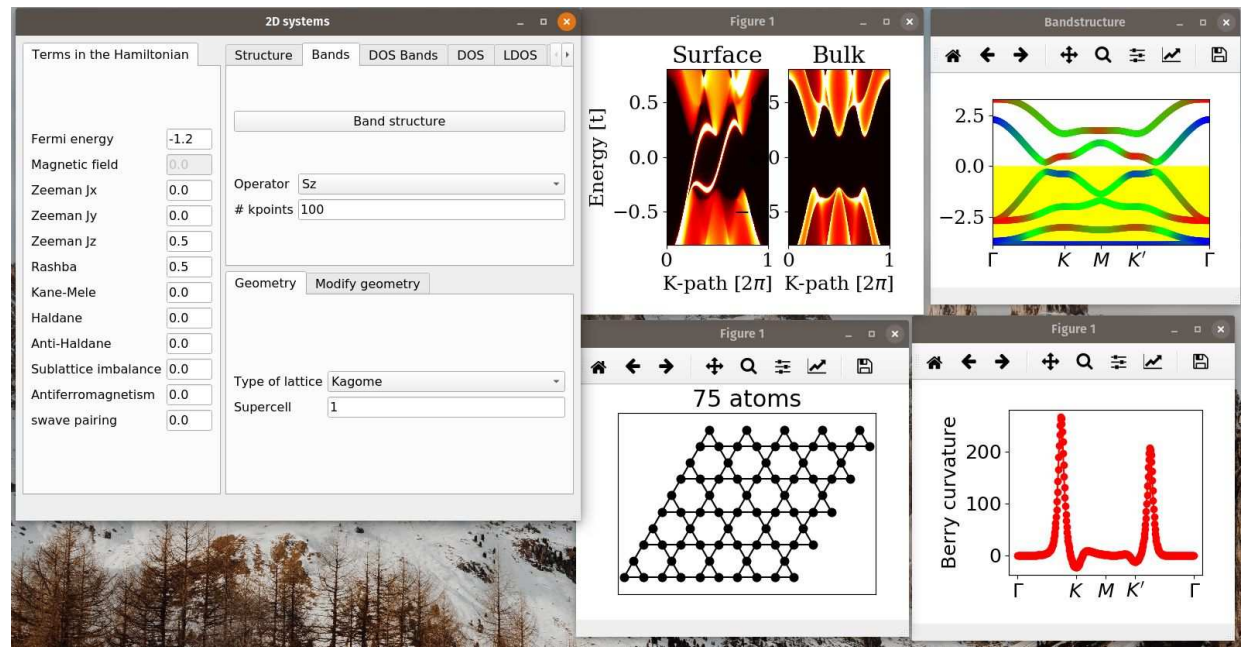
With domain wall



Superconducting domains generate states inside the Dirac gap

Computing electronic
properties

A user interface to compute electronic properties



Quantum Honeycomp: open source interactive interface for tight binding modeling

<https://github.com/joselado/quantum-honeycomp>

Thu Aug 16, 3:00 PM

Activities

```
jose@pop-os:~$ cd /home/jose/.local/share/quantum-honeycomp
jose@pop-os:~$ python3 path.py
/home/jose/.local/share/quantum-honeycomp/path.py
/home/jose/.local/share/quantum-honeycomp/path.py
```

Thu Aug 16, 4:55 PM

Activities

Thu Aug 16, 3:52 PM

Quantum Honeycomp: system selection

Select the system you want to compute

Linux and MacOSX compatible ☒ Only Linux compatible

Islands	Huge islands
Ribbons	Hybrid ribbons
Sheets	Twisted bilayer
Films	Hybrid films
Hofstadter butterflies	3D crystals

Update Quantum Honeycomp

Version 0.14.7

quantum-honeycomp-pyqt

Thu Aug 16, 4:55 PM

```
jose@pop-os:~$ cd /home/jose/.local/share/quantum-honeycomp
jose@pop-os:~$ python3 path.py
```

Quantum Honeycomp: system selection

Select the system you want to compute

Linux and MacOSX compatible ☒ Only Linux compatible

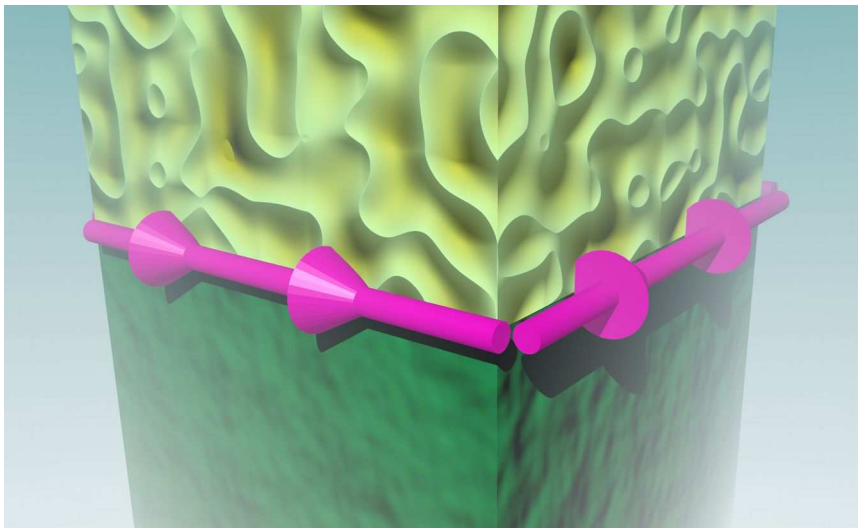
Islands	Huge islands
Ribbons	Hybrid ribbons
Sheets	Twisted bilayer
Films	Hybrid films
Hofstadter butterflies	3D crystals

Update Quantum Honeycomp

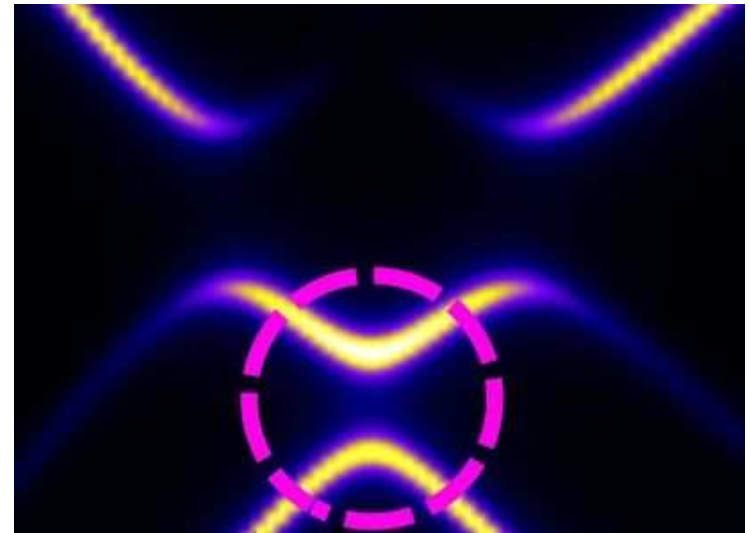
Version 0.14.7

Take home

- **Antiferromagnetically gaped Dirac points allow to generate topological superconductors**
- **ARPES can detect non-unitary multiorbital superconductivity in doped Dirac materials**



Phys. Rev. Lett. 121, 037002 (2018)



Phys. Rev. Research 1, 033107 (2019)