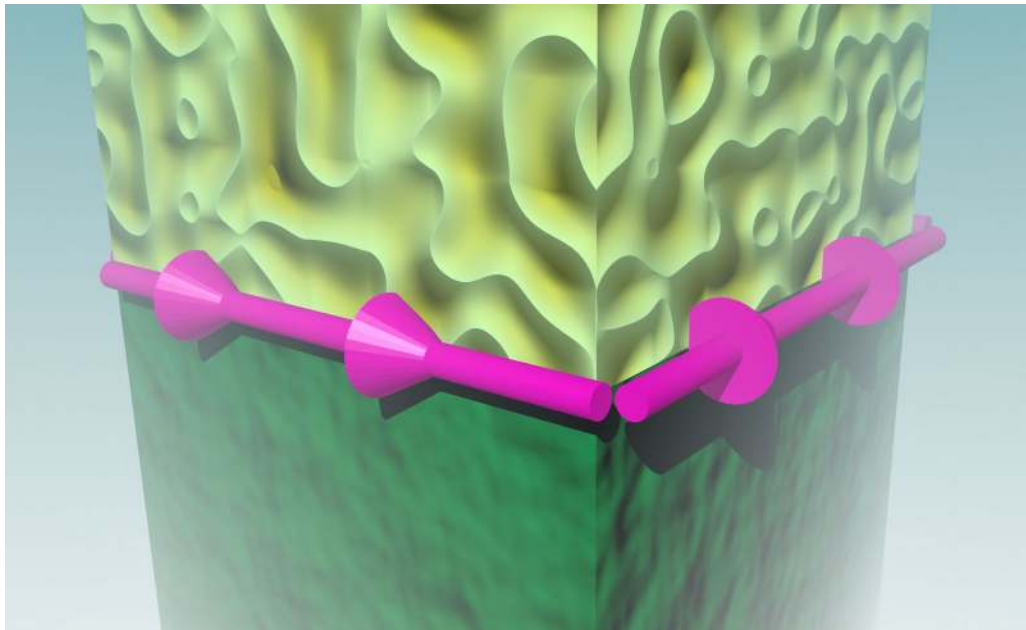


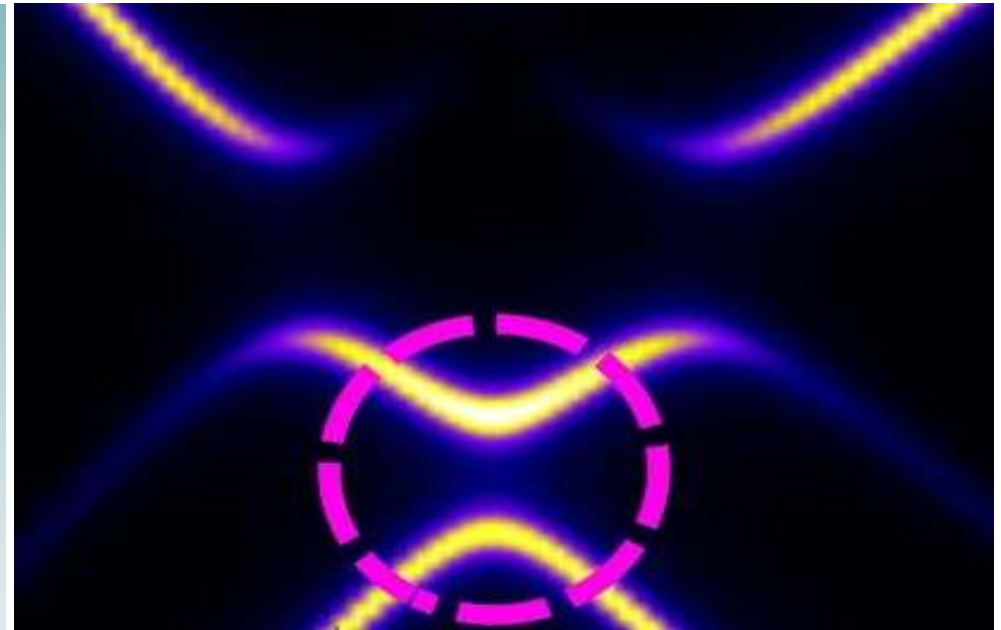
Engineering and detecting unconventional superconductivity with disguised Dirac points

Jose Lado and Manfred Sigrist

Aalto University, Finland & ETH Zurich, Switzerland



Phys. Rev. Lett. 121, 037002 (2018)

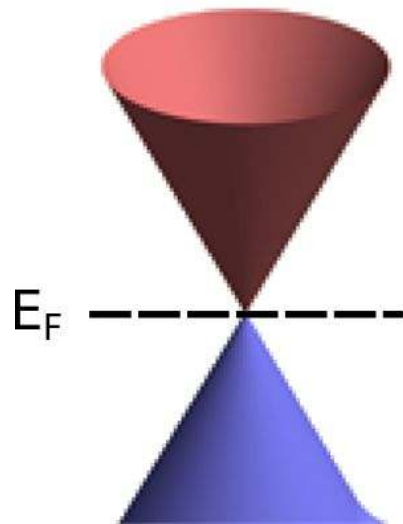


arXiv:1908.04820 (2019)

2019 Zhejiang Workshop on Correlated Matter, September 21st 2019

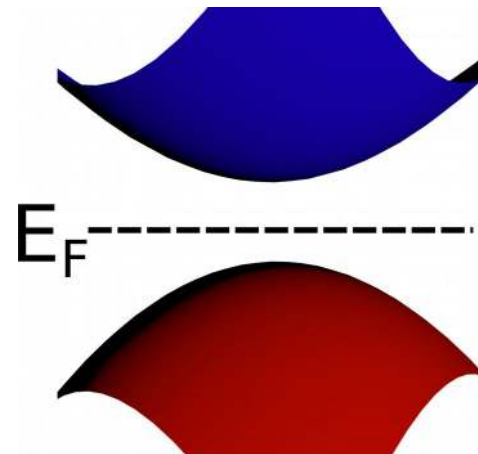
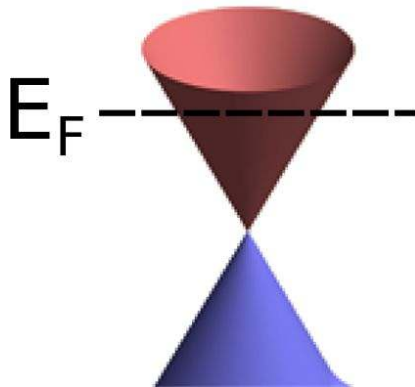
On Dirac crossings

A platform to engineer exotic physics in solid state materials



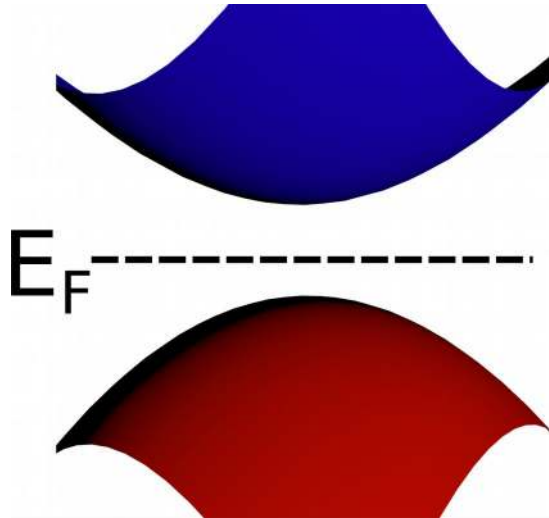
WINDING NUMBER
TOPOLOGY **NODAL LINES**
TWISTED BILAYER GRAPHENE
DIRAC
BERRY FLUX
CHERN NUMBER
BERRY PHASE
Weyl POINTS
NODAL CHAINS
SPINOR
ANOMALOUS HALL

But what if the Dirac points are not observable at the chemical potential?



Exploiting disguised Dirac points

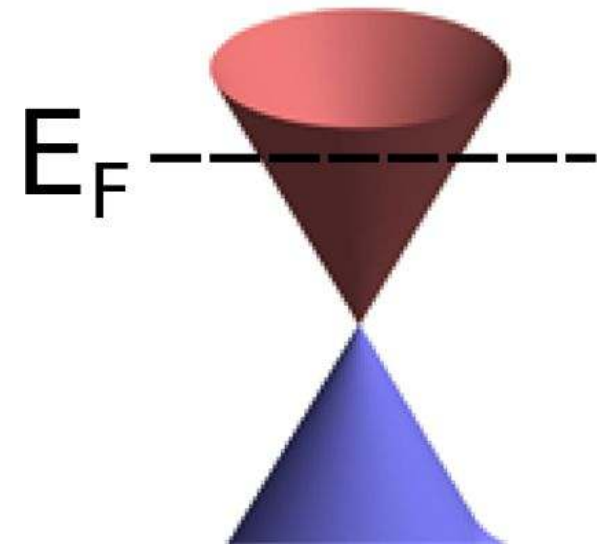
Platform for an artificial
topological superconductor



Gaped out by antiferromagnetism

Phys. Rev. Lett. 121, 037002 (2018)

Platform to detect
non-unitary superconductivity



Away from the Fermi energy

arXiv:1908.04820 (2019)



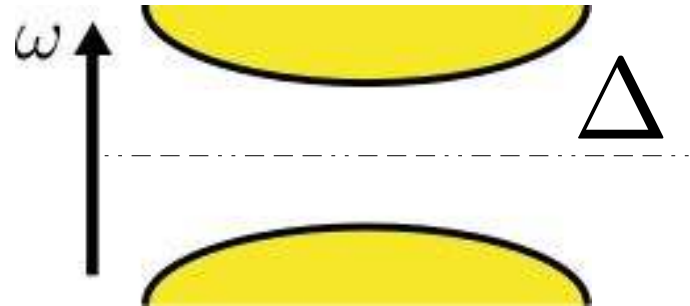
Manfred
Sigrist

**Creating a topological
superconductor with
antiferromagnetically
gaped Dirac points**

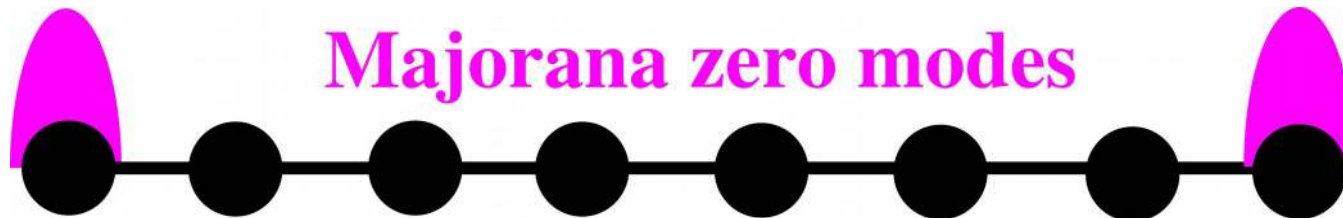
Topological superconductors

Topological superconductors with broken time reversal symmetry

Gapped bulk excitations

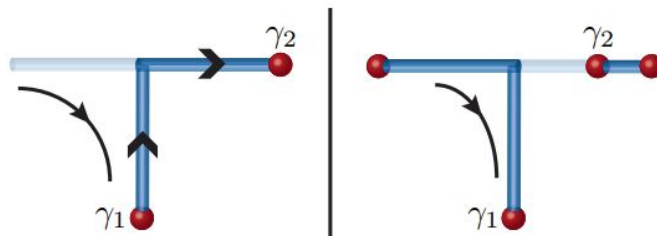


Gapless surface modes, single Majorana mode per edge



A. Y. Kitaev.
Physics-Uspekhi,
44:131, 2001

Anyonic statistics, suitable for quantum computing



Nature Physics 7, 412-417 (2011)

Our objective

Build a 2D topological superconductor with

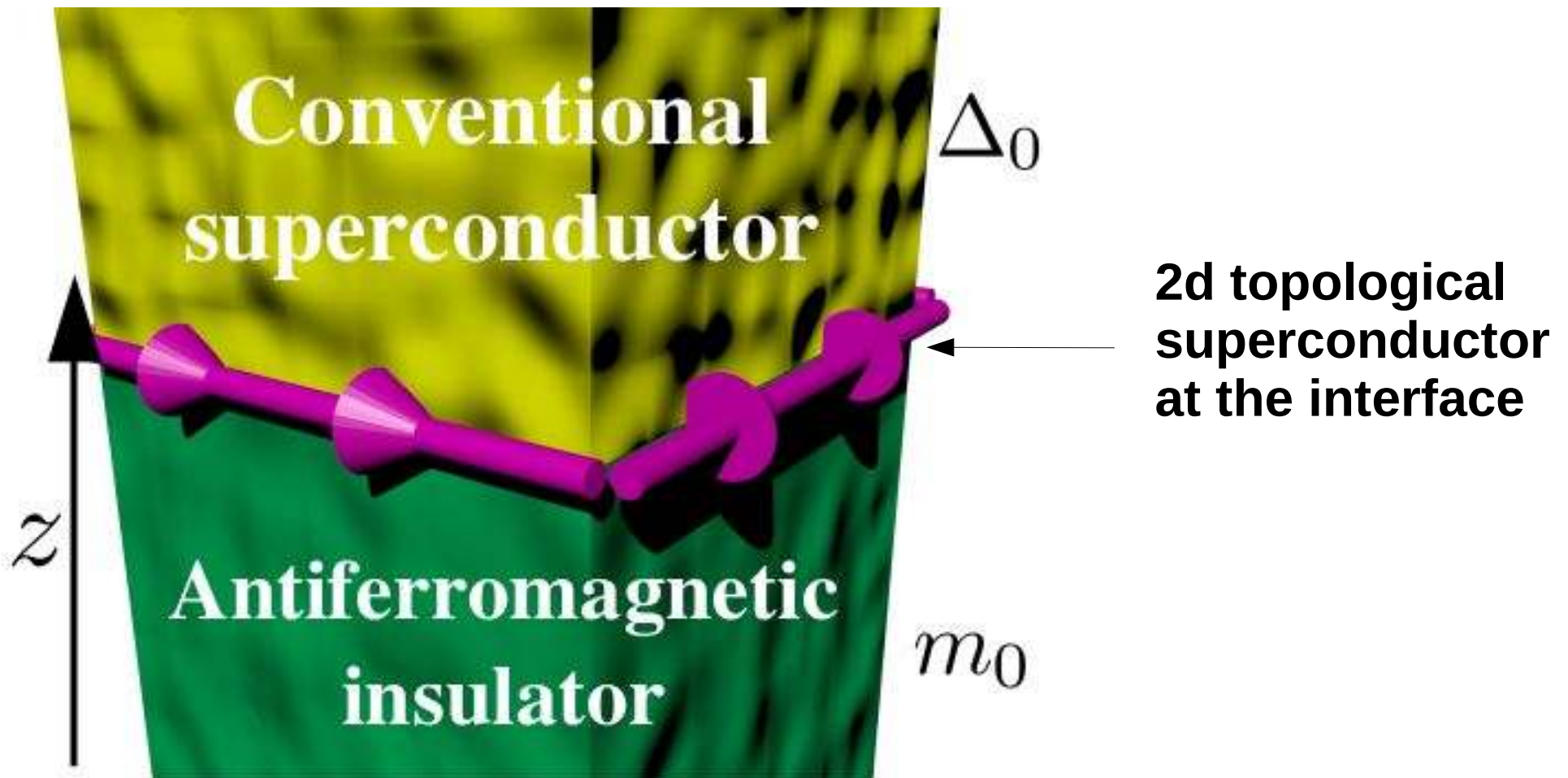
- A conventional (s-wave) 3D superconductor
- A 3D antiferromagnetic insulator

The prize



Bringing a new solid state platform to realize
artificial topological superconductors

The system for TS in an AF insulator



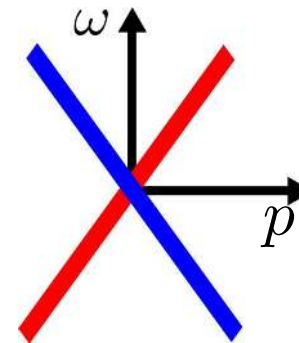
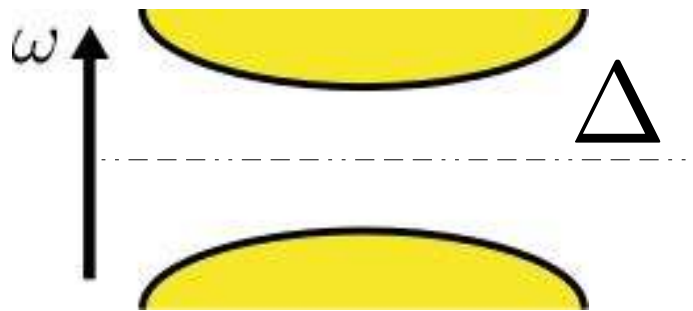
How to do your own topological superconductor



Ingredients

- s-wave pairing
- Helical states

Science 346.6209 (2014)
Science 354.6319 (2016)
Phys. Rev. Lett. 100, 096407 (2008)



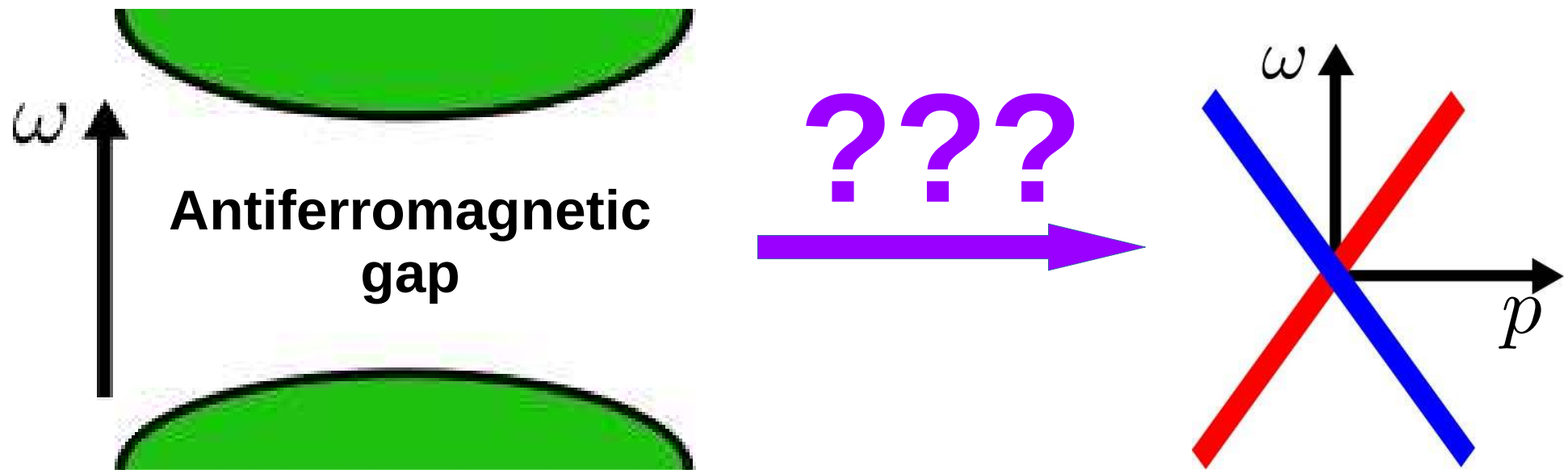
Objective: realize a spin-less superconductor

$$H = \sum_n t c_{n+1}^\dagger c_n + \Delta c_n c_{n+1} + c.c.$$

A. Y. Kitaev. *Physics-Uspekhi*, 44:131, 2001

The initial problem

How can we get a topological phase starting from a trivial insulator?



We need to create a “spinless” gapless state out of an insulator

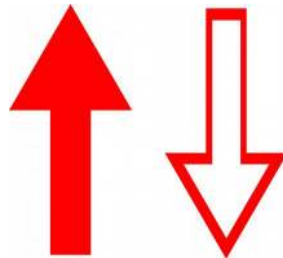
Solitonic modes between Dirac AF and SC

Total Hamiltonian

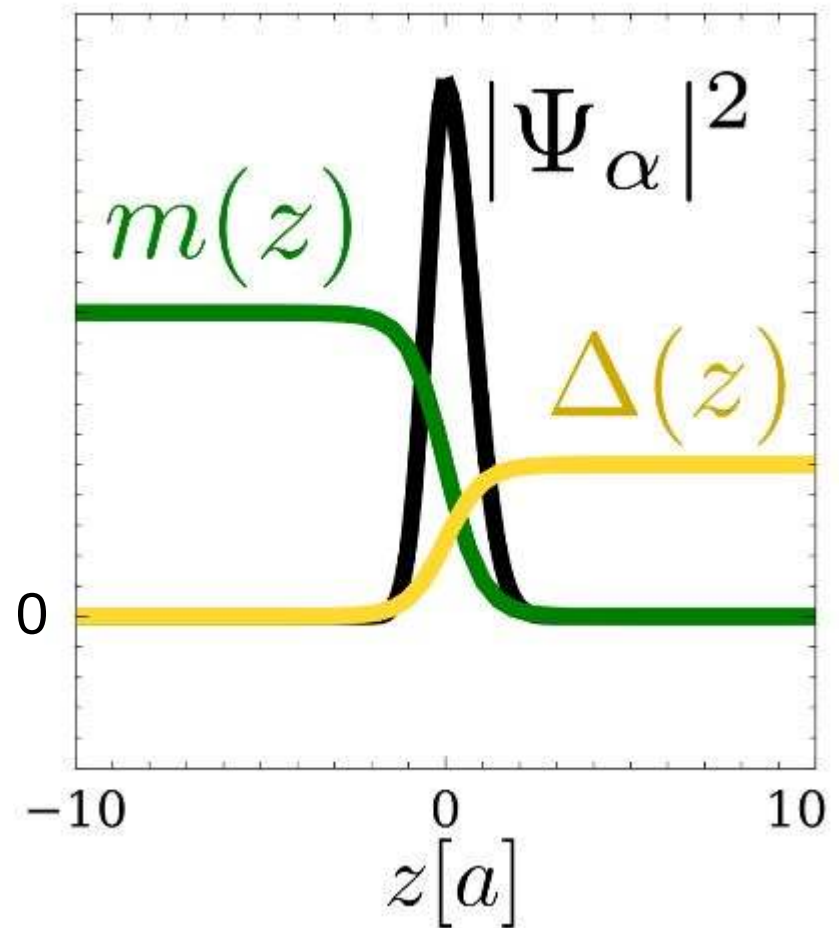
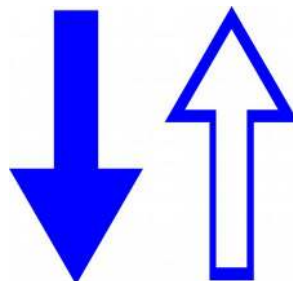
$$H = H_{kin} + H_{AF} + H_{SC}$$

There will be two zero solutions

Sector #1
Up electron, down hole

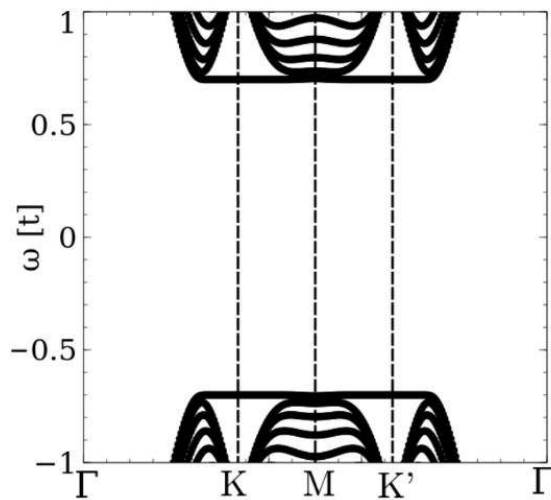


Sector #2
Down electron, up hole

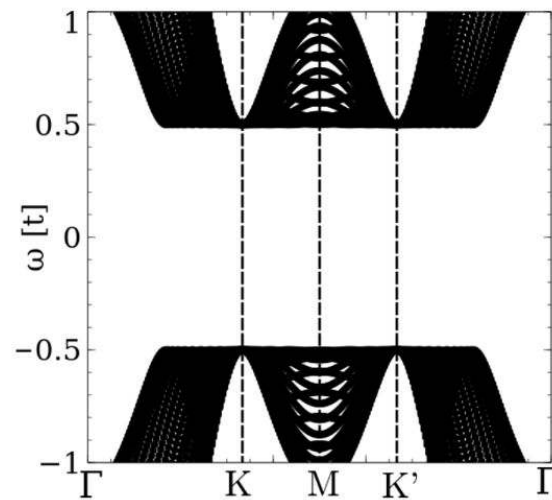


Emergence of interfacial modes, no spin-orbit coupling

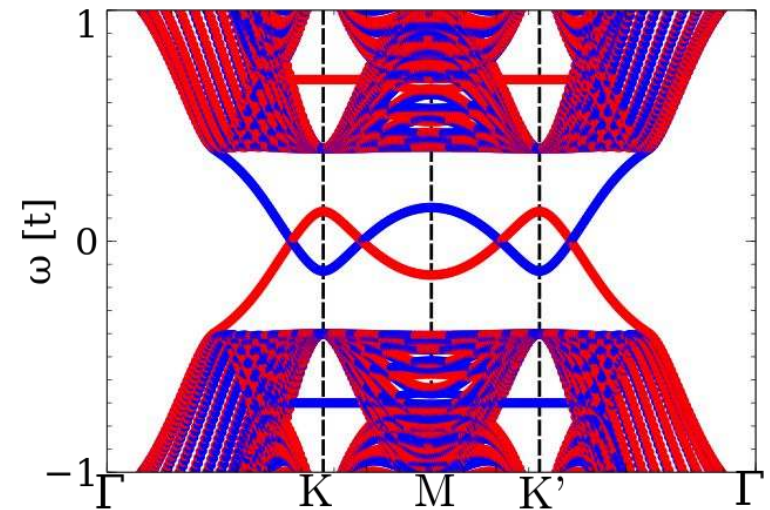
Antiferromagnet



Superconductor



AF/SC heterostructure

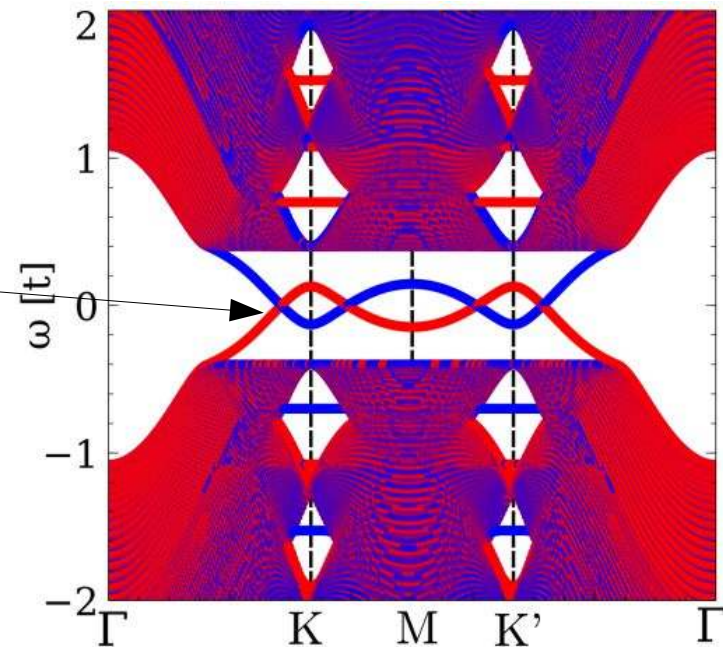
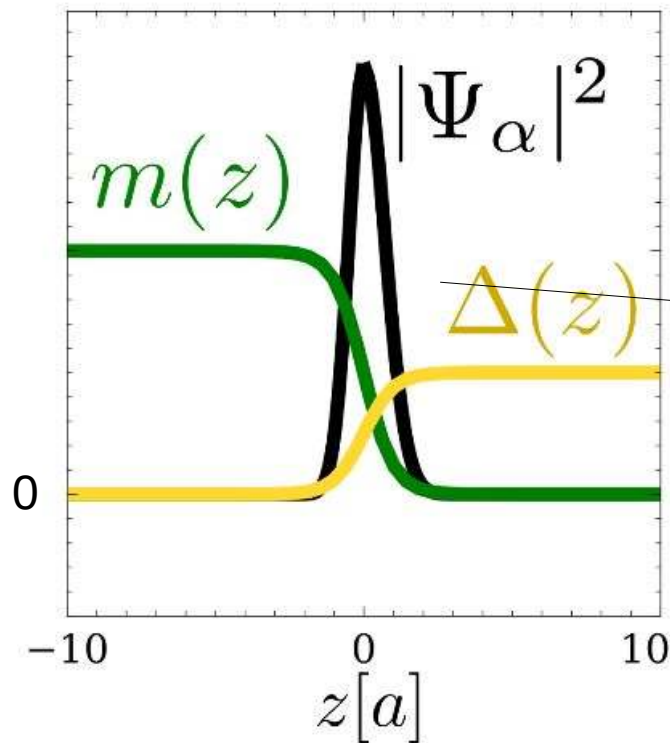
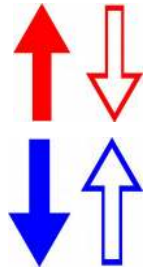


Interface states between AF and SC

Superconducting solitonic excitations

$$\Psi_1(z) = g(z)[c_{A,\vec{k}_{\parallel},\uparrow} - ic_{B,\vec{k}_{\parallel},\uparrow} - c_{A,-\vec{k}_{\parallel},\downarrow}^\dagger - ic_{B,-\vec{k}_{\parallel},\downarrow}^\dagger]$$

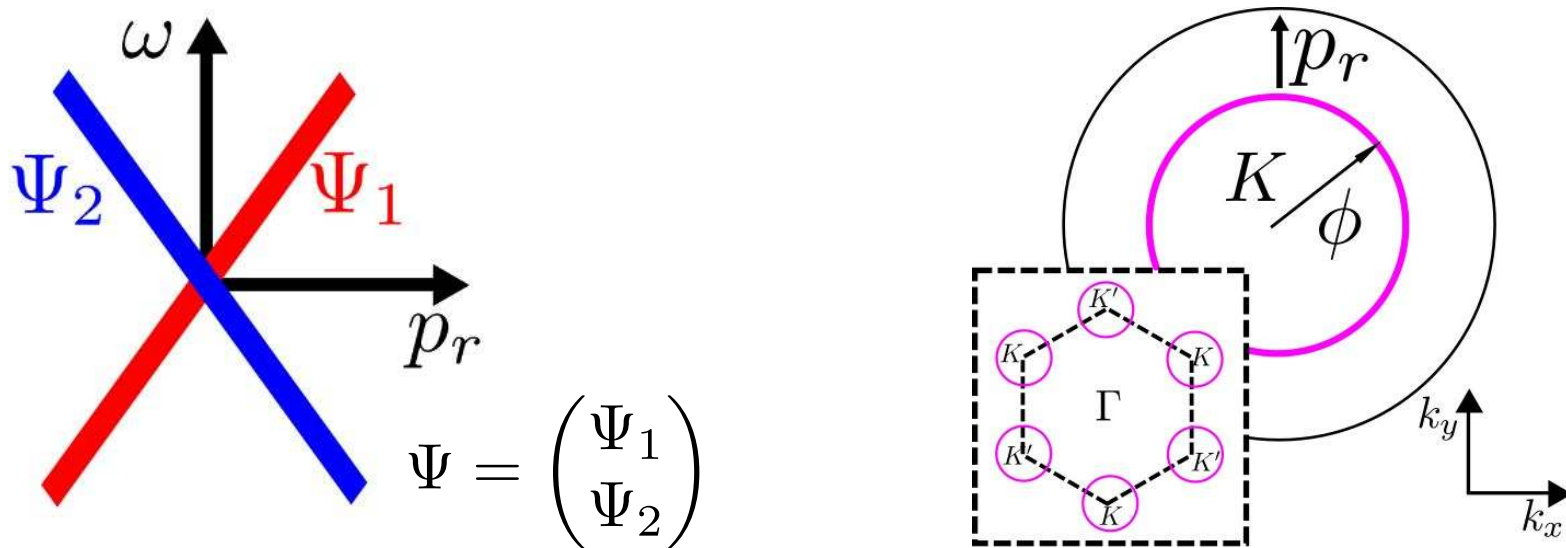
$$\Psi_2(z) = g(z)[c_{A,\vec{k}_{\parallel},\downarrow} + ic_{B,\vec{k}_{\parallel},\downarrow} - c_{A,-\vec{k}_{\parallel},\uparrow}^\dagger + ic_{B,-\vec{k}_{\parallel},\uparrow}^\dagger]$$



Effective model without spin-orbit coupling

Effective Hamiltonian without spin-orbit coupling

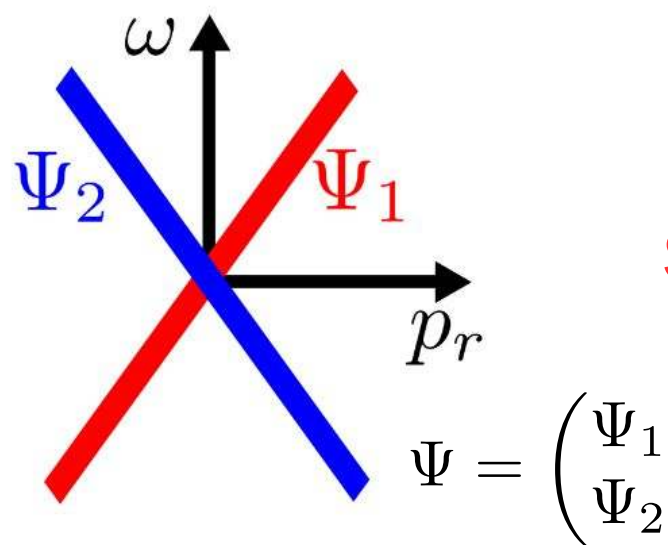
$$\mathcal{H}(p_r, \phi) = \Psi^\dagger \begin{pmatrix} v_r p_r & 0 \\ 0 & -v_r p_r \end{pmatrix} \Psi$$



Every point of the Dirac line creates a zero mode

Effective model with spin-orbit coupling

Effective Hamiltonian with spin-orbit coupling

$$\mathcal{H}(p_r, \phi) = \Psi^\dagger \begin{pmatrix} v_r p_r & -i\lambda e^{i\phi} \\ i\lambda e^{-i\phi} & -v_r p_r \end{pmatrix} \Psi$$


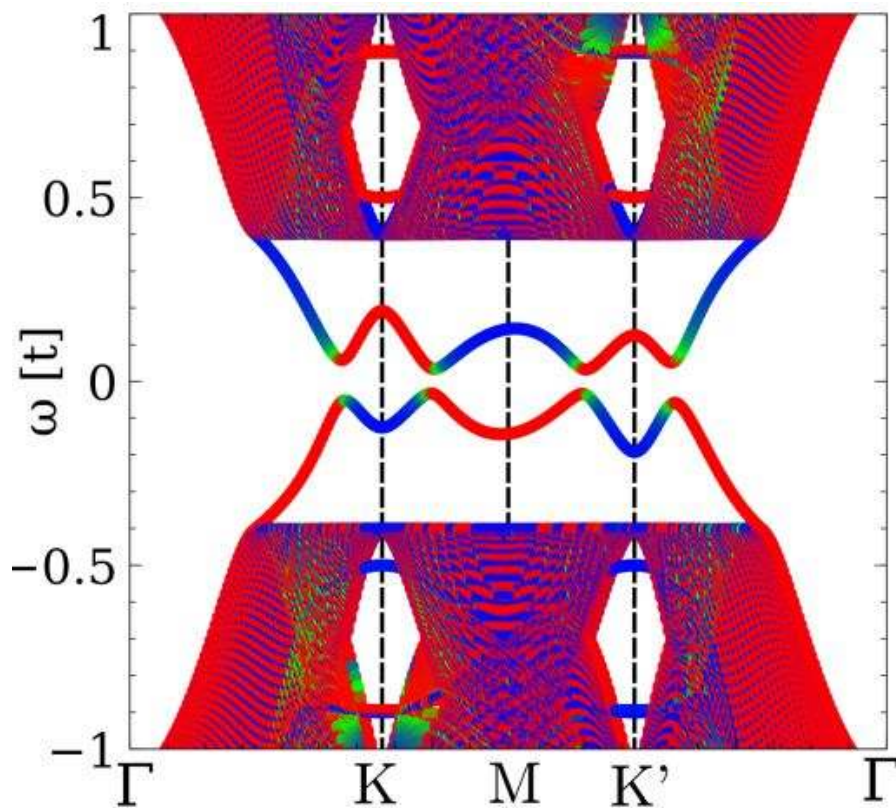
Spin-orbit coupling takes the form of a chiral pairing

$$\Psi = \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix}$$

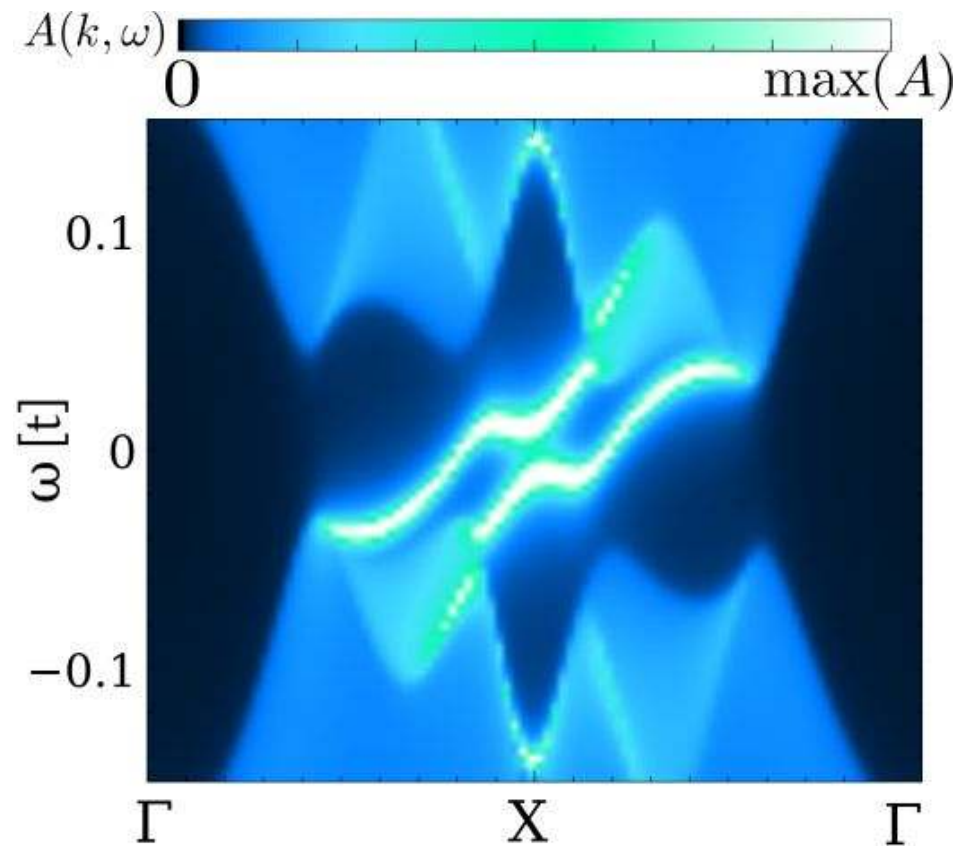
Chern number $\mathcal{C} = \pm 1$ per Dirac line

Adding spin-orbit coupling

Band structure

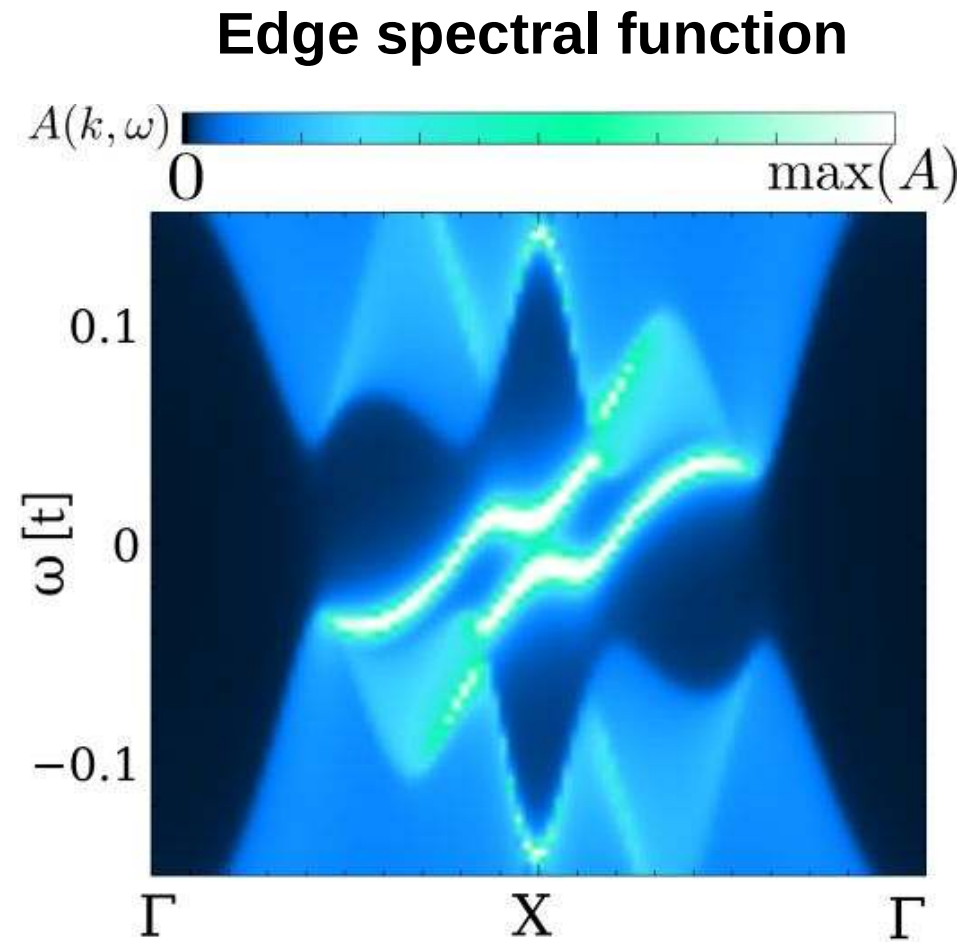
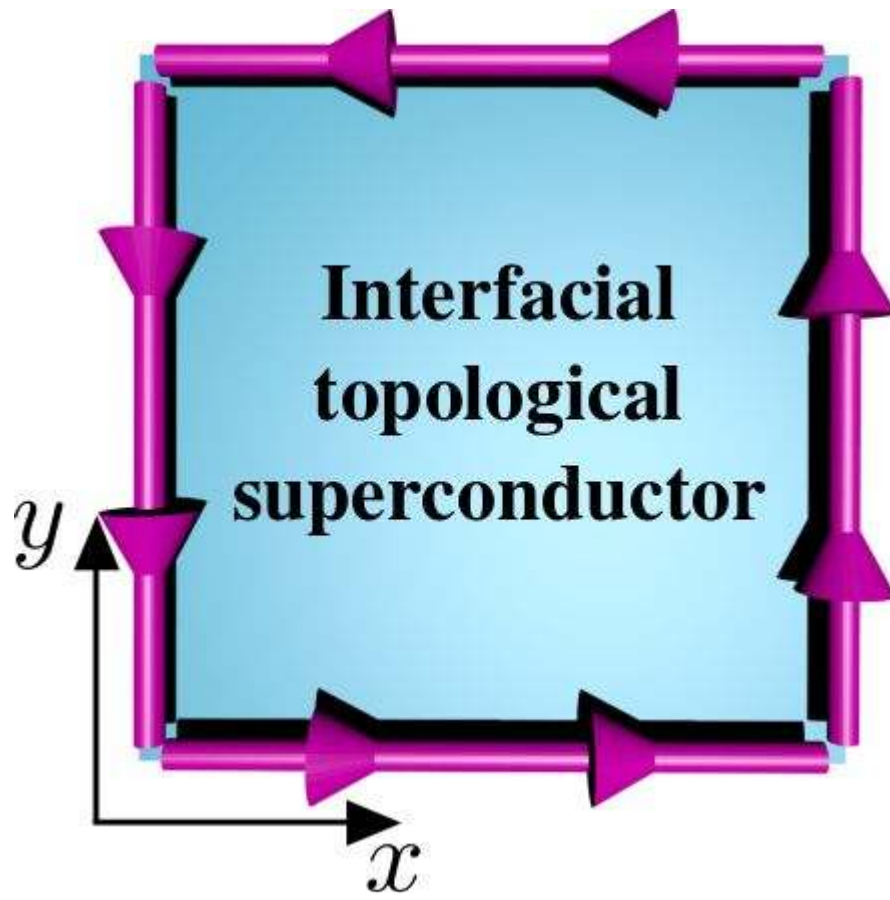


Edge spectral function



Two co-propagating edge modes

An interfacial topological superconductor

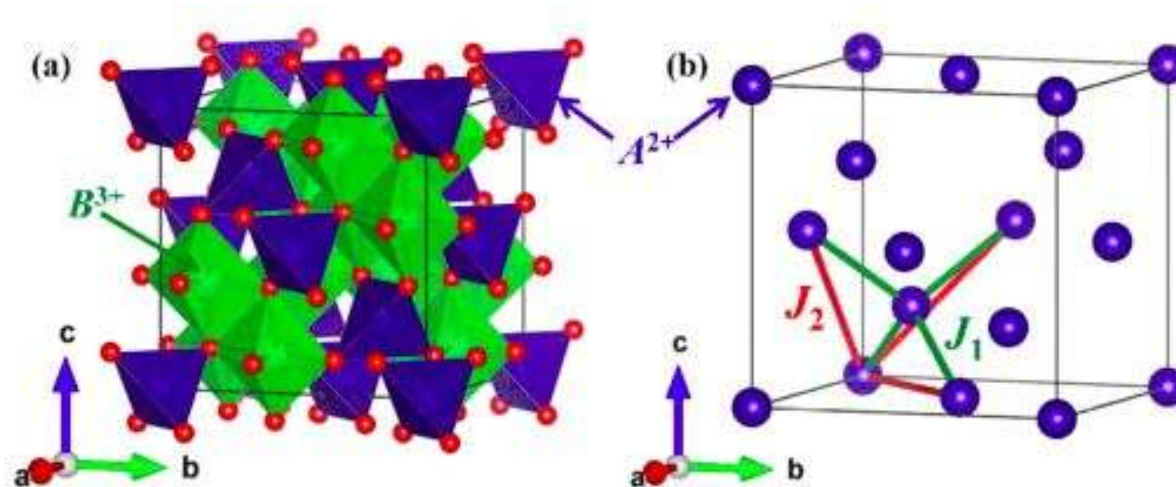


The interface realizes a topological superconductor

Possible materials, spinels

- Antiferromagnet forming a diamond lattice

Antiferromagnetic spinels

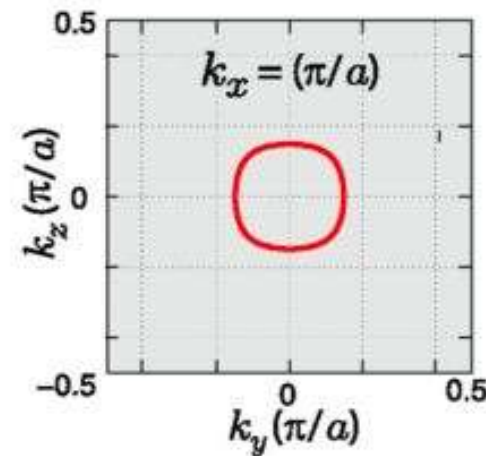
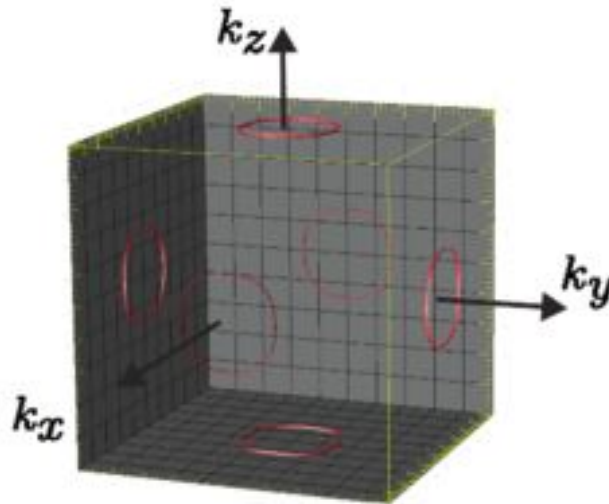


Co atoms form a diamond lattice

Possible materials, Dirac materials in the AF phase

Dirac lines in the absence of spin-orbit coupling and magnetism

Phys. Rev. Lett. 115, 036806 (2015)

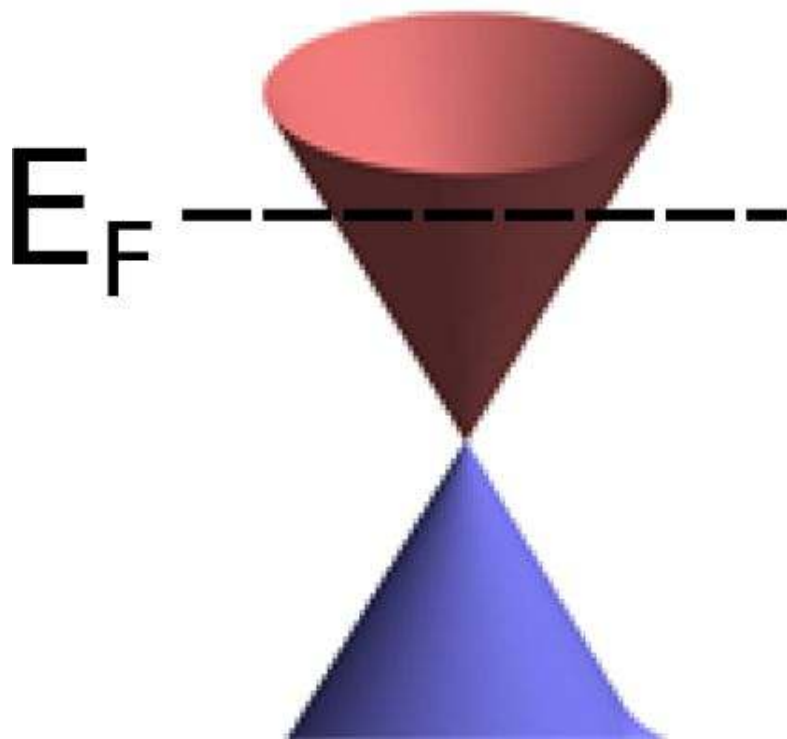


Antiferromagnets whose paramagnetic state hosts Dirac lines

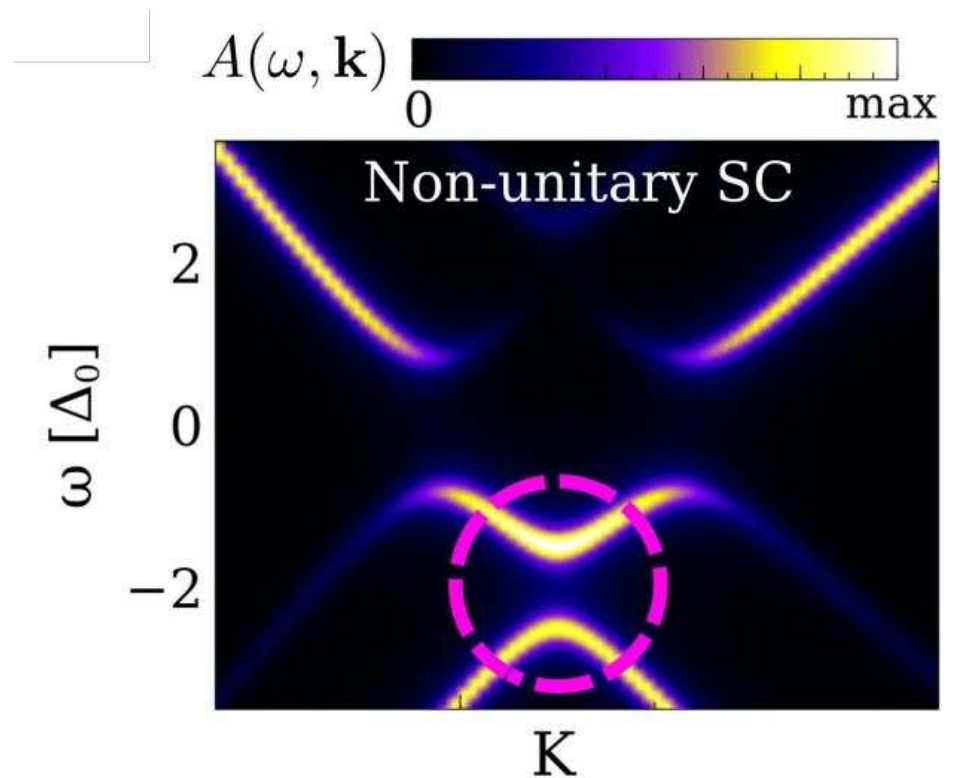
**Detecting non-unitary
superconductivity through
deep Dirac points**

Superconducting symmetry from ARPES experiments

Normal state band structure



ARPES in superconducting state



Can we detect non-unitary superconductivity from an ARPES experiment?

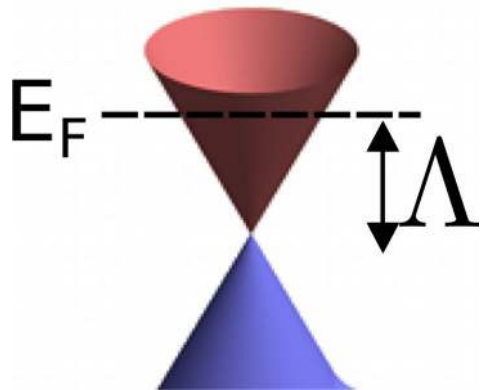
Gap opening by singlet multiorbital superconductivity

Two orbital band structure with a Dirac crossing, singlet pairing

Hamiltonian in normal state

$$\psi_{\mathbf{k}}^\dagger = (c_{\alpha,\mathbf{k}}^\dagger, c_{\beta,\mathbf{k}}^\dagger);$$

$$H_0^{DP}(\mathbf{k}) = -\Lambda\tau_0 + \tau_x k_x + \tau_y k_y$$



Hamiltonian in SC state

$$\Psi_{\mathbf{k}}^\dagger = (c_{\alpha_0,\uparrow,\mathbf{k}}^\dagger, c_{\beta_0,\uparrow,\mathbf{k}}^\dagger, c_{\alpha_0,\downarrow,-\mathbf{k}}, c_{\beta_0,\downarrow,-\mathbf{k}}).$$

$$H_{BdG}(\mathbf{k}) = \begin{pmatrix} H_0^{DP}(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta(\mathbf{k}) & -H_0^{DP}(-\mathbf{k}) \end{pmatrix}$$

Effective Hamiltonian after Schrieffer-Wolff transformation

$$H^{DP} = H_0^{DP} + \frac{\Delta\Delta^\dagger}{2\Lambda}$$

SC gap opening at the Dirac point if $\Upsilon = \text{Tr}[\tau_z \Delta\Delta^\dagger] \neq 0$

Conditions for a Dirac gap opening

$$\Upsilon = \text{Tr}(\tau_z \Delta \Delta^\dagger)$$

No gap opening $\Upsilon = 0$

Unitary states

Intraorbital

$$\Delta \sim \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Interorbital

$$\Delta \sim \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

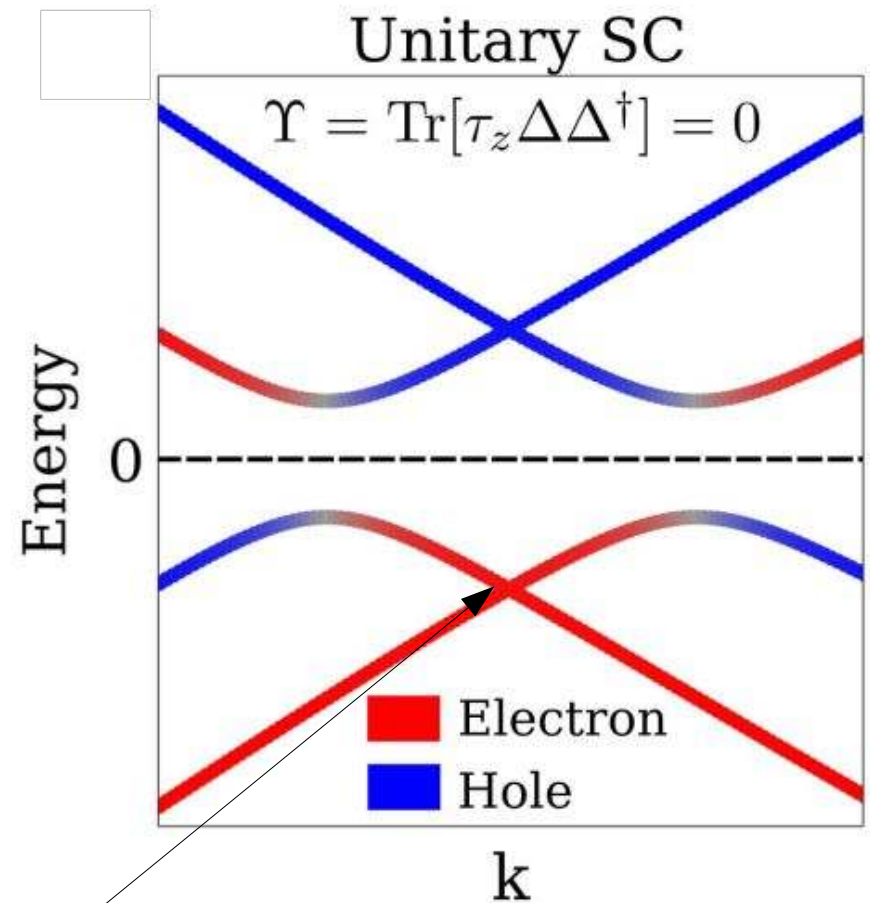
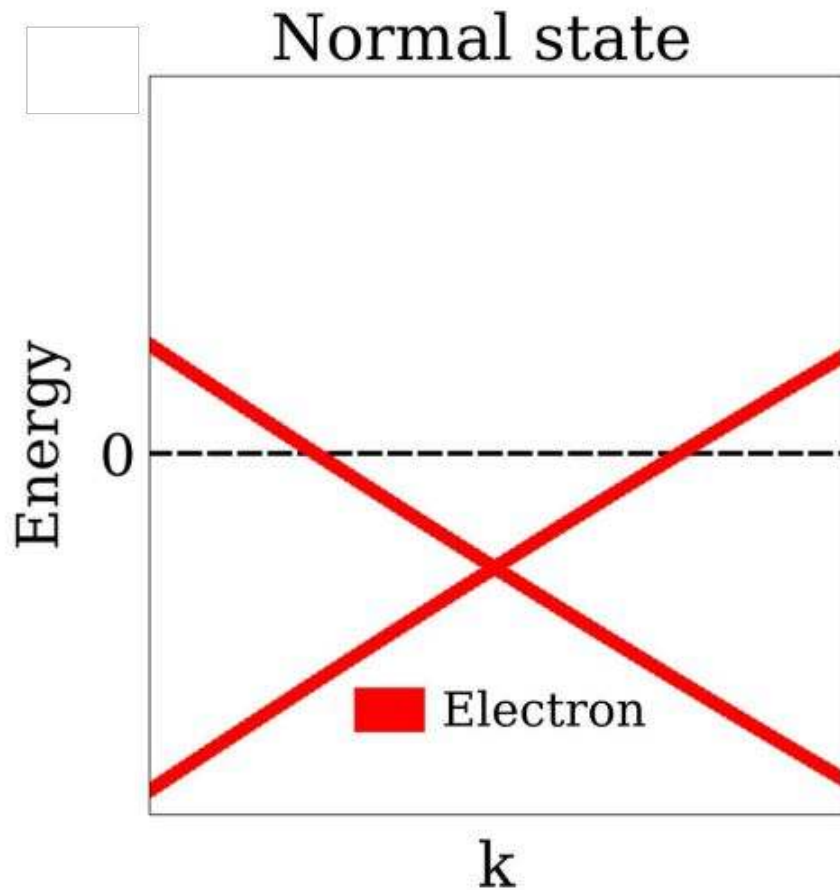
Gap opening $\Upsilon \neq 0$

Non-unitary states

Non-unitary intraorbital

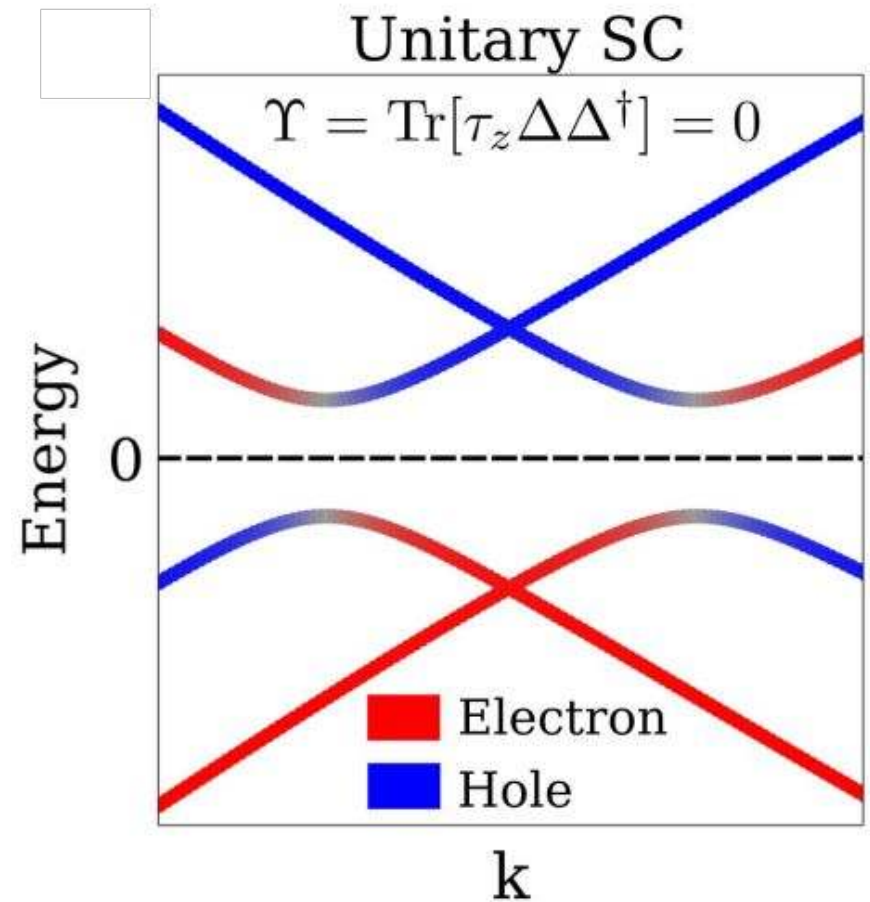
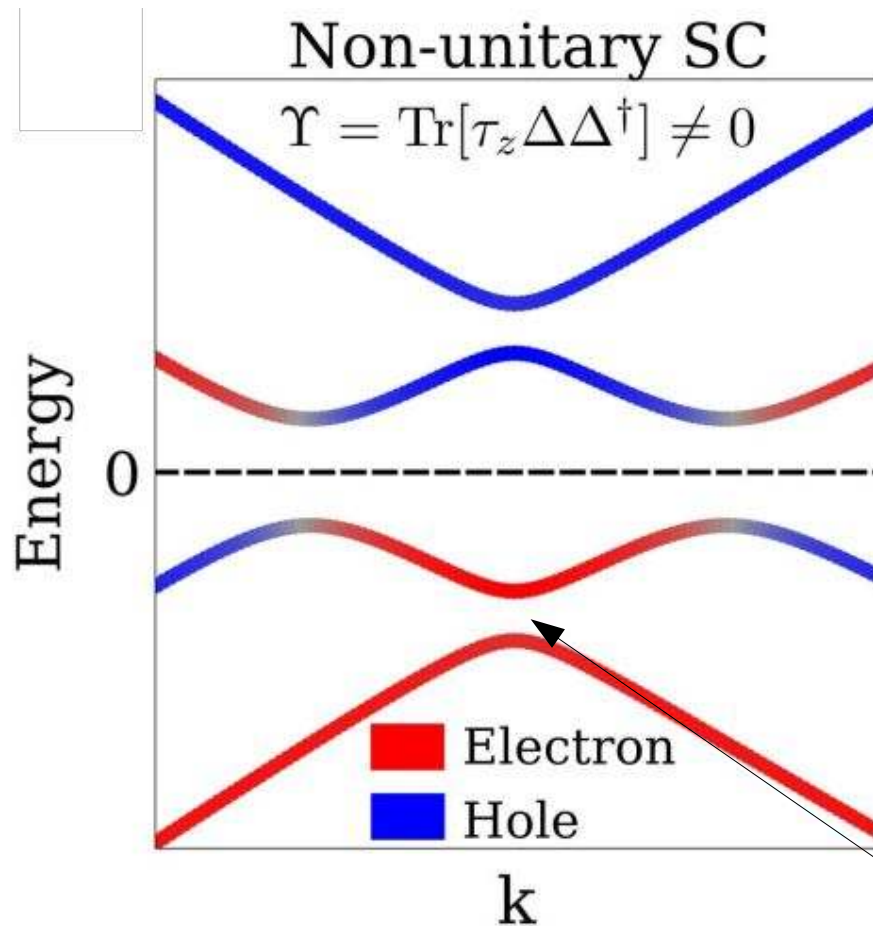
$$\Delta \sim \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Dirac crossing in the superconducting state



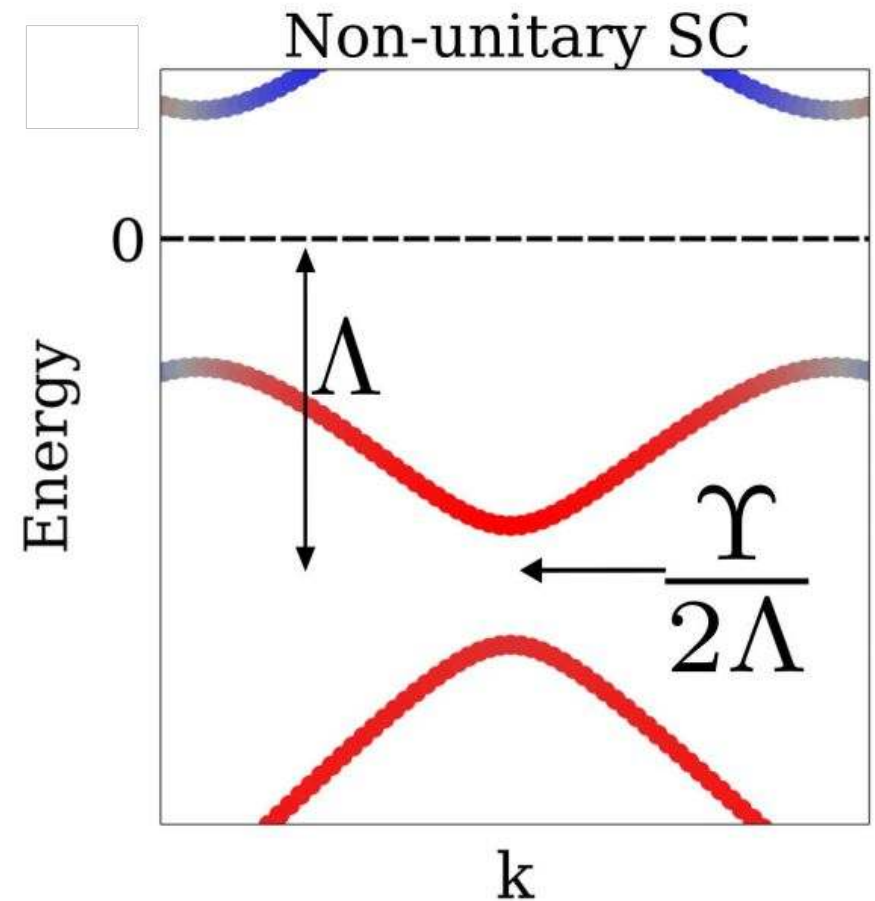
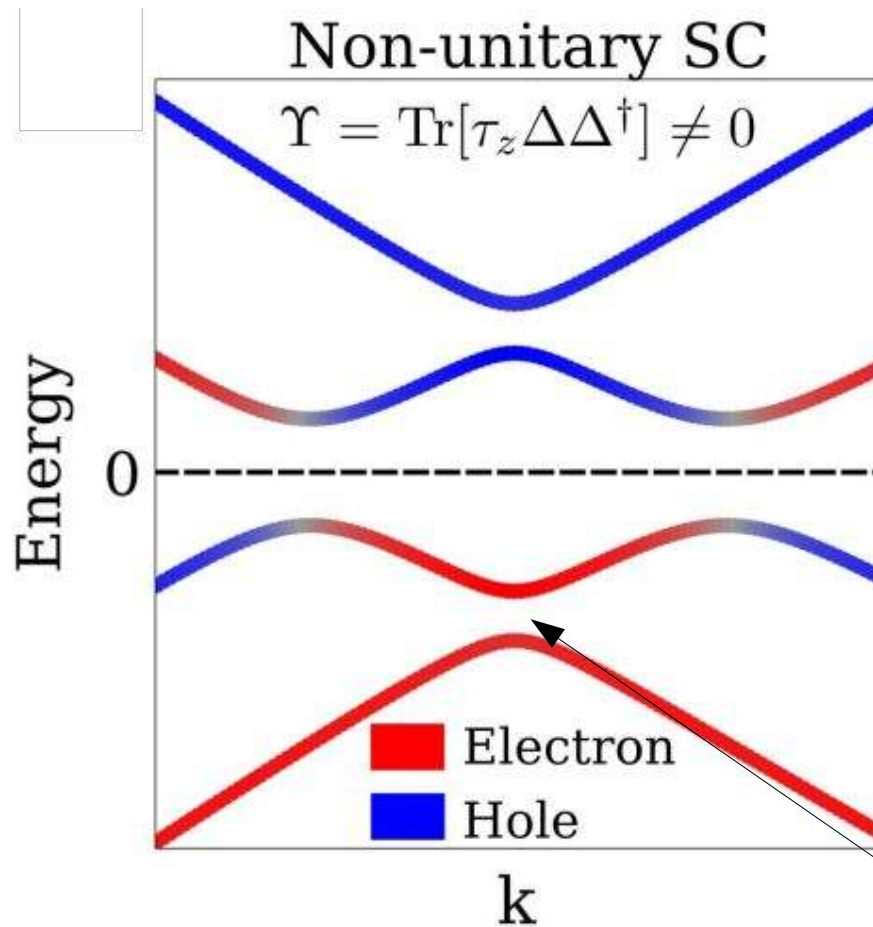
Conventional superconducting state: Dirac point remains closed

Dirac crossing in the superconducting state



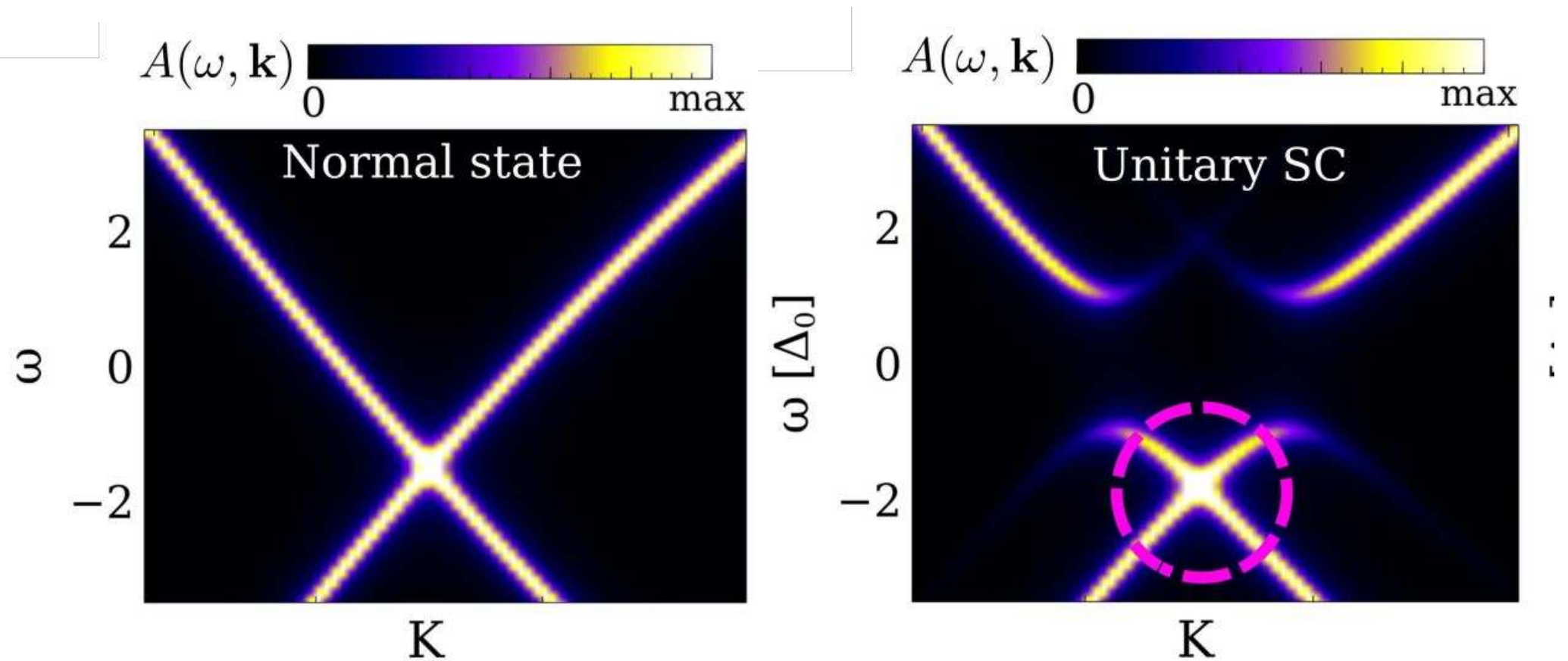
Dirac gap generation is a hallmark of non-unitary superconductivity

Dirac crossing in the superconducting state



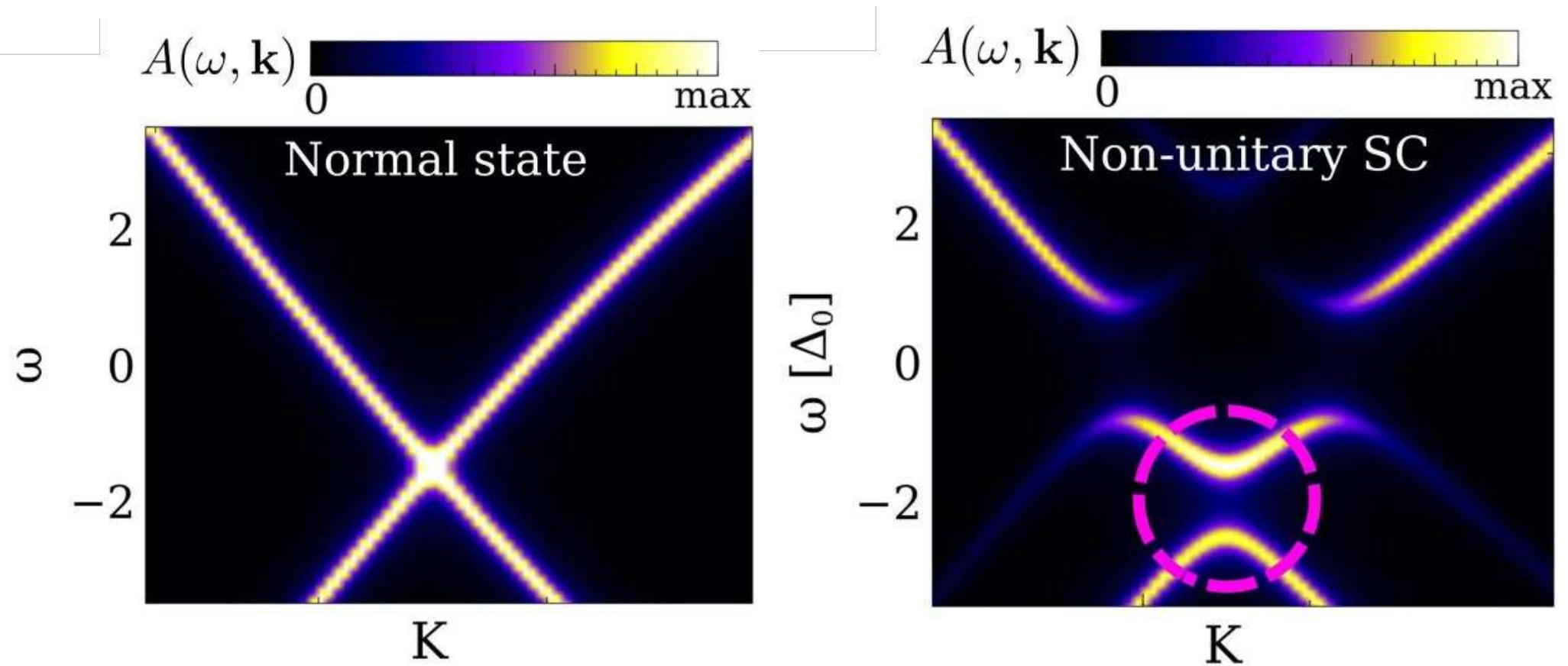
Dirac gap generation is a hallmark of non-unitary superconductivity

ARPES of a unitary SC



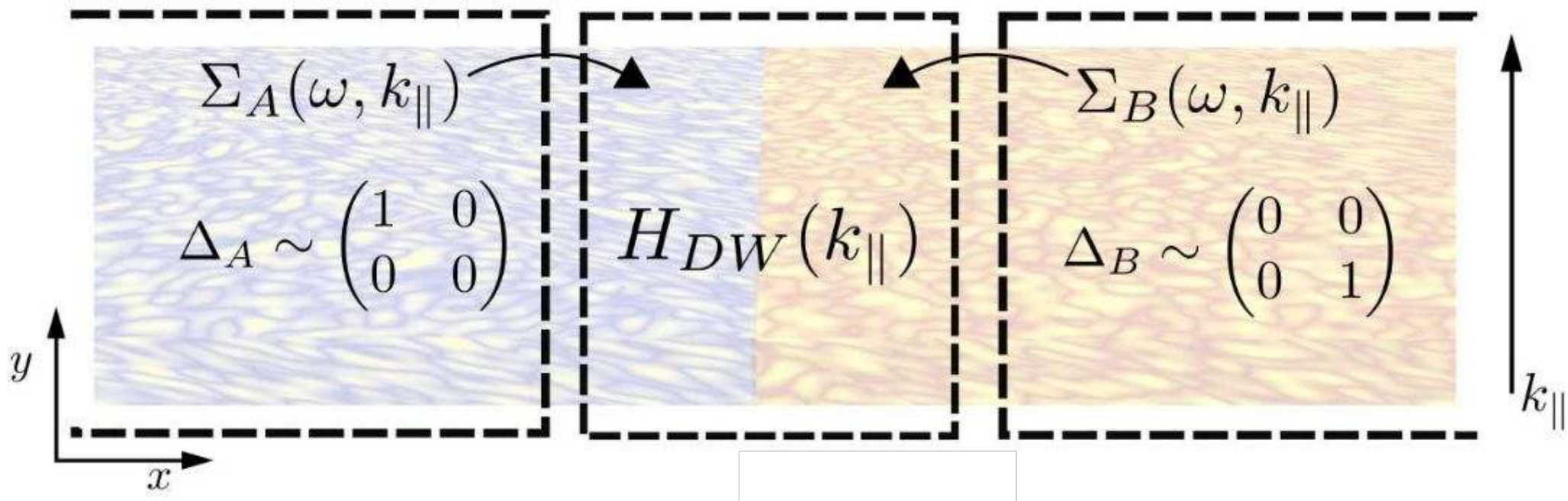
No gap generated at the Dirac point for a unitary state

ARPES of a non-unitary SC



Finite gap generated at the Dirac point for a non-unitary state

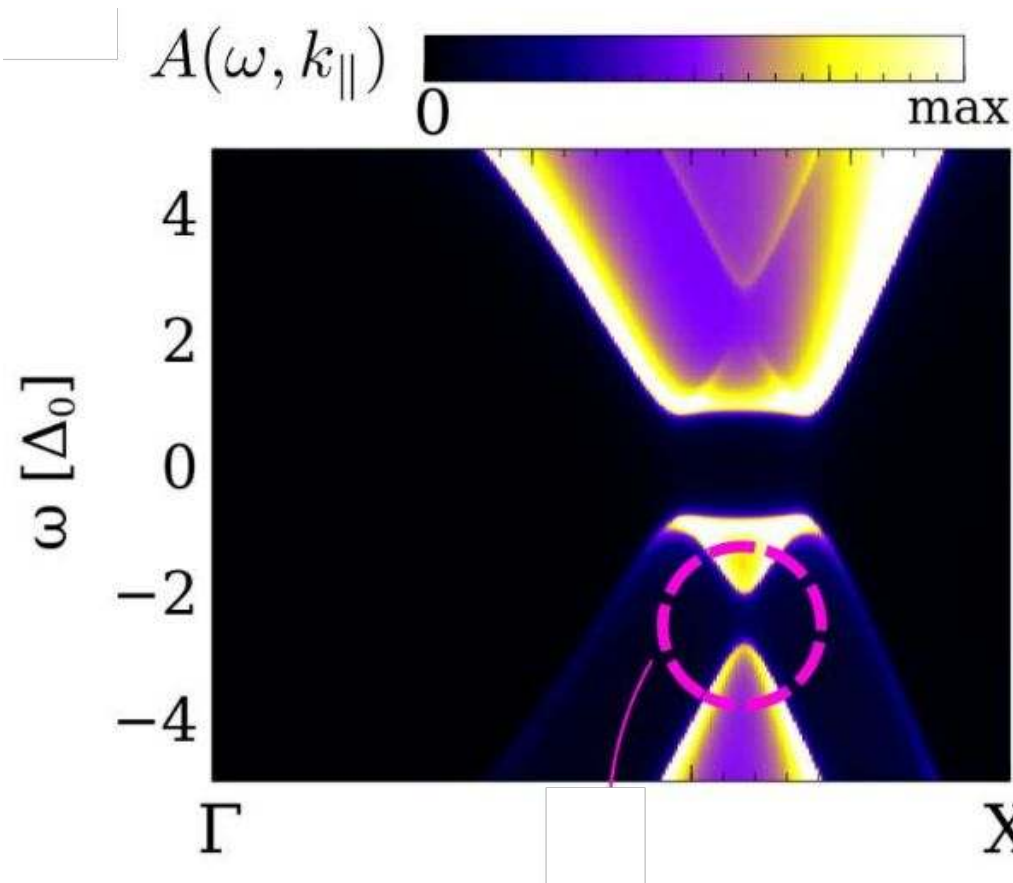
Superconducting orbital domain wall



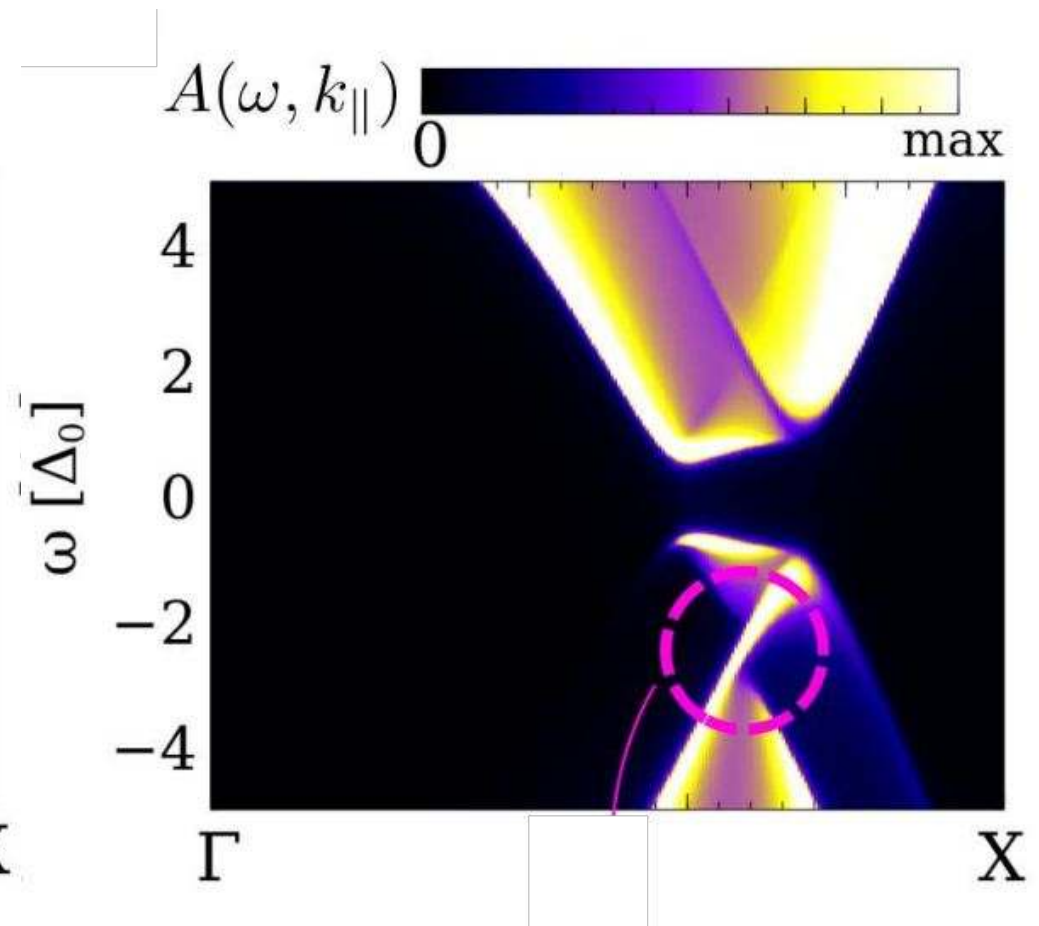
Is a superconducting domain wall visible in a spectral function?

ARPES of the SC with and without SC domain wall

Without domain wall



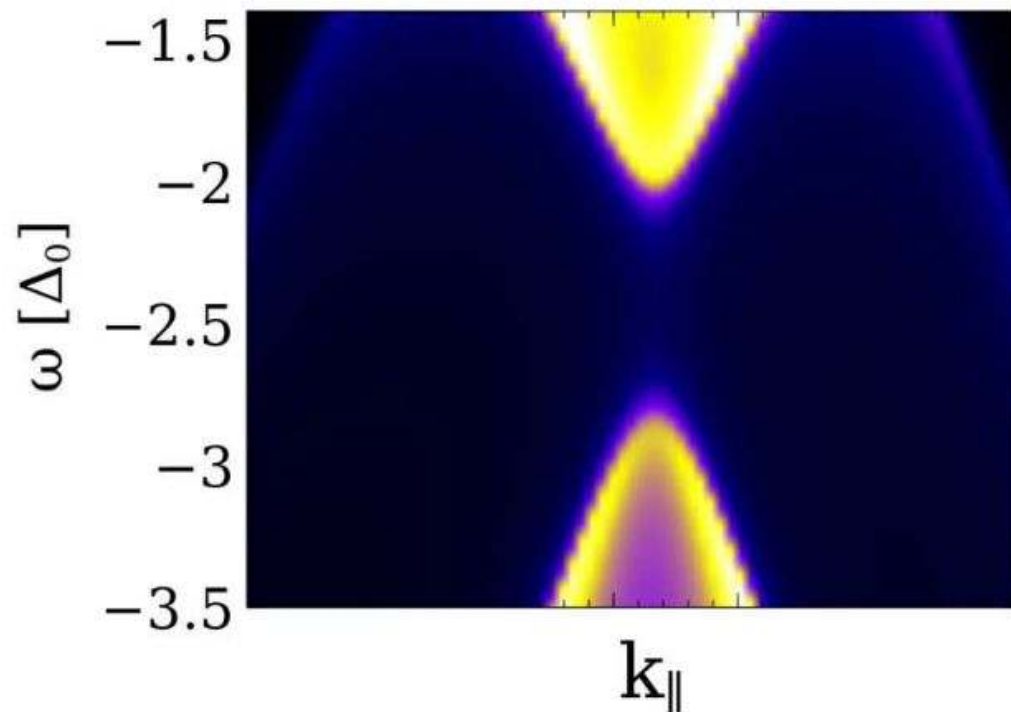
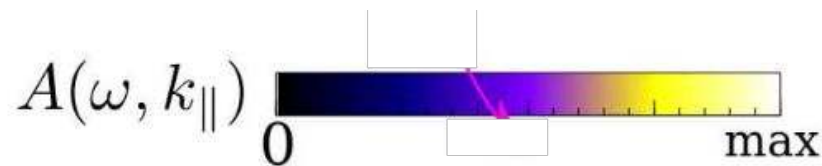
With domain wall



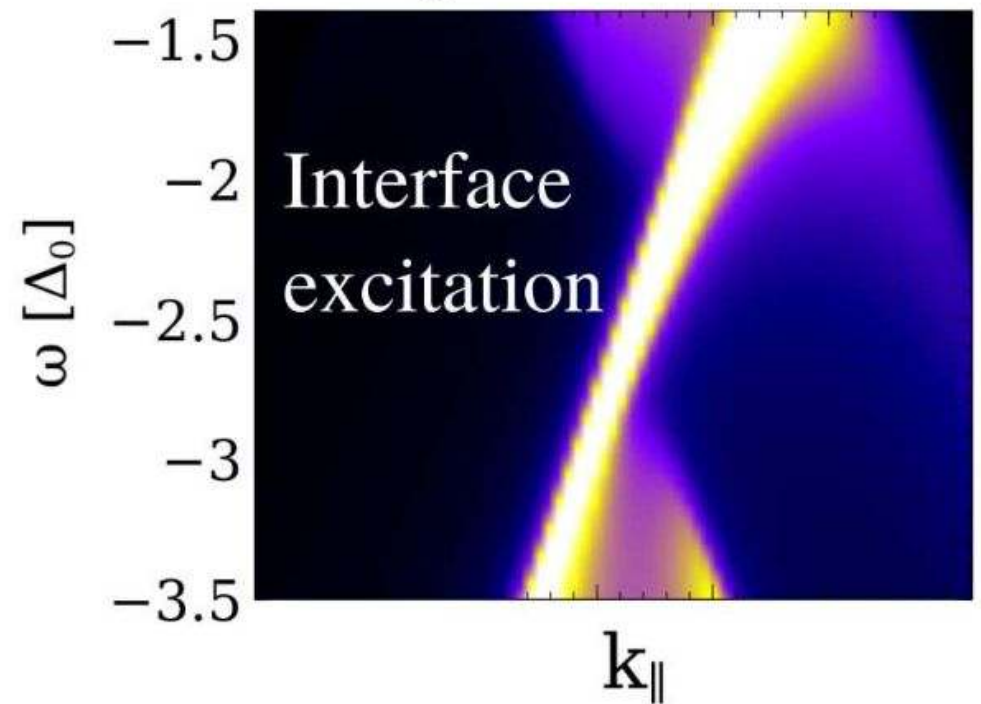
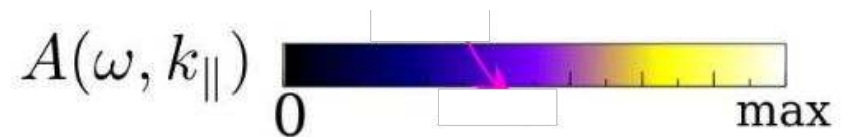
Superconducting domains generate states inside the Dirac gap

ARPES of the SC with and without SC domain wall

Without domain wall



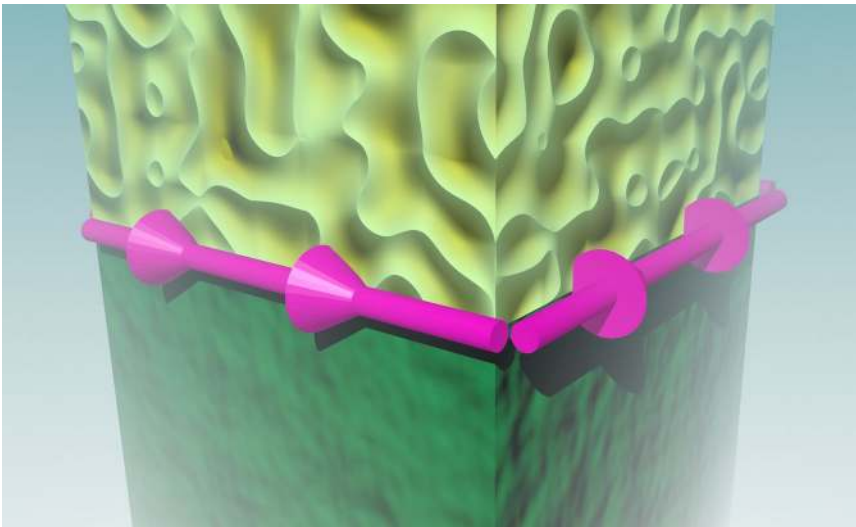
With domain wall



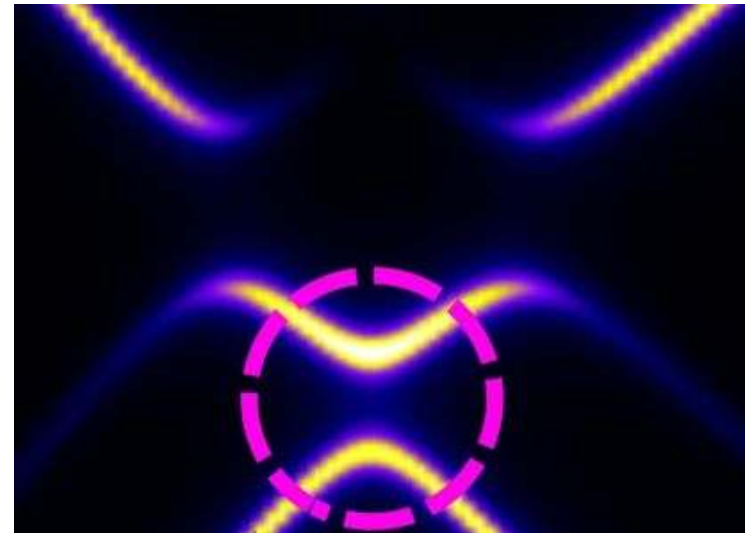
Superconducting domains generate states inside the Dirac gap

Take home

- **Antiferromagnetically gaped Dirac points allow to generate topological superconductors**
- **ARPES can detect non-unitary multiorbital superconductivity in doped Dirac materials**



Phys. Rev. Lett. 121, 037002 (2018)



arXiv:1908.04820 (2019)