Moire physics in graphene/magnetic van der Waals heterostructures

Jose Lado

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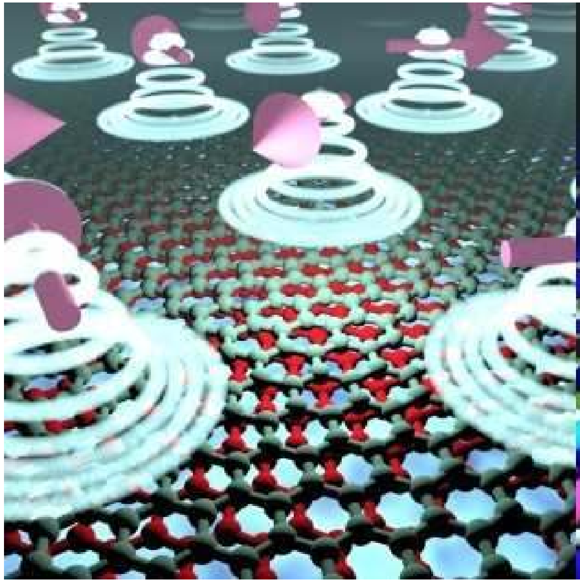
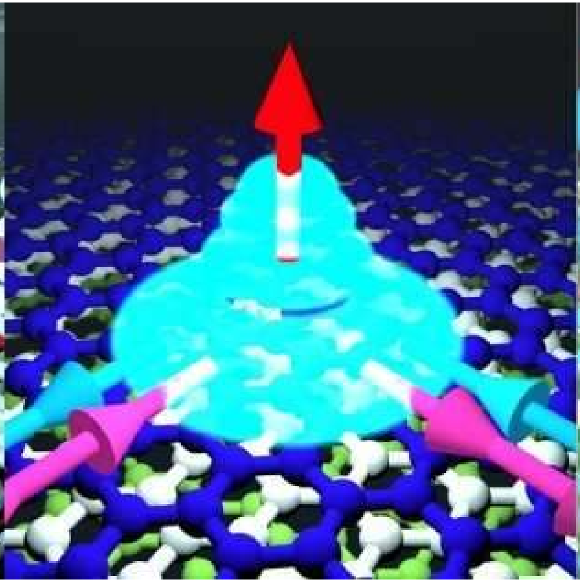
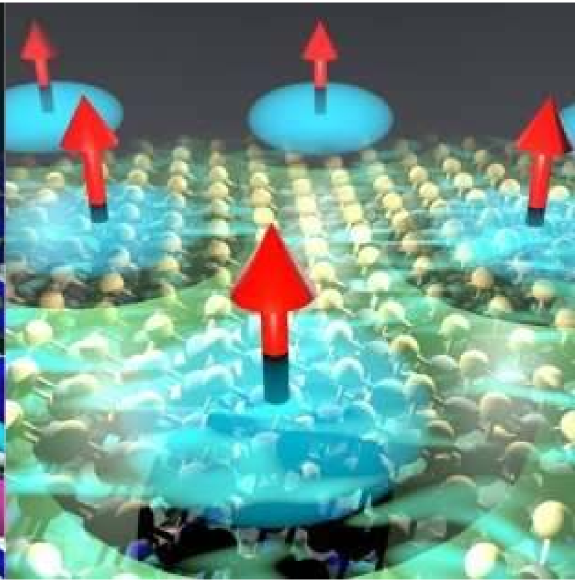


Phys. Rev. Lett. 126, 056803 (2021)



Phys. Rev. Lett. 127, 026401 (2021)

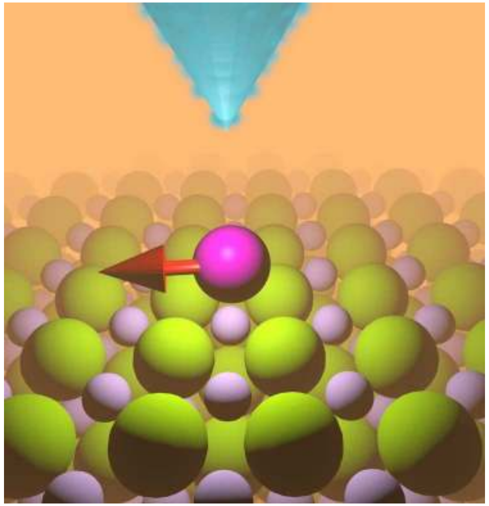


arXiv:2103.11989 (2021)

ICSM 2021, Milas-Bodrum, Turkey, October 23rd 2021

Building quantum matter with artificial lattices

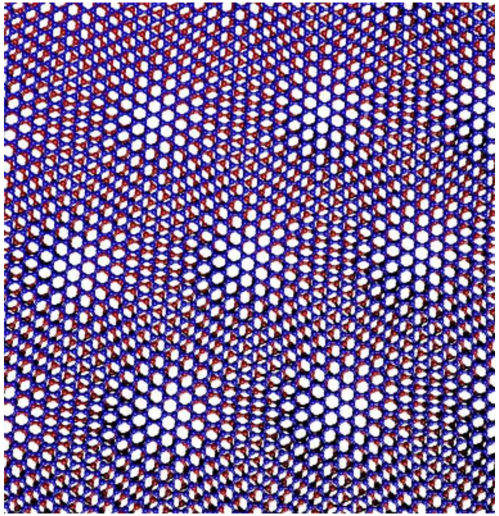
Atomic lattices



a = 0.5 nm

Nature Physics 12 (7), 656 (2016)
Nature Physics 13 (7) 668 (2017)
Science 335 (6065), 196-199 (2012)

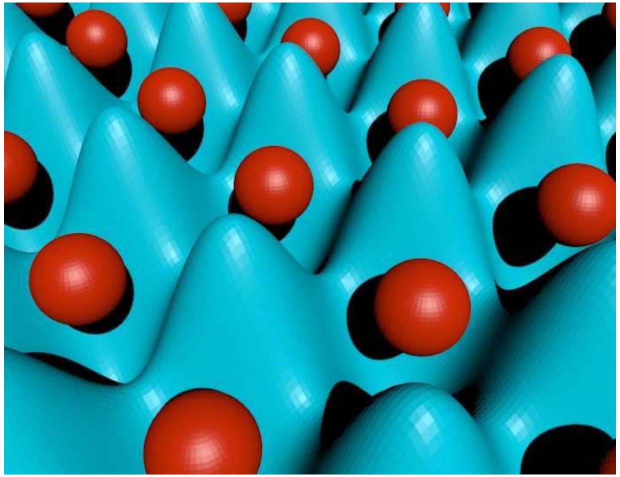
Twisted 2D moire materials



a = 30 nm

Nature 556, 80–84 (2018)
Science 363 (6431) 1059-1064. (2019)

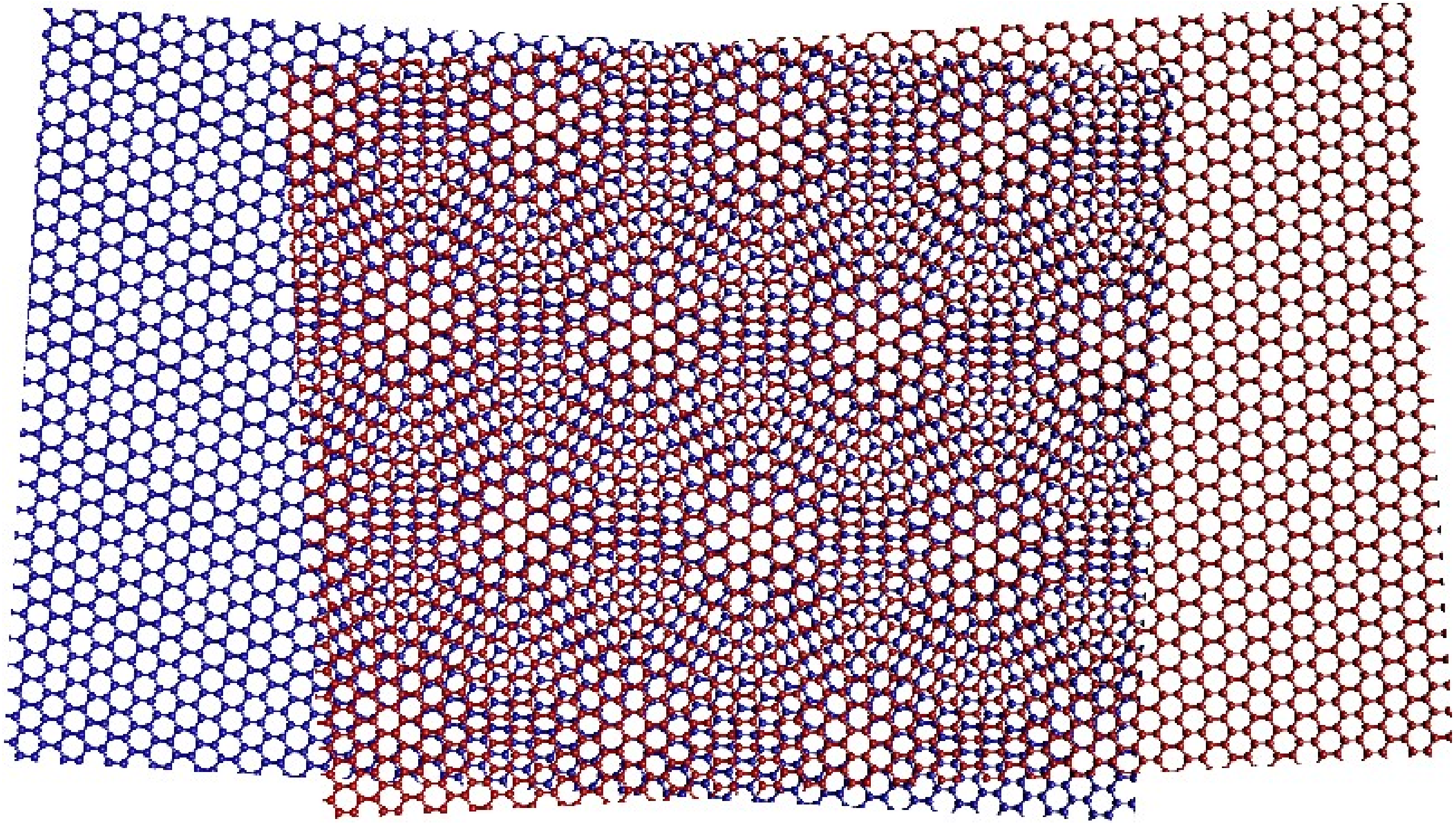
Cold atoms



a = 1000 nm

Nature 545.7655 (2017)
Nature 455.7210 (2008)
Phys. Rev. Lett. 111, 185307 (2013)

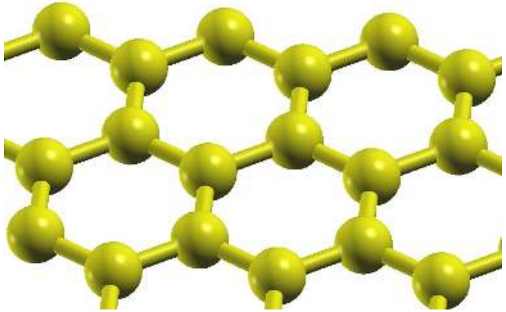
Twisted 2D materials as a platform for artificial quantum matter



The two-dimensional materials world

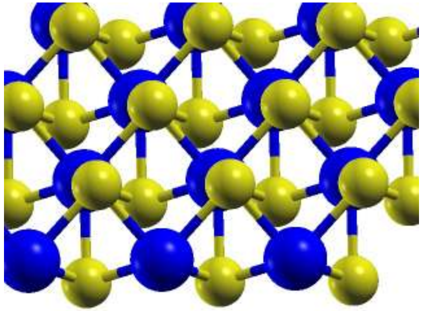
Semimetal

Graphene



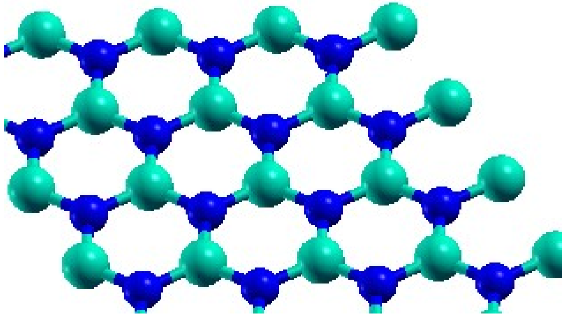
Superconductor

$NbSe_2$



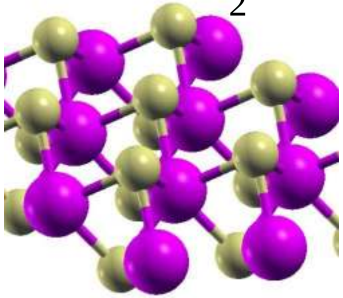
Insulator

BN



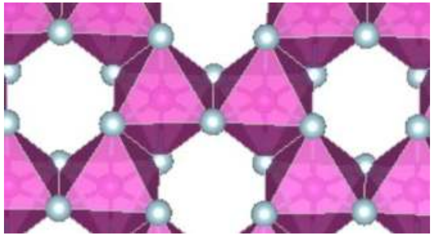
Semiconductor

WSe_2



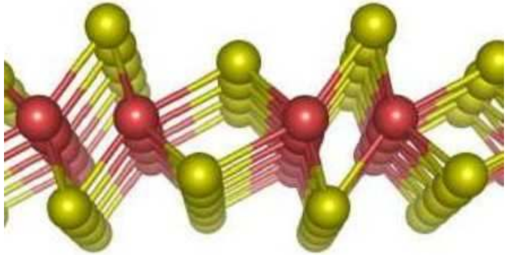
Ferromagnet

CrI_3



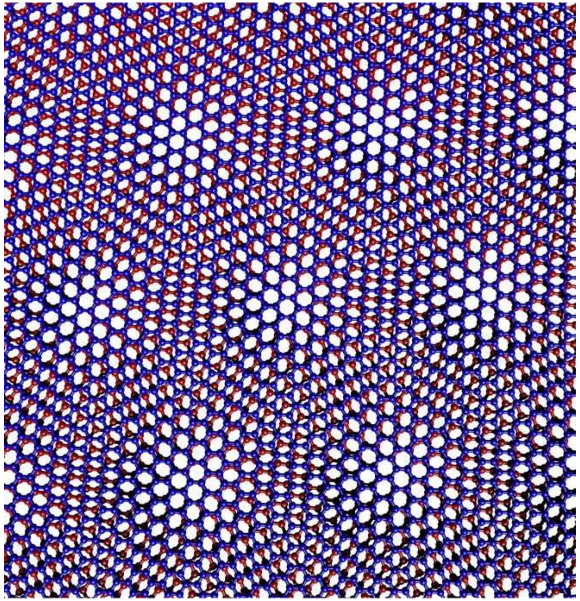
Quantum spin Hall insulator

WTe_2

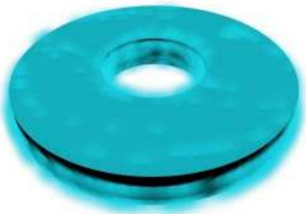


One material, a zoo of electronic phases

Twisted bilayer graphene



Superconductivity

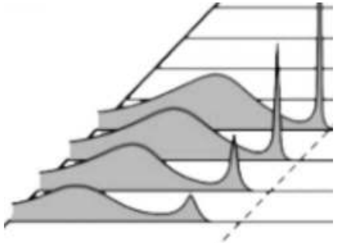


Nature 556, 43–50 (2018)

Topological networks

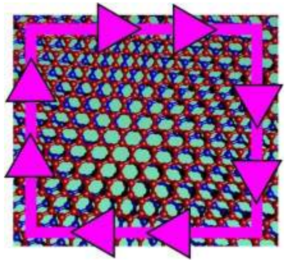
Nano Lett. 18, 11, 6725-6730 (2018)

Strange metals



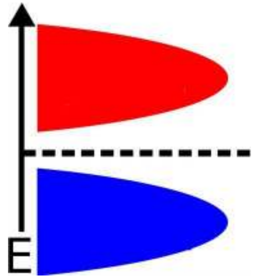
Phys. Rev. Lett. 124, 076801 (2020)

Chern insulators



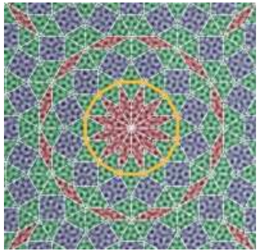
Science 365, 605-608 (2019)

Correlated insulators



Nature 556, 80–84 (2018)

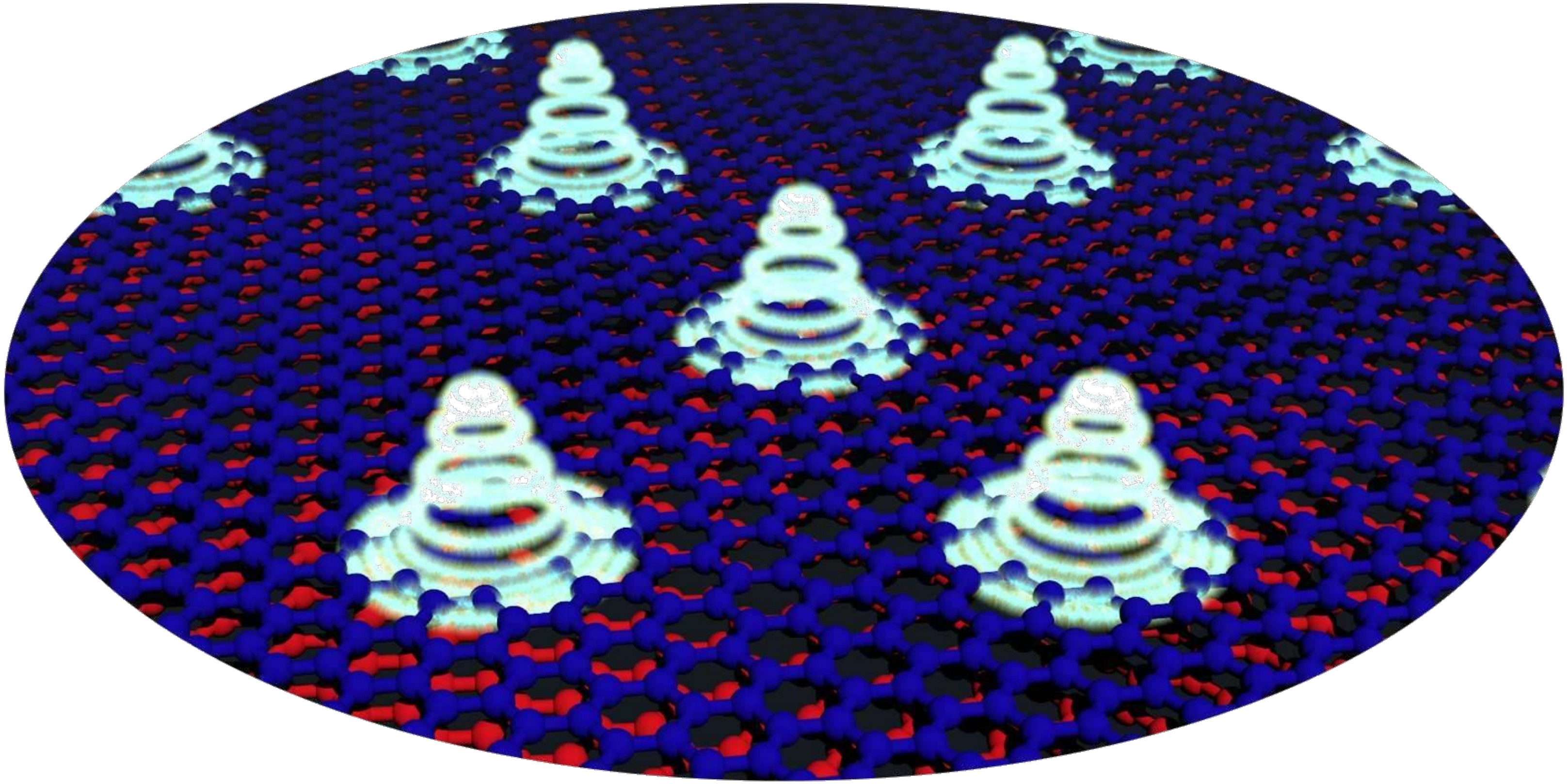
Quasicrystalline physics



Science 361 782-786 (2018)

Twisted graphene multilayers provide a powerful platform for emergent phenomena

A new universe in each new twisted material



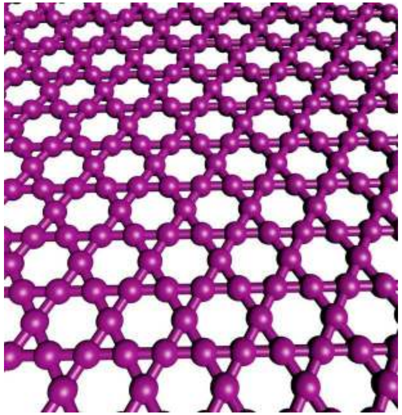
Each twisted material is a new universe for electrons, with laws that change from twist to twist

Designing quantum matter in twisted materials

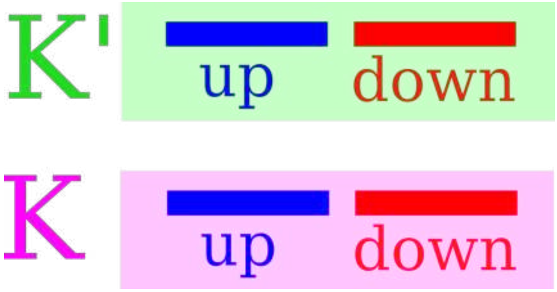
The strength of electronic correlations

$$V_{ijkl} c_i^{\dagger} c_j c_k^{\dagger} c_l$$

The details of the electronic dispersion



The local quantum degrees of freedom



Twist angle control

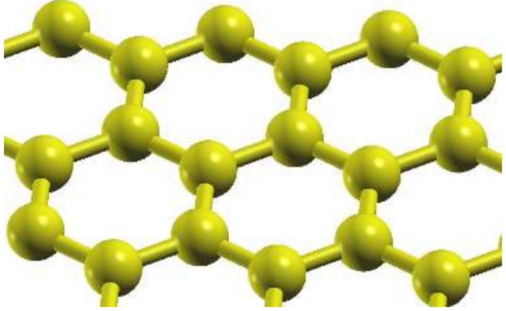
Interlayer bias

<u>Magnetic encapsulation</u>

Goal: control local degrees of freedom via magnetic encapsulation a twisted material

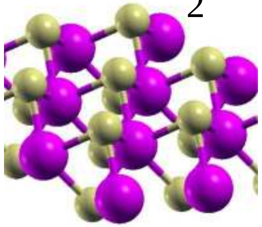
Designing quantum matter with twisted magnetic van der Waals materials

Graphene



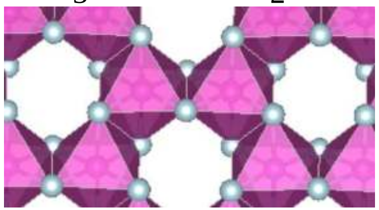
Semiconductor

WSe_2



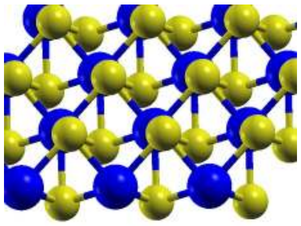
Magnet

$CrBr_3$, 1T-TaS_2, CrI_3



Metal

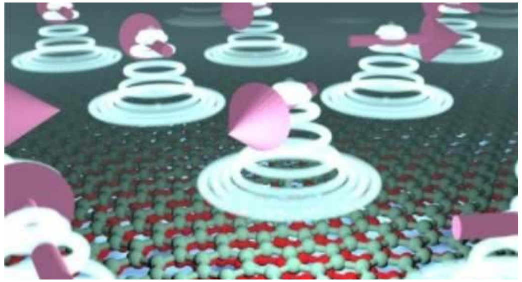
1H-TaS_2



Ingredients: Twist + exchange proximity + spin-orbit coupling + interactions

Behind the scenes

Valley-mixing correlations in CrI_3/twisted BG/CrI_3



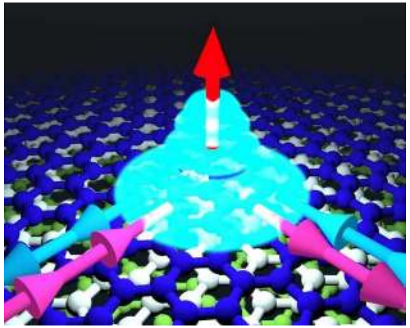
Tobias Wolf

Gianni Blatter

Oded Zilberberg

Phys. Rev. Lett. 126, 056803 (2021)

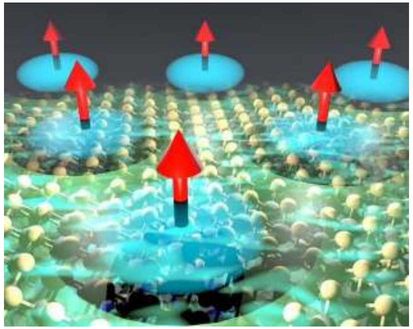
Heavy fermions in twisted graphene trilayers



Aline Ramires

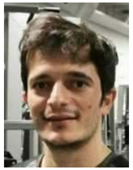
Phys. Rev. Lett. 127, 026401 (2021)

Heavy fermions in twisted dichalcogenide bilayers



S. Ganguli

M. Amini



G. Chen

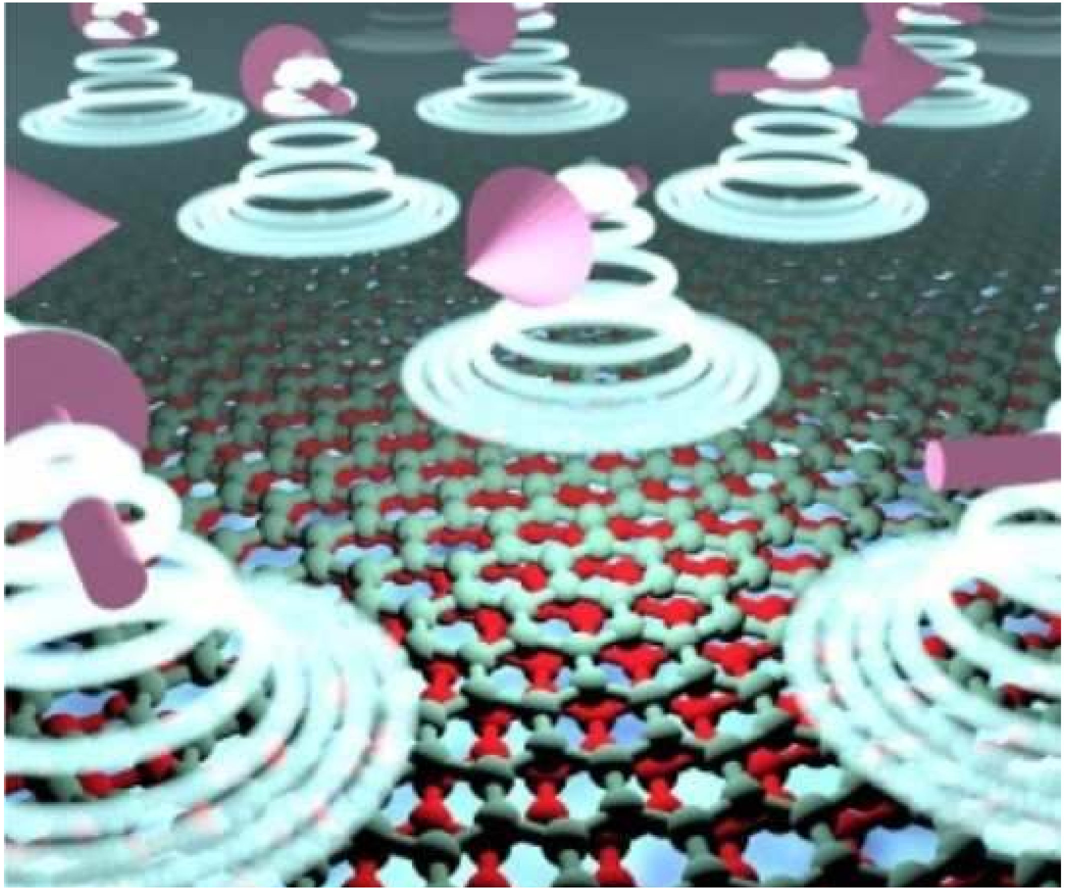
V. Vaňo

S. Kezilebieke

P. Liljeroth

arXiv:2103.11989 (2021)

Valley-mixing correlations in magnetically encapsulated tBG



Tobias Wolf

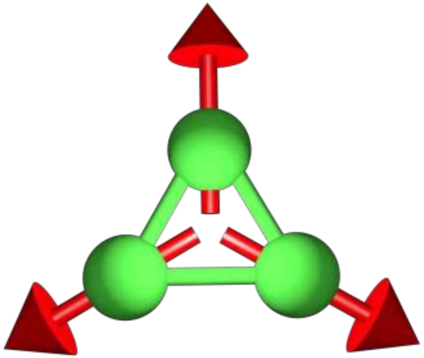
Gianni Blatter

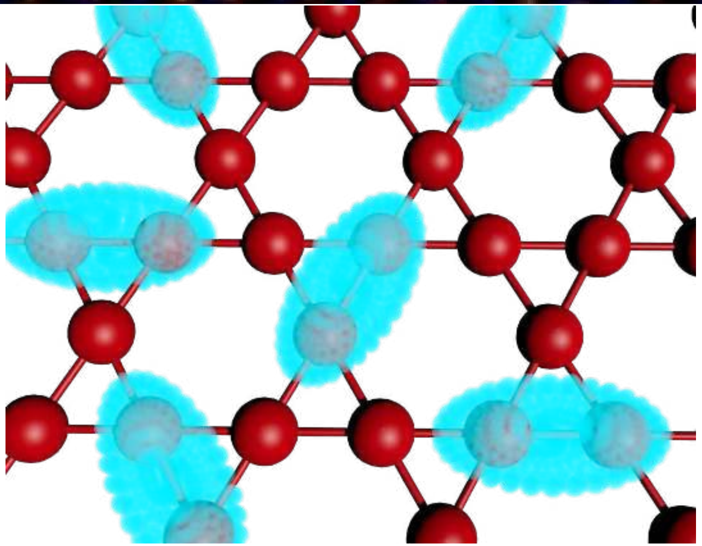
Oded Zilberberg

Phys. Rev. Lett. 126, 056803 (2021)

Quantum frustration in the valley sector

Frustrated magnetism





Quantum spin liquids

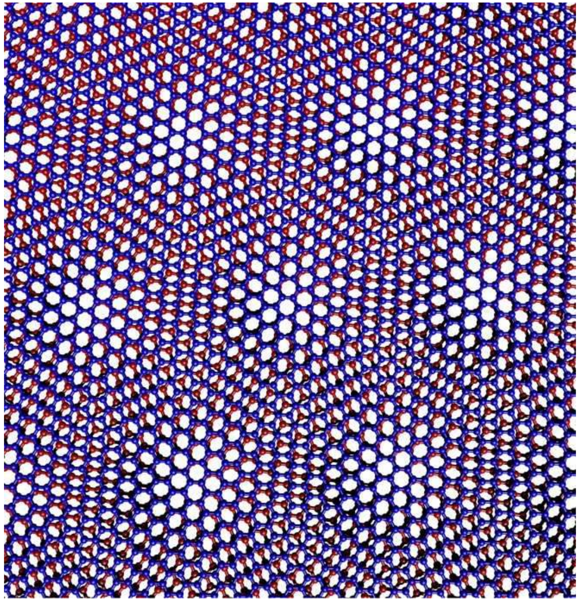
Non-trivial topological properties
Fractional excitations
Macroscopic quantum state
“Ultra-quantum” matter

$\equiv |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$

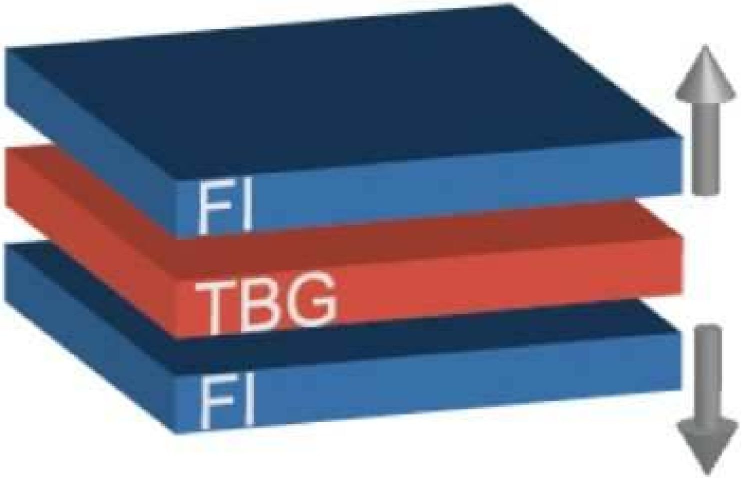
Can we create a controllable frustrated magnet, not in the spin but in the “valley” degree of freedom?

Magnetically encapsulated twisted graphene bilayer

Moire structure



Magnetic encapsulation



CrI_3/twisted bilayer graphene/CrI_3

Let us focus on the antiferromagnetic configuration of the encapsulation

Magnetically encapsulated twisted graphene bilayer

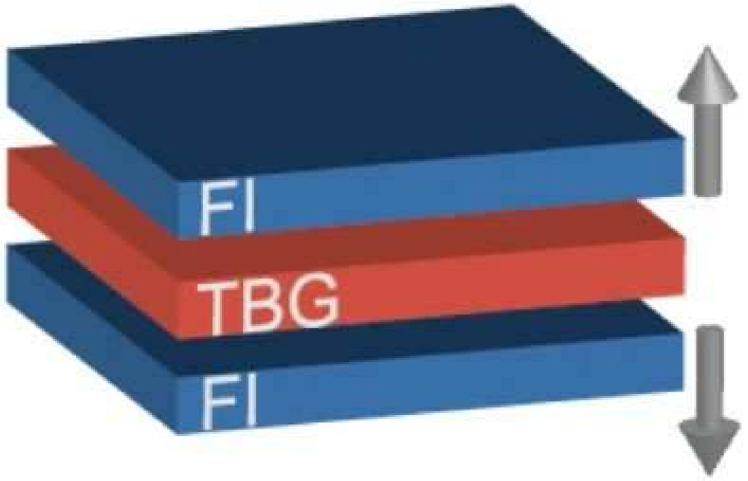
$$H = H_0 + H_\mathrm{J} + H_\mathrm{R}$$

Kinetic energy

Exchange proximity effect

Rashba spin-orbit coupling

Magnetic encapsulation



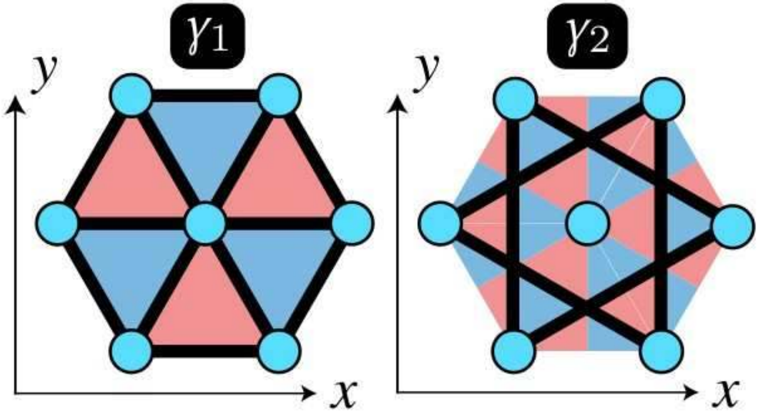
Valley frustration from combined twist engineering, exchange proximity and spin-orbit coupling.

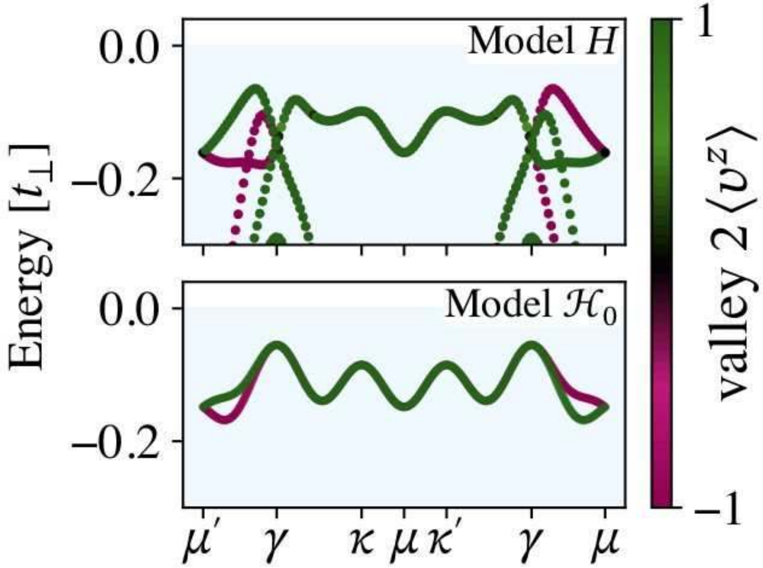
Lets focus on the regime $\alpha = 1.3^\circ$

Effective low energy valley model

Flux model in the triangular lattice

$$\mathcal{H}_0 = \sum_{\langle IJ \rangle} \gamma_1 \, \psi_I^\dagger e^{i\sigma^z \nu_{IJ} \phi_1} \psi_J + \sum_{\langle\langle IJ \rangle\rangle} \gamma_2 \, \psi_I^\dagger e^{i\sigma^z \nu_{IJ} \phi_2} \psi_J$$





Two-band valley-only model in the triangular lattice

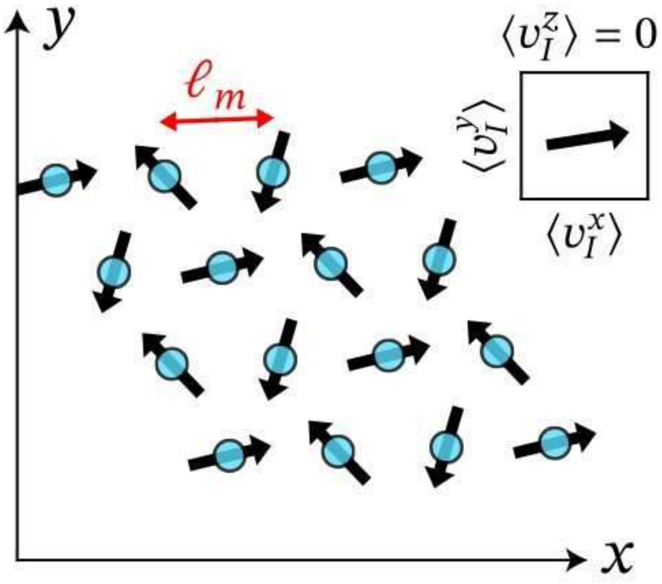
Impact of interactions

Lets add electronic interactions

Free Hamiltonian

Electronic interactions

$$\bar{H} = H + H_{int}$$

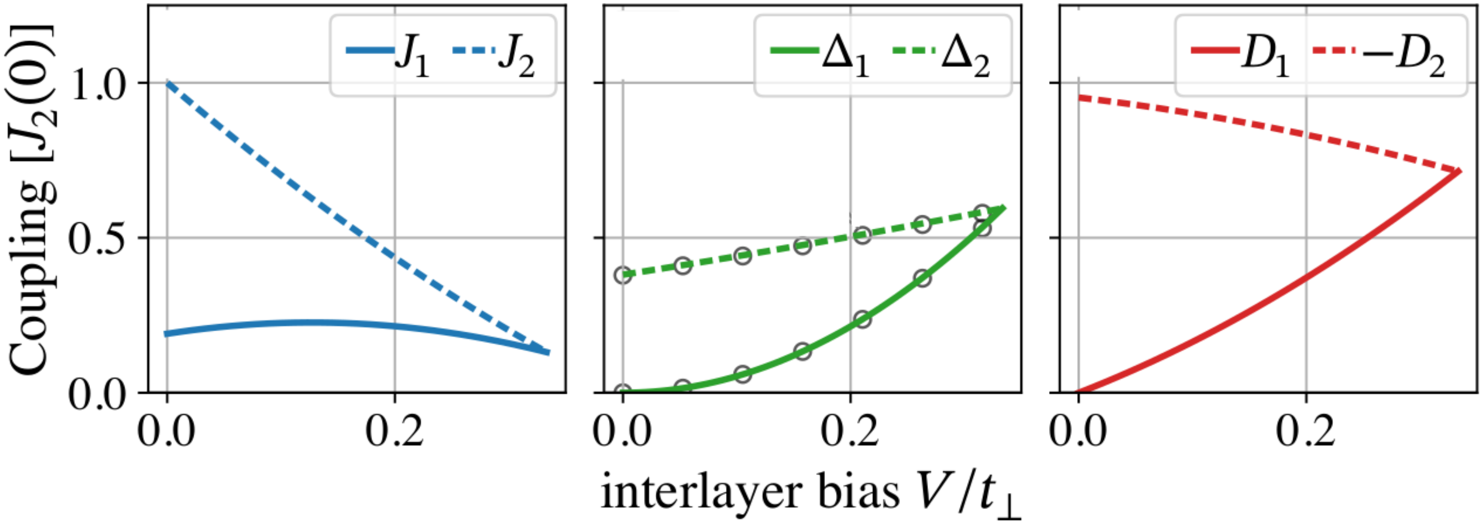


And compute the effective valley-Heisenberg Hamiltonian in the strong coupling limit

$$\mathcal{H}_v = \sum_{IJ} J_{IJ}\, \boldsymbol{v}_I \cdot \boldsymbol{v}_J + \Delta_{IJ}\, v_I^z v_J^z + \nu_{IJ}\, D_{IJ}\, (\boldsymbol{v}_I \times \boldsymbol{v}_J)_z$$

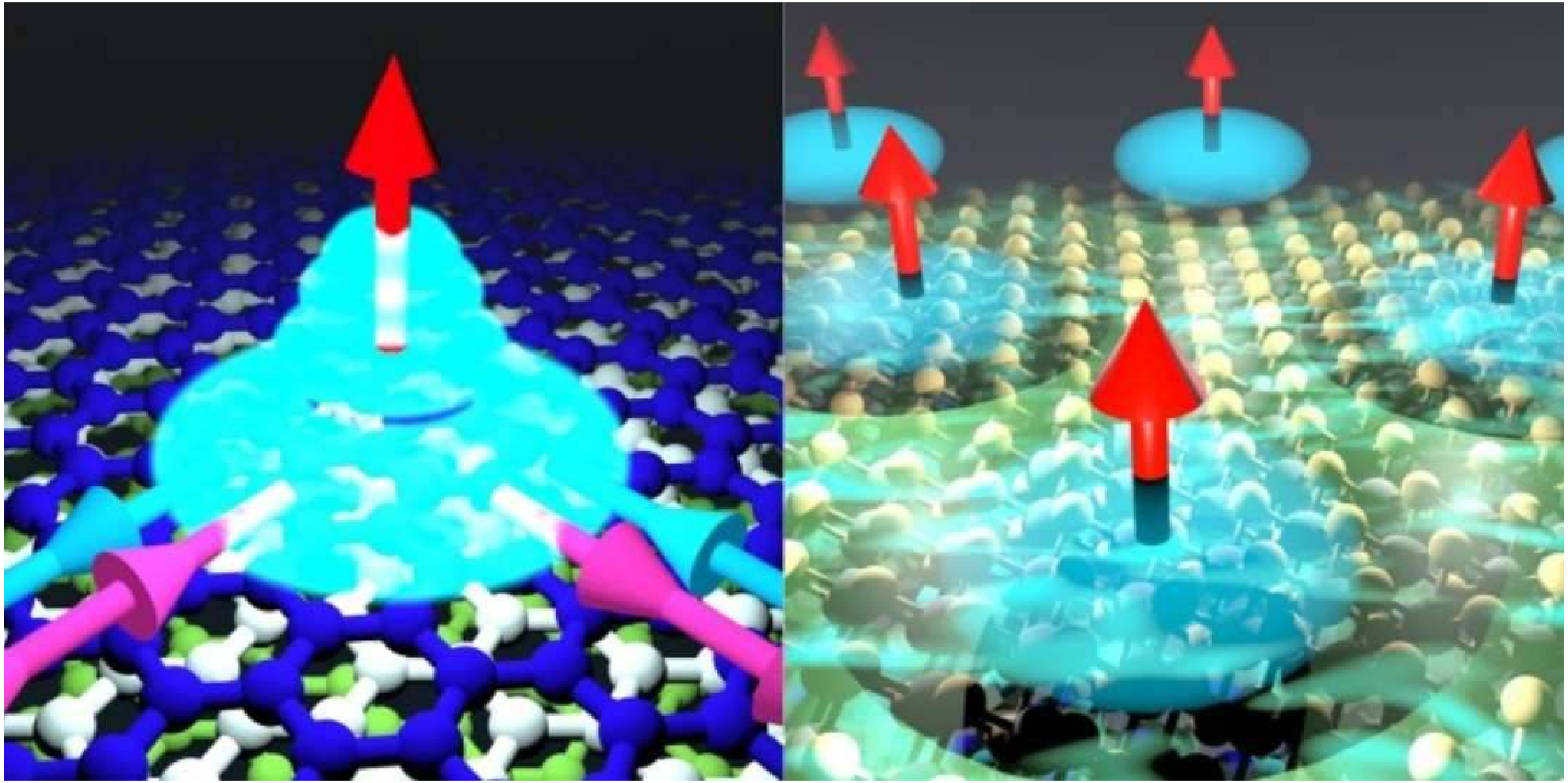
Controlling a valley-Heisenberg model electrically

$$\mathcal{H}_v = \sum_{IJ} J_{IJ}\, \boldsymbol{v}_I \cdot \boldsymbol{v}_J + \Delta_{IJ}\, v_I^z v_J^z + \nu_{IJ}\, D_{IJ}\, (\boldsymbol{v}_I \times \boldsymbol{v}_J)_z$$



An interlayer bias provides a tunable anisotropic valley-Heisenberg model

A new family of correlated states based on quantum magnetism: heavy fermions in twisted materials



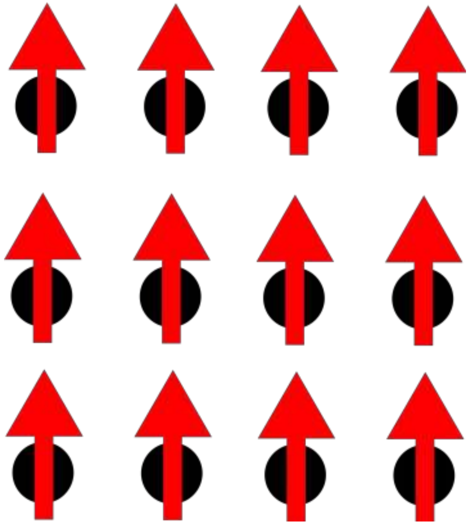
Basics of heavy fermion physics

- Strongly correlated matter, dominated by Kondo lattice physics
- Unique system to realize quantum criticality and unconventional superconductivity
- Found in rare-earth compounds, such as UTe_2 or UPt_3

Can we realize heavy-fermion physics with twisted two-dimensional materials?

Basics of heavy fermion physics

Lattice of Kondo impurities



f-electrons in
rare-earth compounds

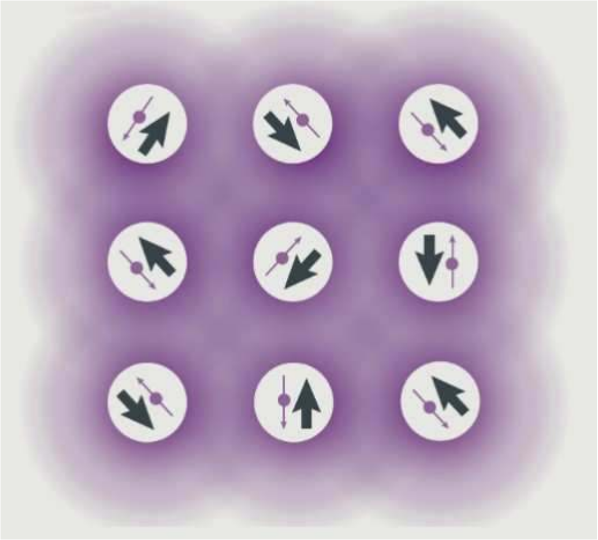
+

Dispersive electron gas

s/p/d-electrons in
rare-earth compounds

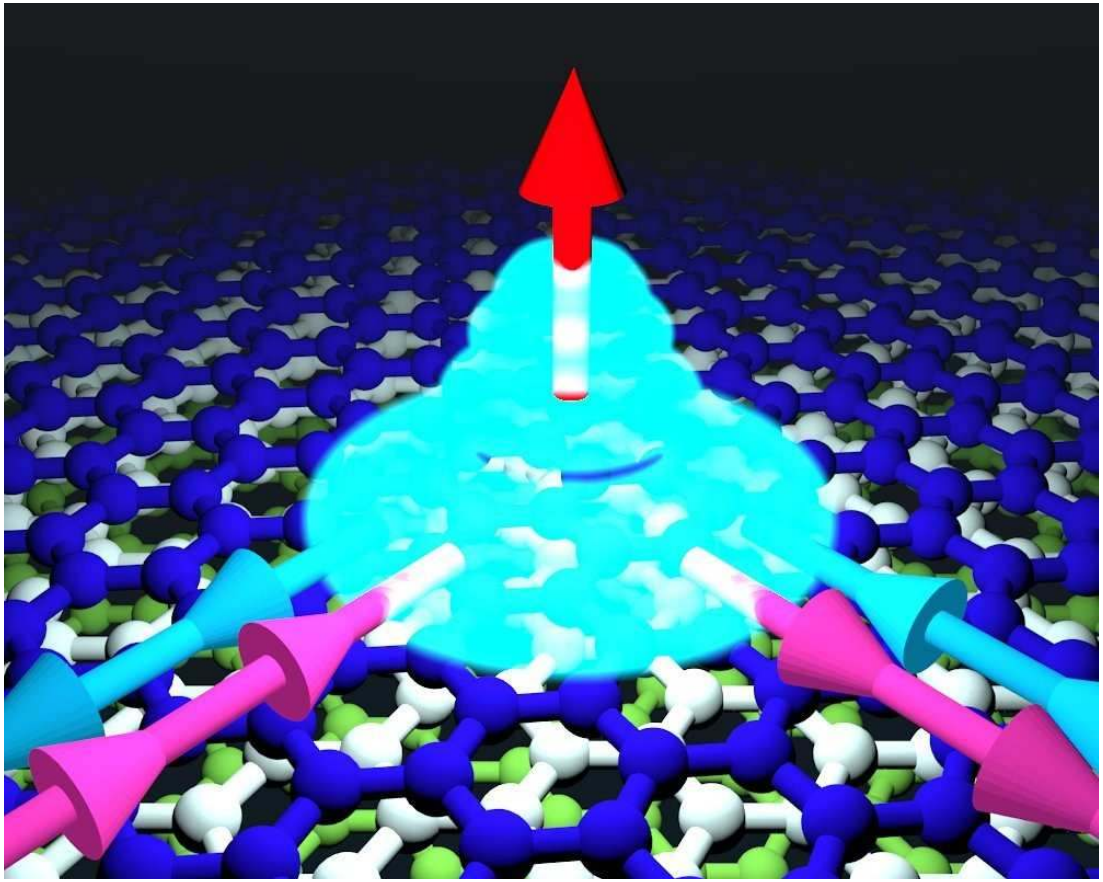
=

Kondo-lattice model



Associated with Kondo lattice physics:
- Colossal mass enhancement
- Quantum criticality
- Unconventional superconductivity

Heavy-fermions in twisted graphene trilayers

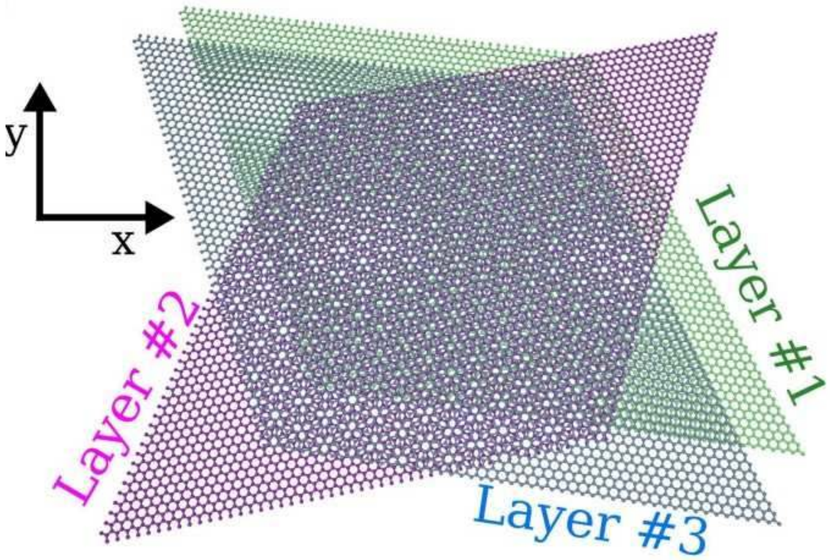


Aline Ramires

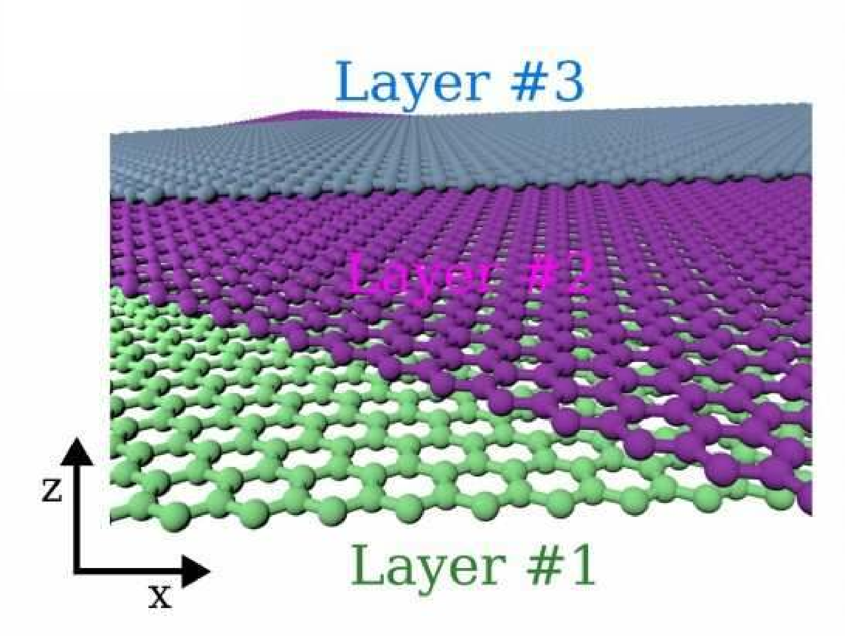
Phys. Rev. Lett. 127, 026401 (2021)

Heavy-fermions in twisted graphene trilayers

Top view

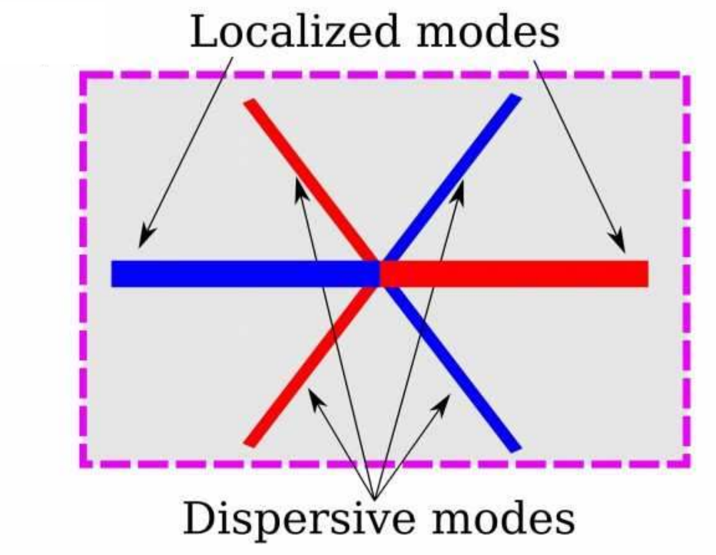


Lateral view

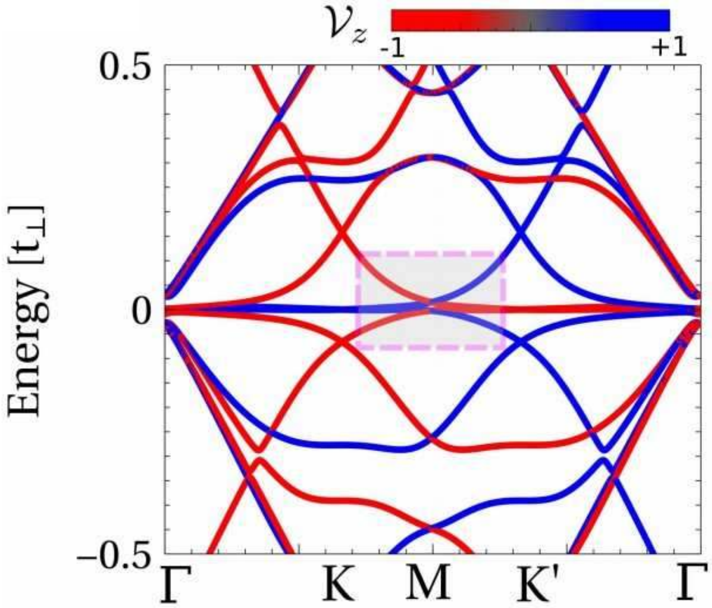


Heavy-fermions in twisted graphene trilayers

Low energy model

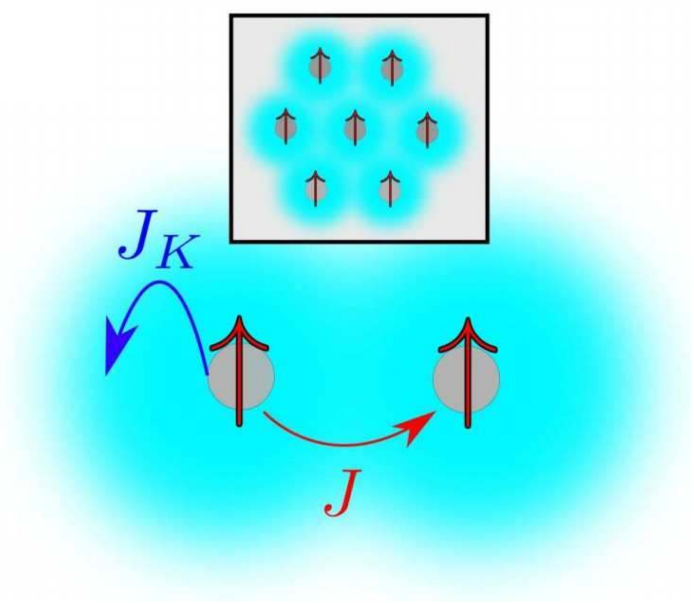


Electronic structure in the presence of an electric bias

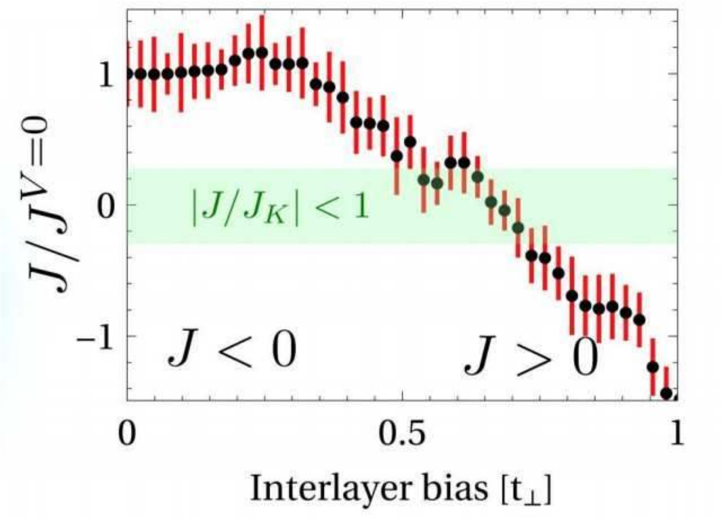


Kondo lattice model in the presence of interactions

Interactions of the Kondo lattice model



Electric control of the Kondo regime

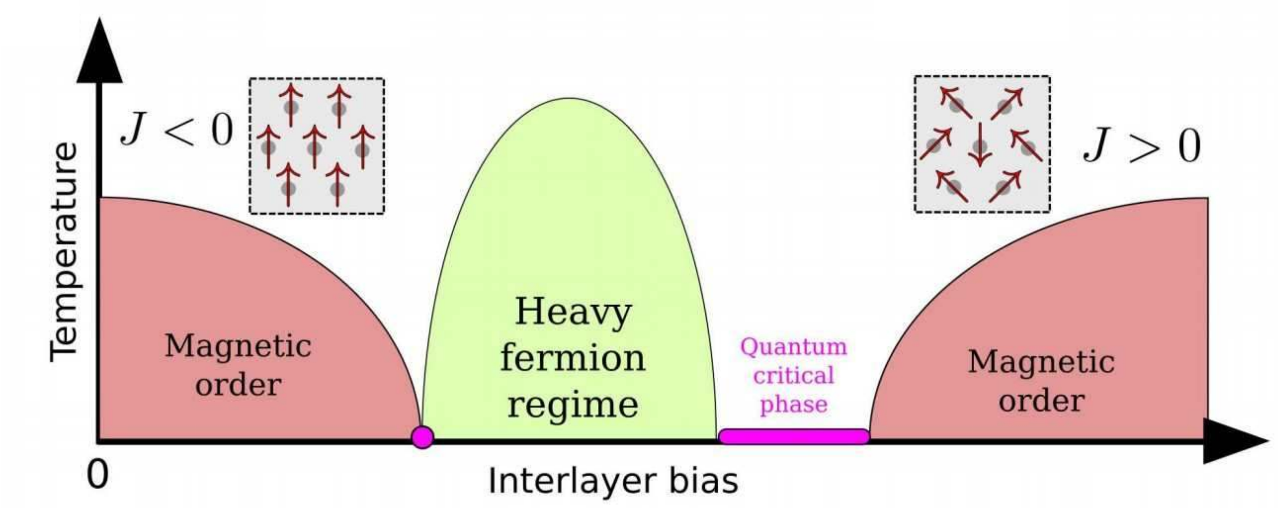


A heavy fermion regime appears when the Kondo coupling dominates over the exchange coupling

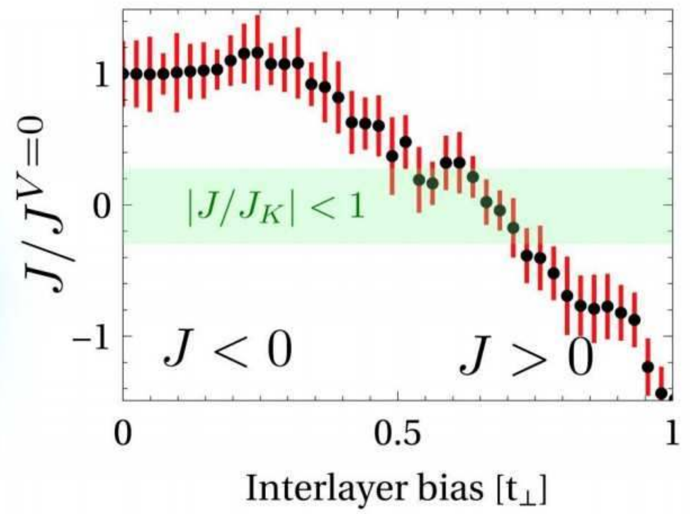
$$|J/J_K| < 1$$

Heavy-fermions in twisted graphene trilayers

Heavy-fermion (Doniach) phase diagram

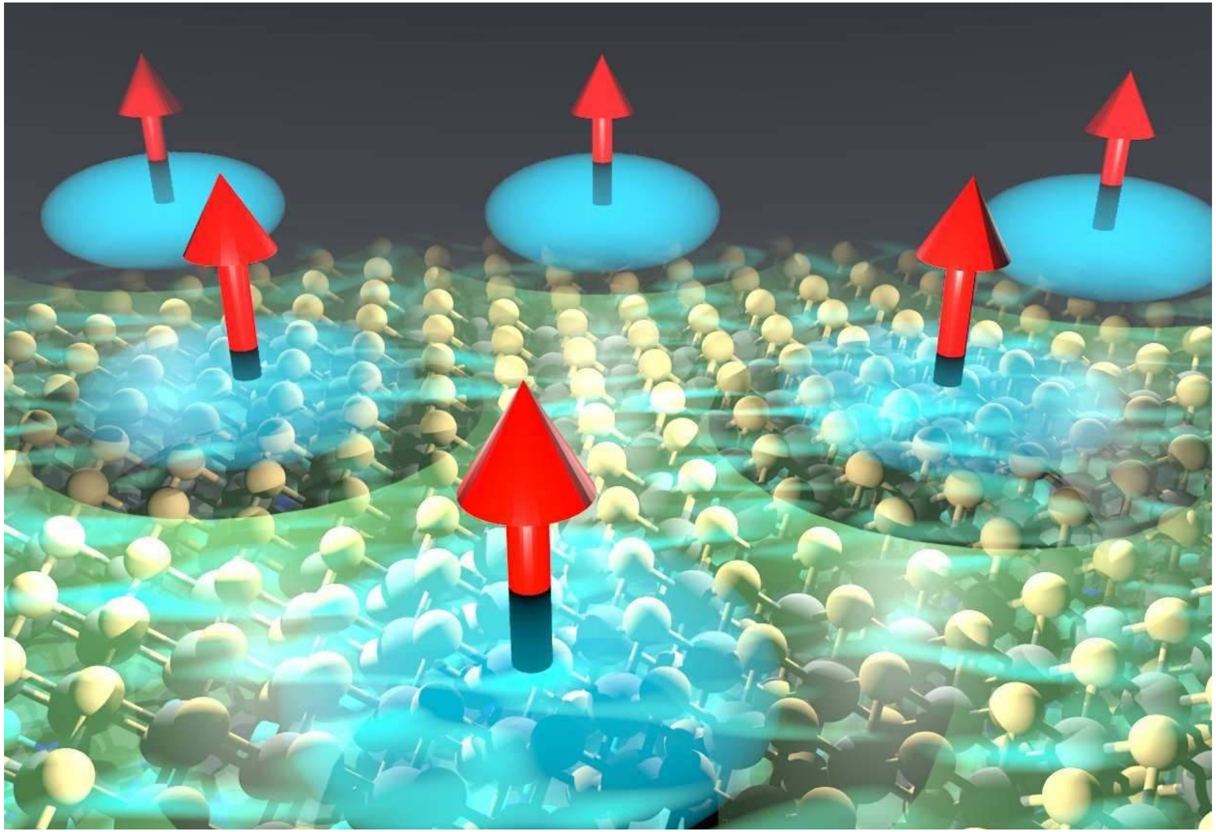


Electric control of the Kondo regime



The full phase diagram of the heavy fermion system can be explored in a single twisted graphene, by tuning the system with an electric bias

Heavy-fermions in a van der Waals dichalcogenide heterostructure



S. Ganguli

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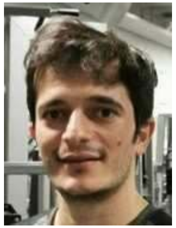


G. Chen

V. Vaňo

S. Kezilebieke

P. Liljeroth

arXiv:2103.11989 (2021)

Heavy-fermions in dichalcogenide bilayers

Triangular lattice of local magnetic moments

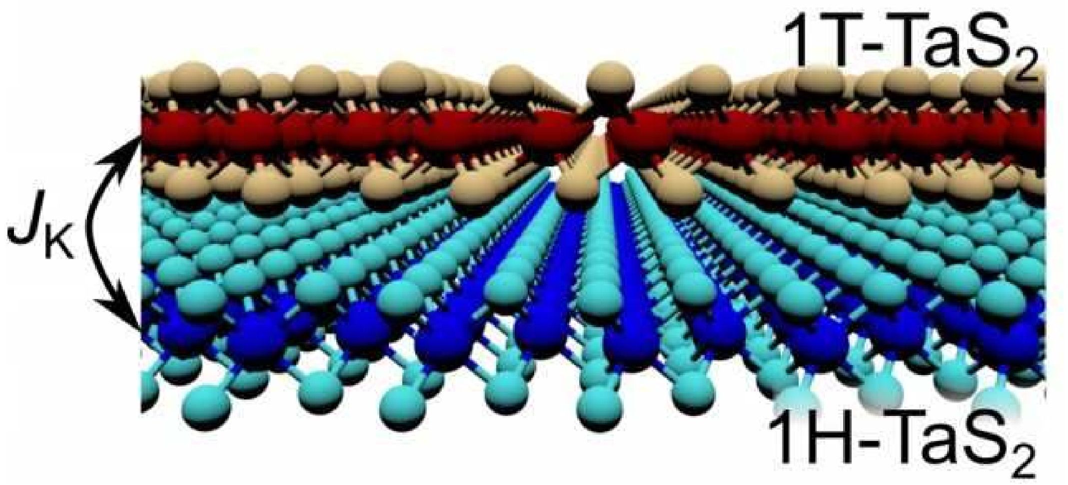
1T-TaS_2

Two-dimensional electron gas

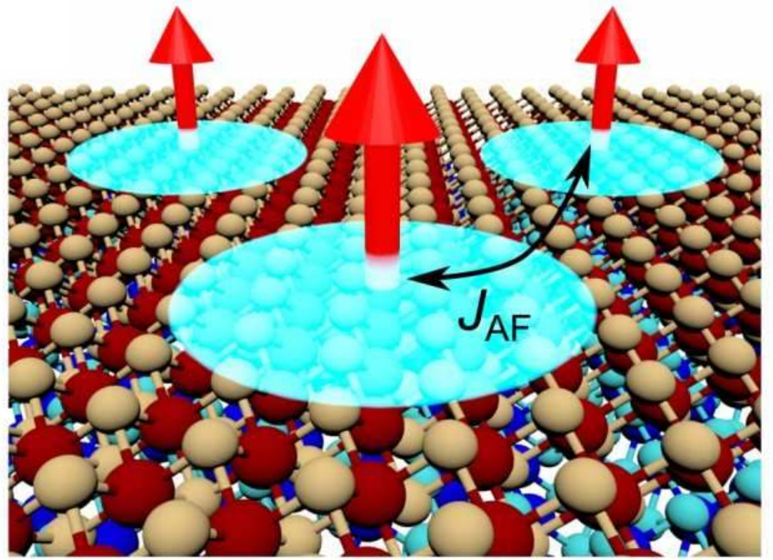
1H-TaS_2

Heavy-fermions in dichalcogenide bilayers

Interlayer coupling create the Kondo coupling



Intralayer exchange between local moments

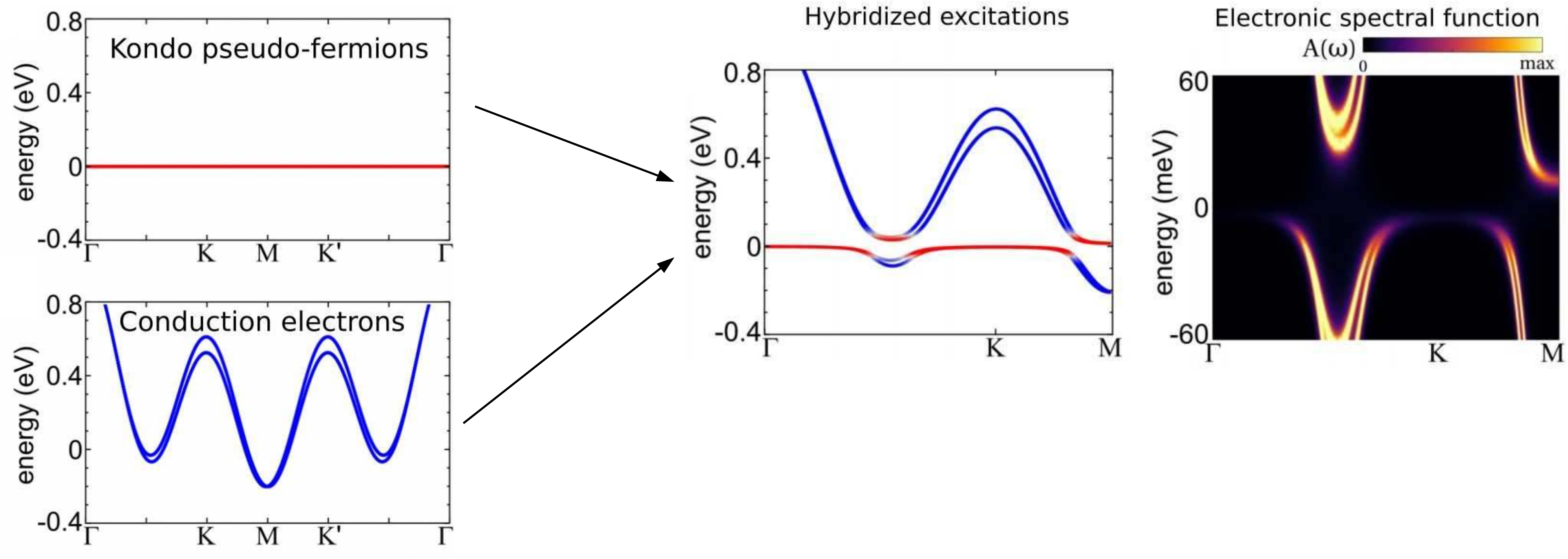


The bilayer realizes all the ingredients for heavy-fermion physics

Brief theory of heavy-fermions

Kondo physics introduces resonant pseudo-fermions at the chemical potential

Leading to the opening of a heavy-fermion gap

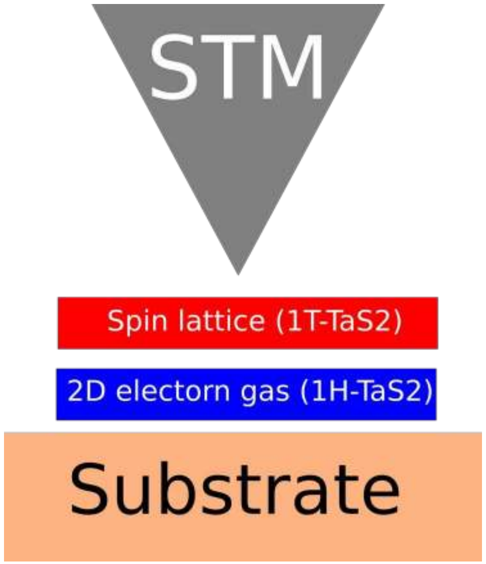


Experimental signatures of heavy-fermion physics

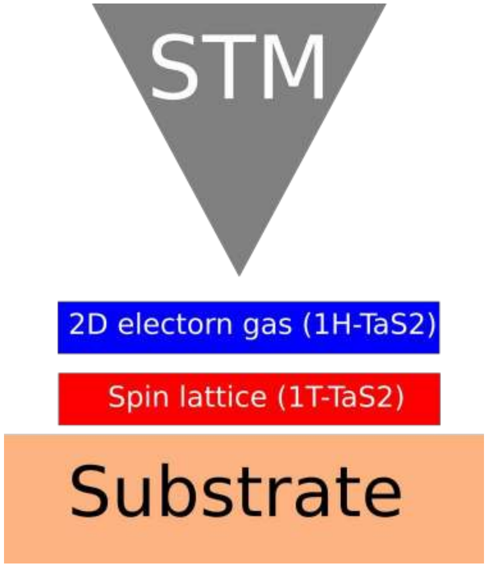
- Kondo physics in the magnetic lattice
- Gap opening in the metallic layer

Both signatures can be probed with scanning tunnel microscopy by growing two heterostructures

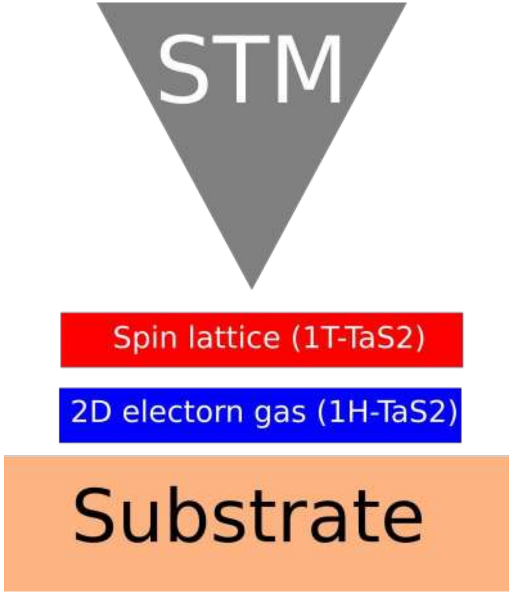
Probing the Kondo peak



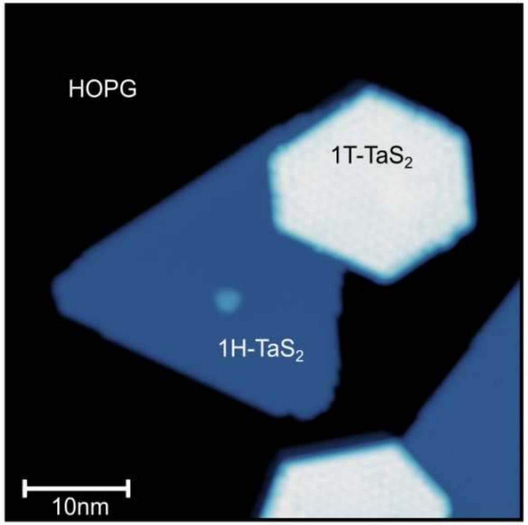
Probing the heavy-fermion gap



Kondo physics in the magnetic lattice

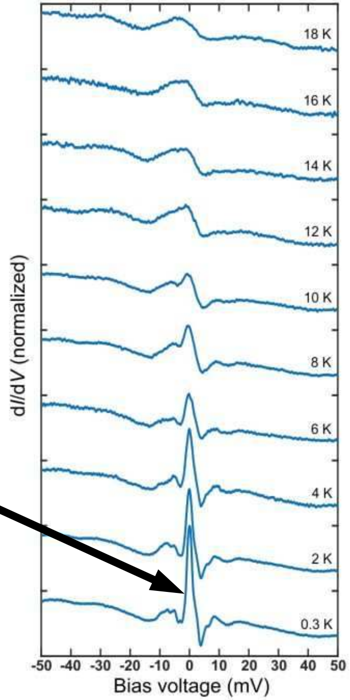


Sample



Differential conductance

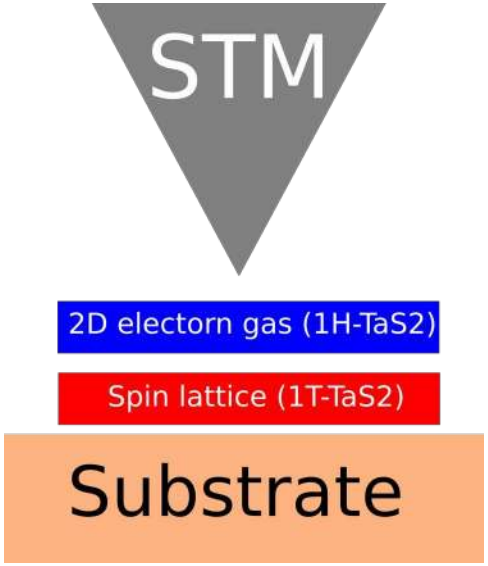
Kondo peak



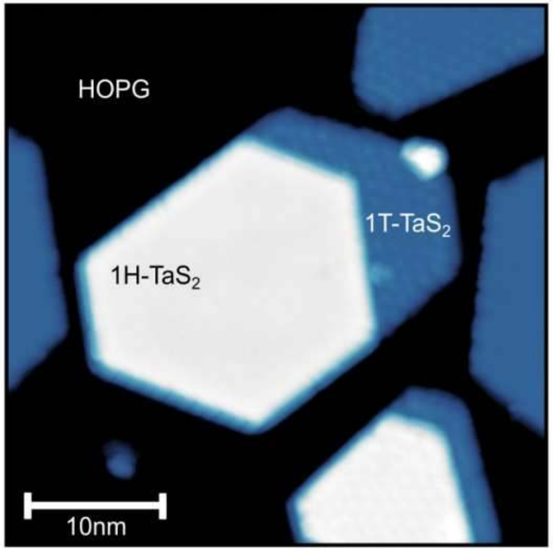
The system shows a Kondo peak when the local moments are probed

Temperature and field dependence consistent with Kondo

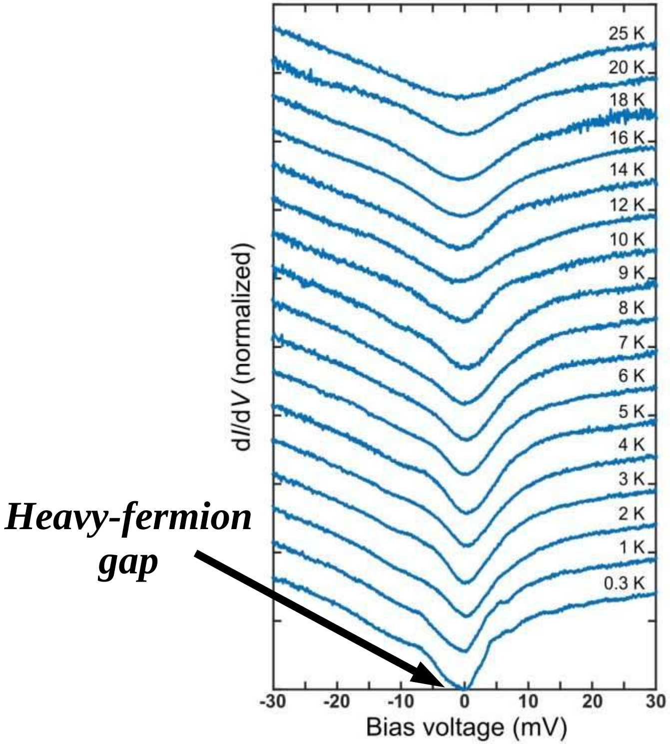
Heavy-fermion gap in the 2D electron gas



Sample



Differential conductance

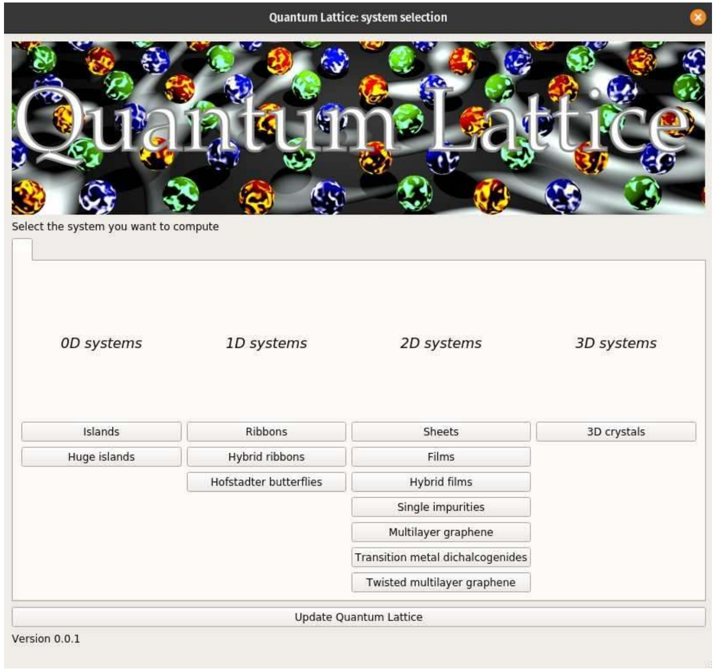


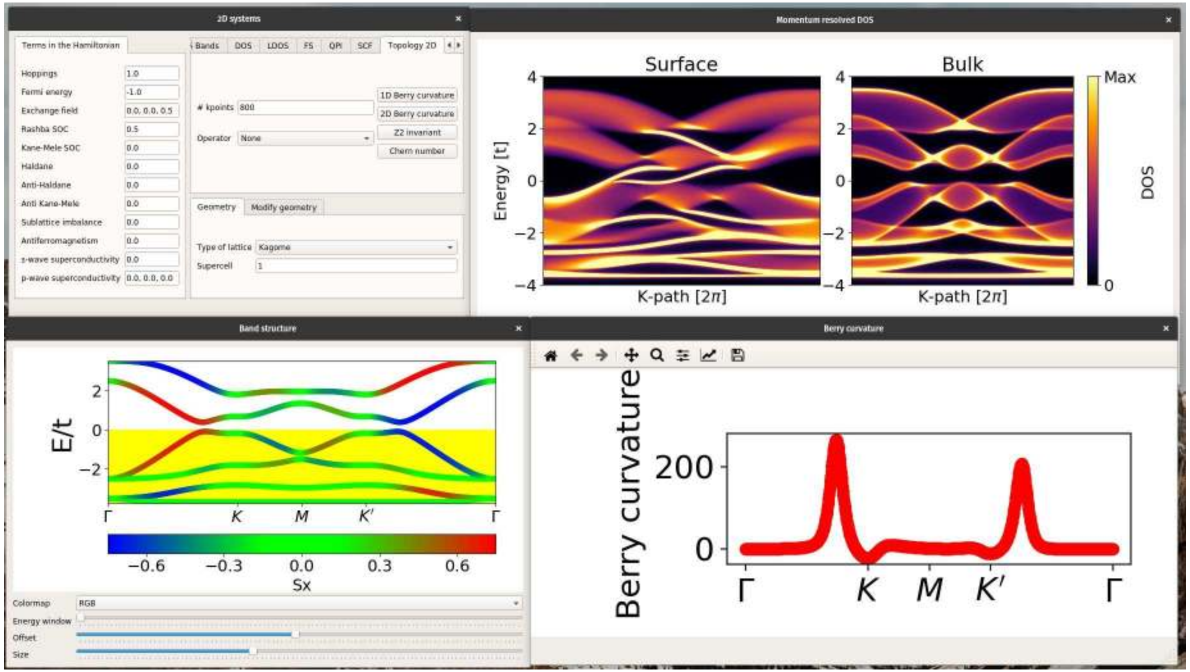
A heavy-fermion gap appears on the metal side

Temperature and field dependence consistent with a heavy-fermion gap

Computing electronic properties

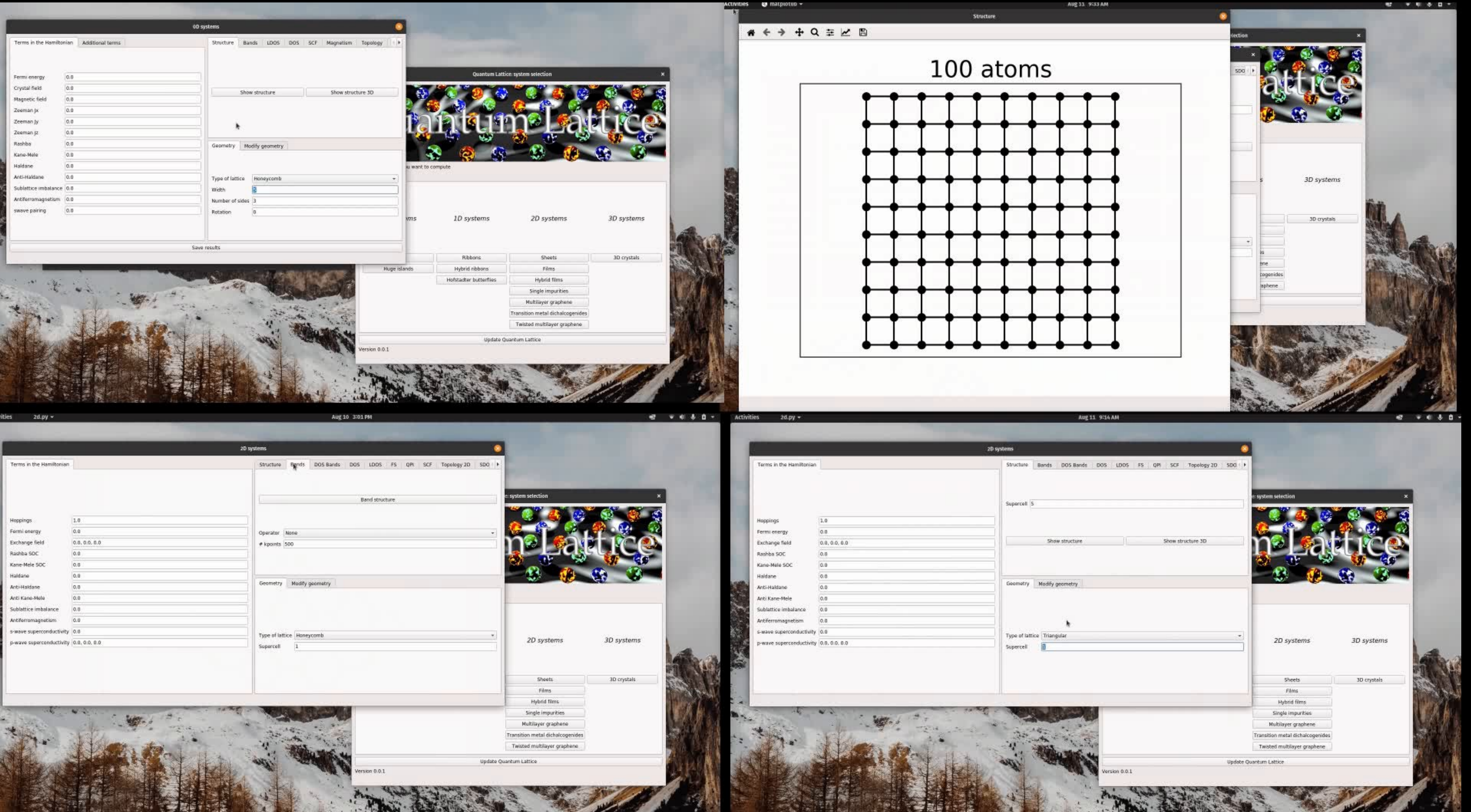
Quantum Lattice: A user interface to compute electronic properties





Quantum Lattice: open source interactive interface for tight binding modeling

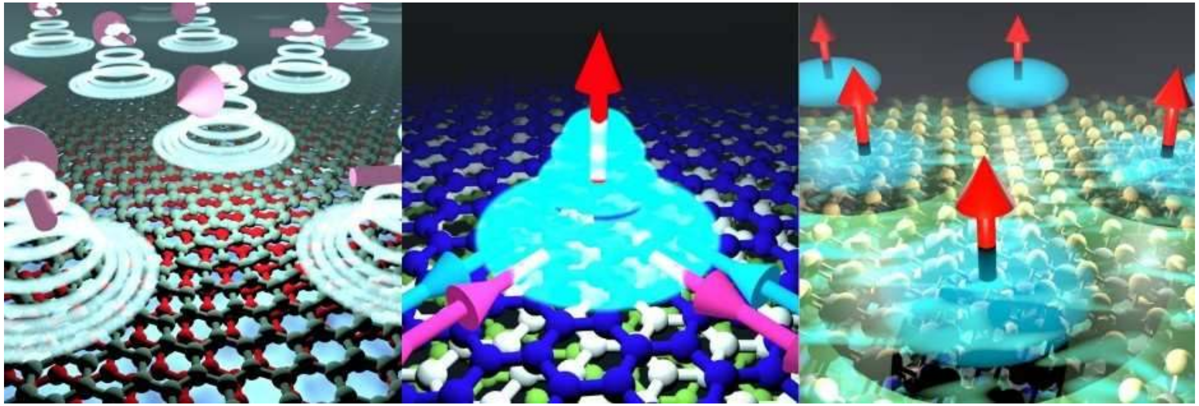
https://github.com/joselado/quantum-lattice



0D systems
Terms in the Hamiltonian
Additional terms
Structure
Bands
LDOS
DOS
SCF
Magnetism
Topology
Fermi energy
Crystal field
Magnetic field
Zeeman Jx
Zeeman Jy
Zeeman Jz
Rashba
Kane-Mele
Haldane
Anti-Haldane
Sublattice imbalance
Antiferromagnetism
swave pairing
Show structure
Show structure 3D
Geometry
Modify geometry
Type of lattice
Honeycomb
Width
Number of sides
Rotation
Save results
Quantum Lattice: system selection
1D systems
2D systems
3D systems
Ribbons
Sheets
3D crystals
Huge islands
Hybrid ribbons
Films
Hofstadter butterflies
Hybrid films
Single impurities
Multilayer graphene
Transition metal dichalcogenides
Twisted multilayer graphene
Update Quantum Lattice
Version 0.0.1
Structure
100 atoms
2D systems
Structure
Bands
DOS Bands
DOS
LDOS
FS
QPI
SCF
Topology 2D
SDO
Band structure
Operator
None
kpoints
500
Hoppings
Fermi energy
Exchange field
Rashba SOC
Kane-Mele SOC
Haldane
Anti-Haldane
Anti Kane-Mele
Sublattice imbalance
Antiferromagnetism
s-wave superconductivity
p-wave superconductivity
Type of lattice
Honeycomb
Supercell
Triangular

Take home

Magnetic twisted van der Waals materials provide a platform to engineer a new family of correlated states, including exotic heavy-fermion physics



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Phys. Rev. Lett. 127, 026401 (2021)
arXiv:2103.11989 (2021)

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SUOMEN AKATEMIA