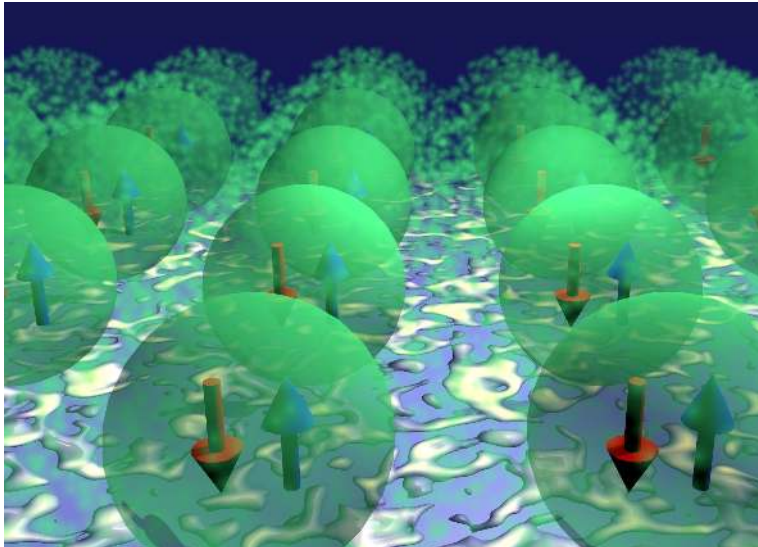


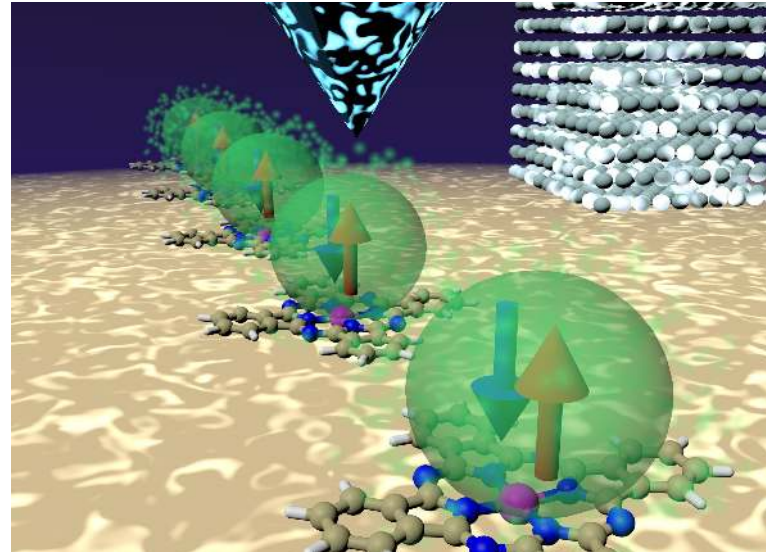
Hamiltonian learning quantum magnets with scanning tunnel microscopy

Jose Lado

Department of Applied Physics, Aalto University, Finland



Phys. Rev. Lett. 131, 086701 (2023)



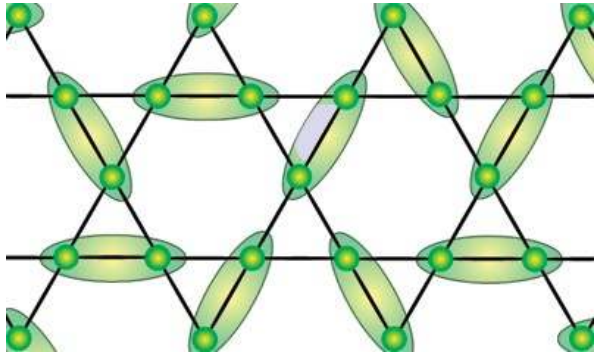
arXiv:2504.20711 (2025)

Low-Temperature Quantum Detectors, Helsinki

August 6th 2025

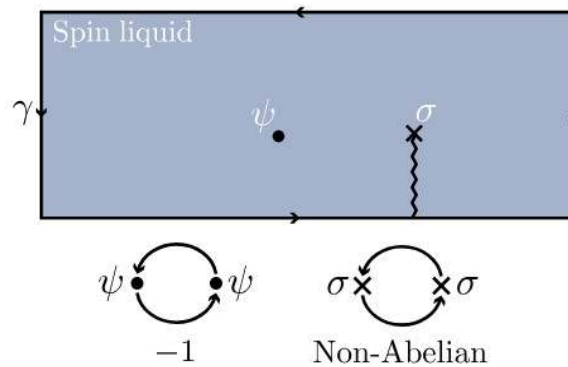
Quantum spin liquid quantum materials: opportunity and challenge

Macroscopically entangled magnetic state



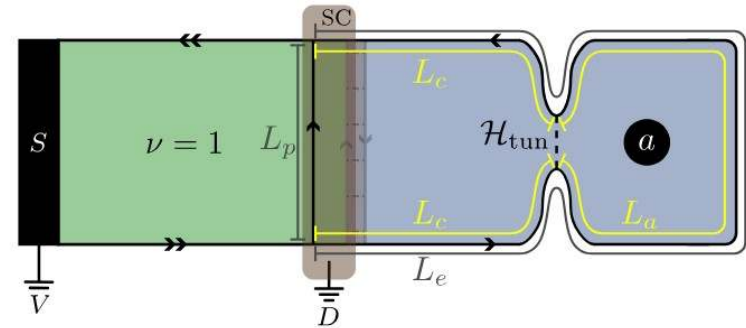
Rep. Prog. Phys. 80, 016502 (2017)

Featuring fractional excitations



Phys. Rev. X 10, 031014 (2020)

Potential candidates for topological quantum computing



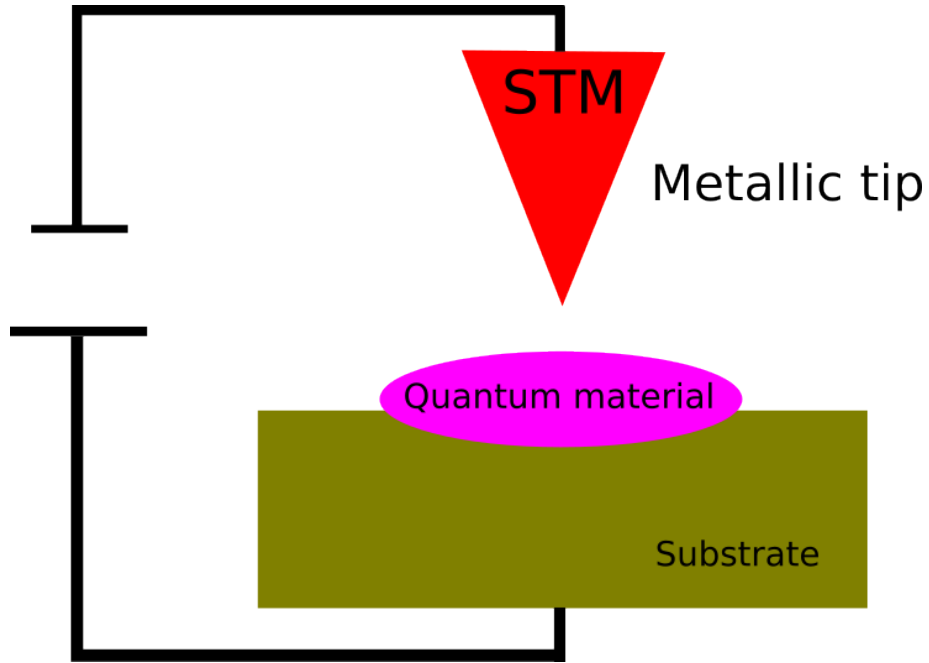
Phys. Rev. X 12, 011034 (2022)

Many potential QSL candidates (Herbertsmithite, 1T-TaS₂, RuCl₃)

However, quantum spin liquids are exceptionally challenging to identify experimentally

How can we learn the microscopic description of quantum magnets?

Sensing atomic-scale quantum matter with scanning tunneling microscopy

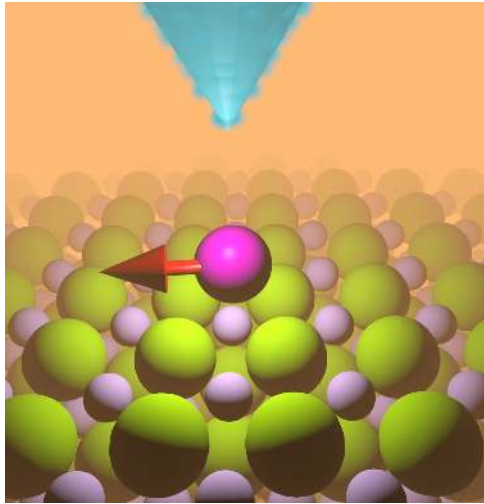


Atomic real-space resolution
Energy/frequency and time resolution

- Electronic states
- Phonons
- Superconducting pairing
- Orbital order
- Spin excitations

Probing quantum matter with scanning tunneling microscopy

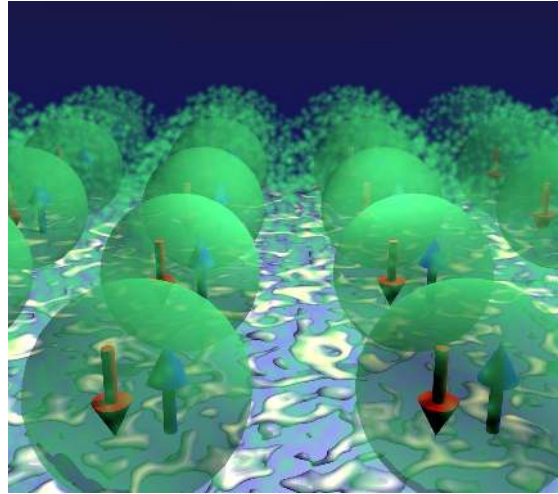
Atomic lattices



$a = 0.5 \text{ nm}$

Nature Materials 13, 782–785 (2014)
Rev. Mod. Phys. 91, 041001 (2019)

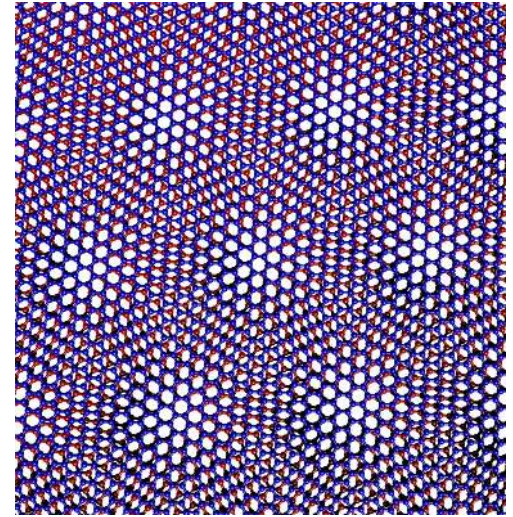
Molecular lattices



$a = 3 \text{ nm}$

Nature Rev. Mat. 5 (2), 87–104 (2020)
Nature 598, 287–292 (2021)

Twisted 2D moire materials

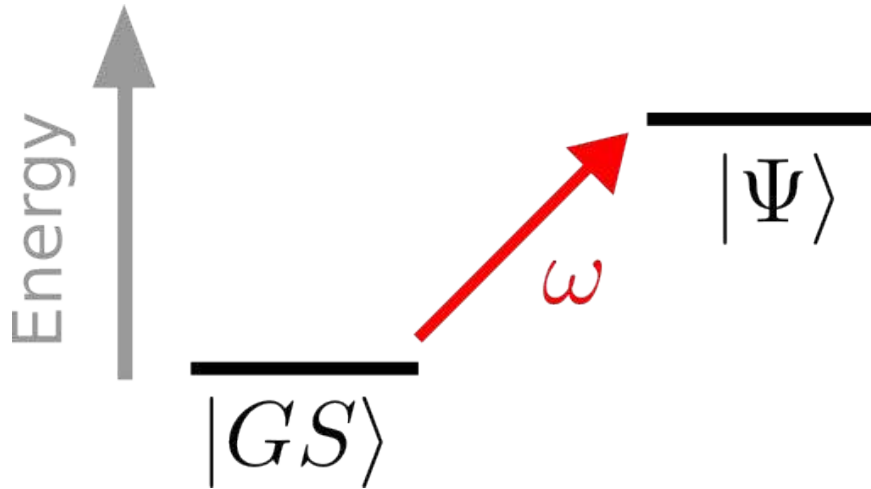


$a = 30 \text{ nm}$

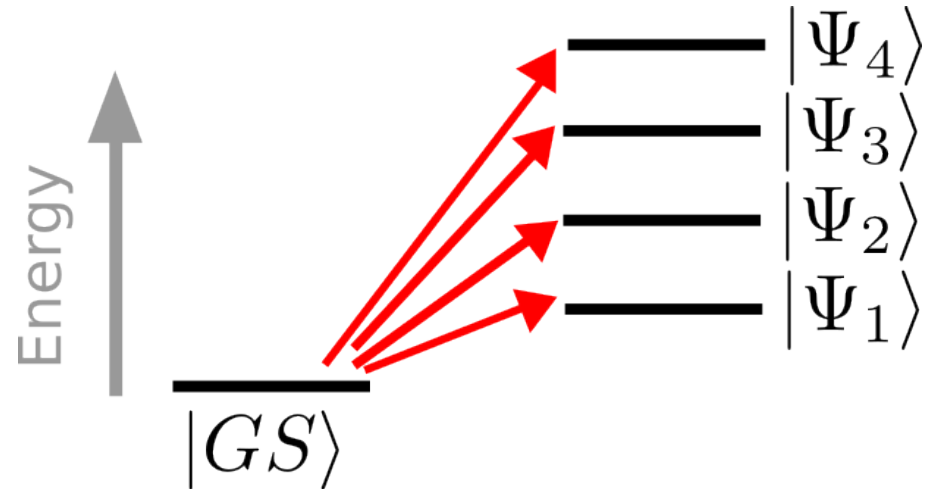
Nature 600, 240–245 (2021)
Nat. Rev. Materials 9, 460–480 (2024)

Spin excitations in collective quantum matter

Minimal quantum spin system



Complex quantum many-body system



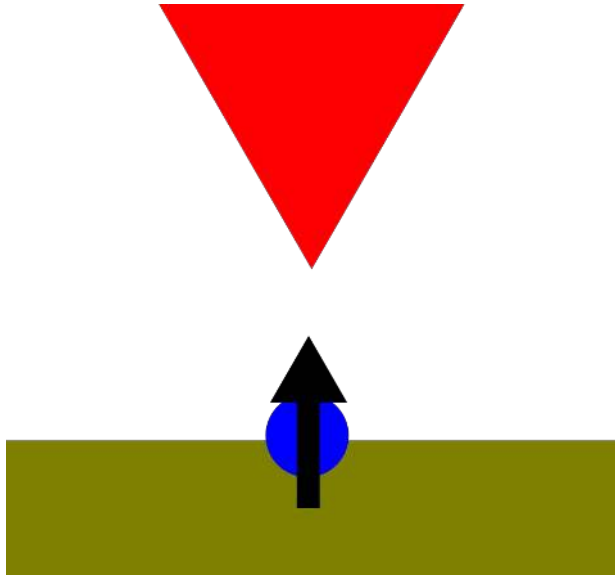
Spin excitations in quantum materials are described by the dynamical spin correlator

$$A(\mathbf{r}, \omega) = \sum_{\alpha} \langle GS | S_{\mathbf{r}}^{\alpha} \delta(\omega - H + E_{GS}) S_{\mathbf{r}}^{\alpha} | GS \rangle$$

That can be measured with scanning tunneling microscopy

Two ways of measuring spin-excitations in atomic-scale magnets

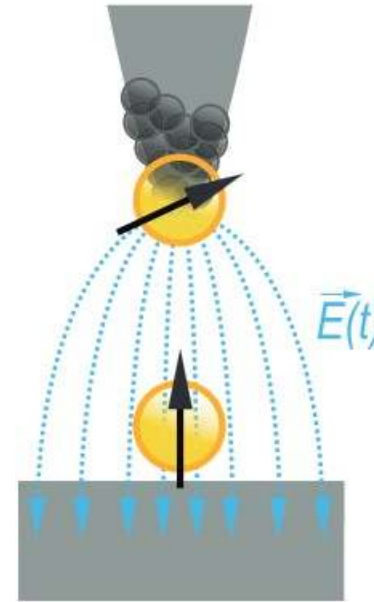
Inelastic electron tunnelling
Spectroscopy (IETS)



Science 306 (5695), 466-469 (2004)

Without spin-polarized tip

Electrically driven electron
spin resonance (ESR)

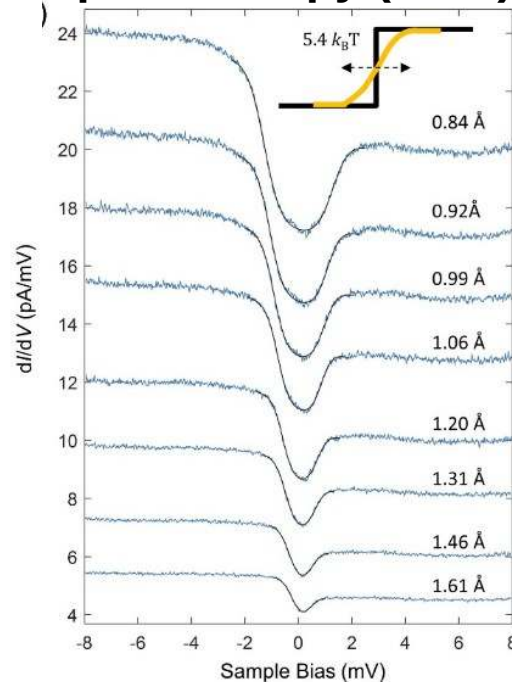
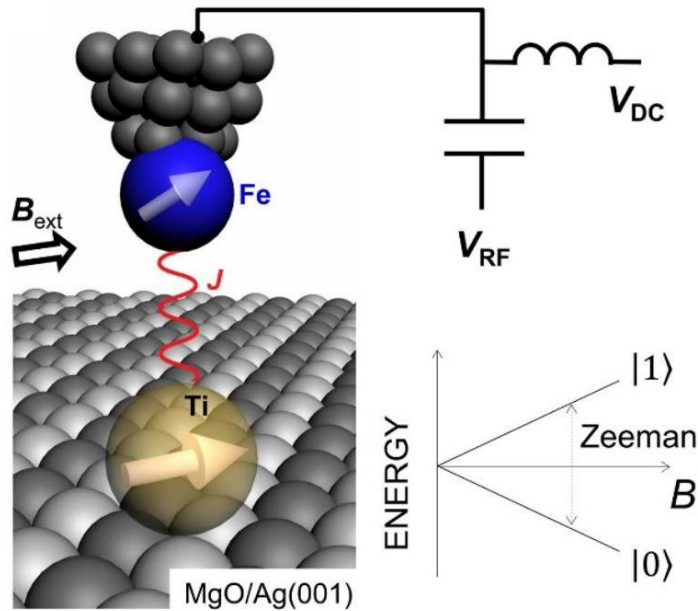


Science 350 (6259), 417-420 (2015)

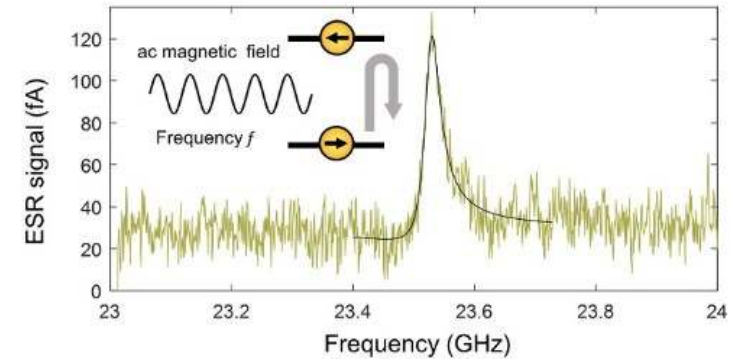
With spin-polarized tip

Measuring spin excitations with IETS and ESR

Inelastic electron tunneling spectroscopy (IETS)



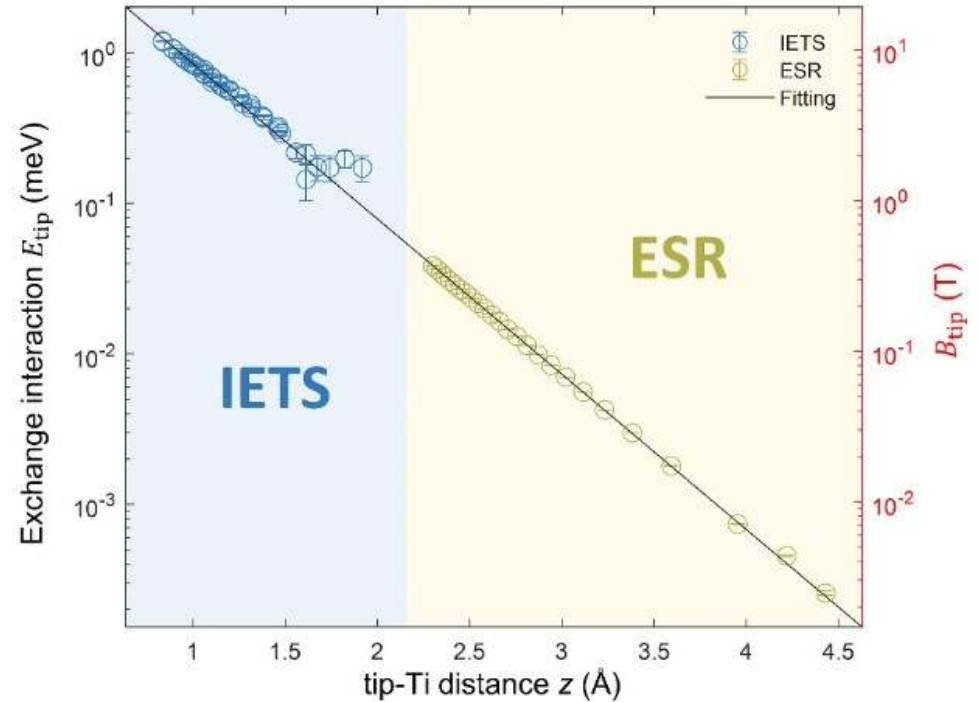
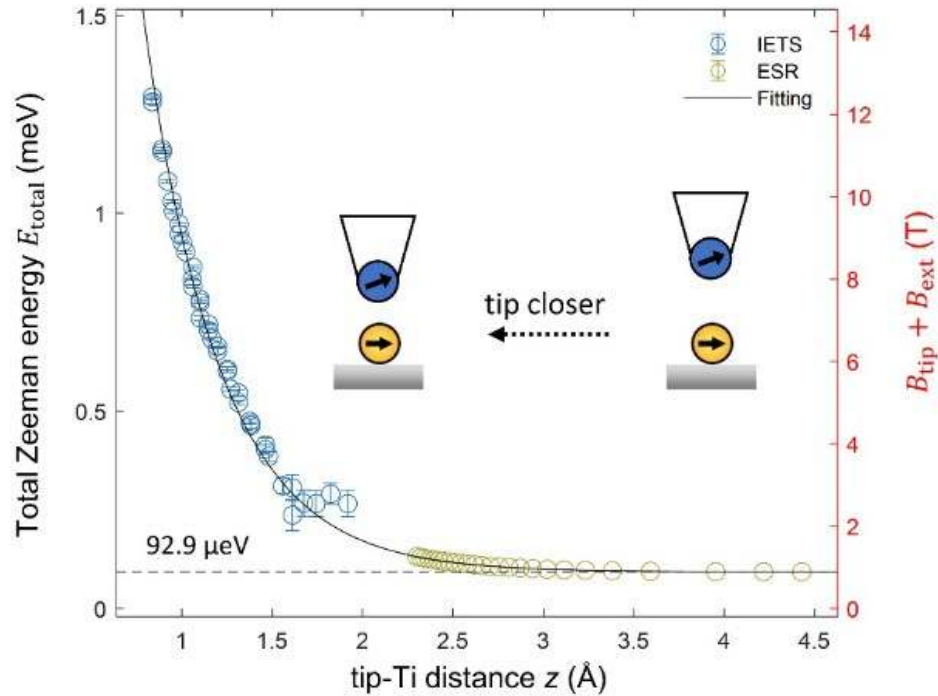
Electron spin resonance (ESR) with STM



Spin excitations can be probed with IETS or ESR

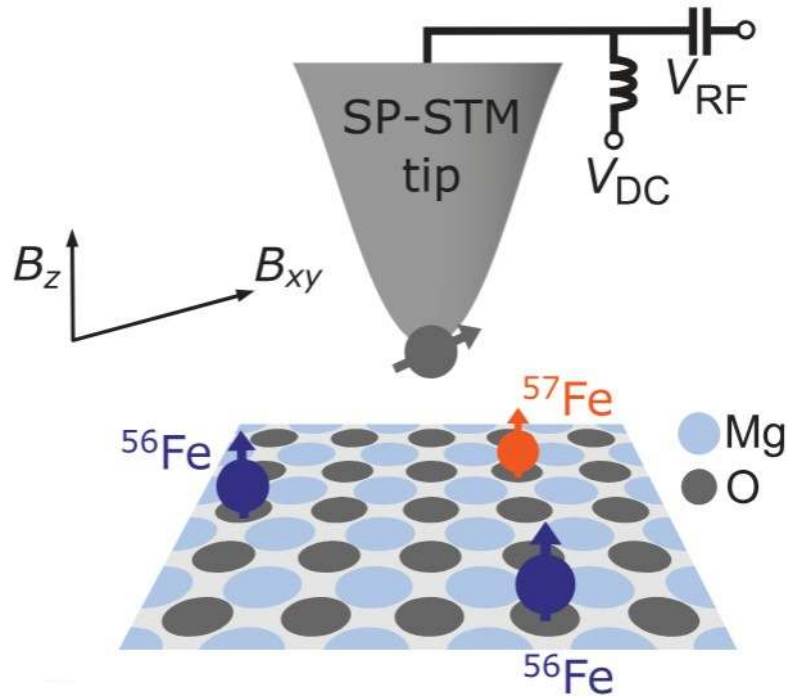
Measuring spin excitations with IETS and ESR

Excitation energies ranging 3 orders of magnitude can be measured with STM

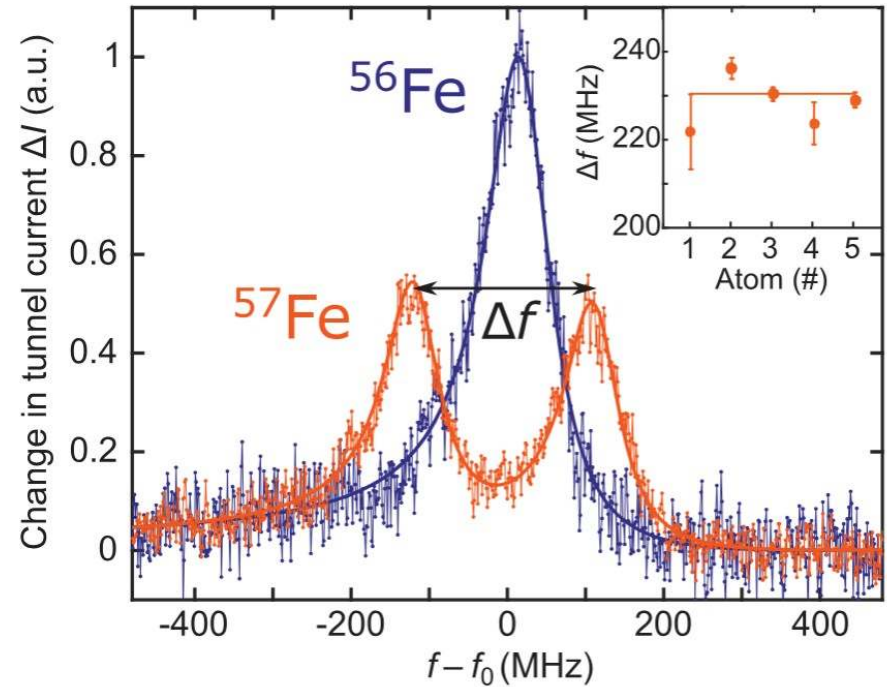


Hyperfine coupling in individual atoms with STM

Individually addressing each isotope

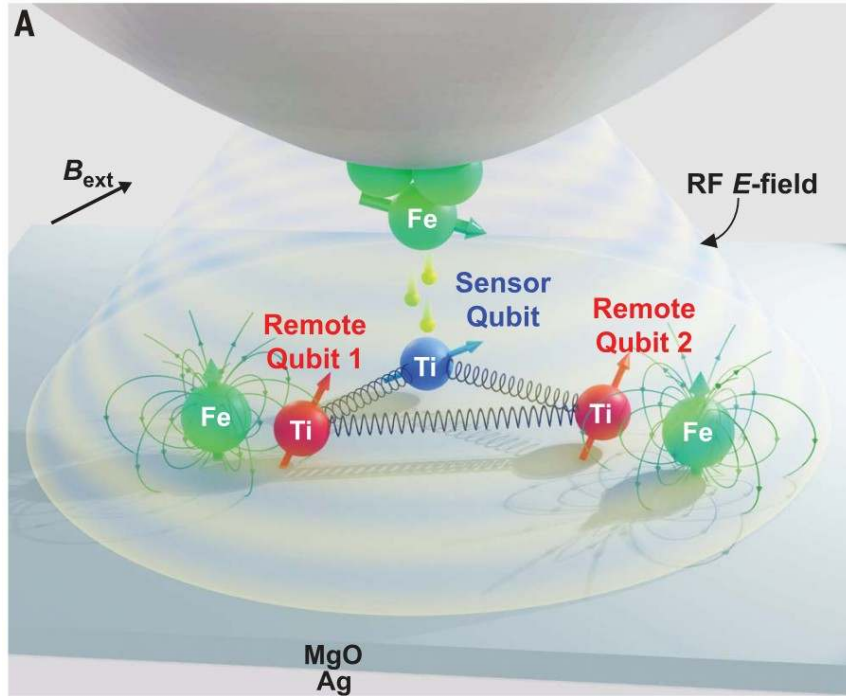


Hyperfine splitting in ESR

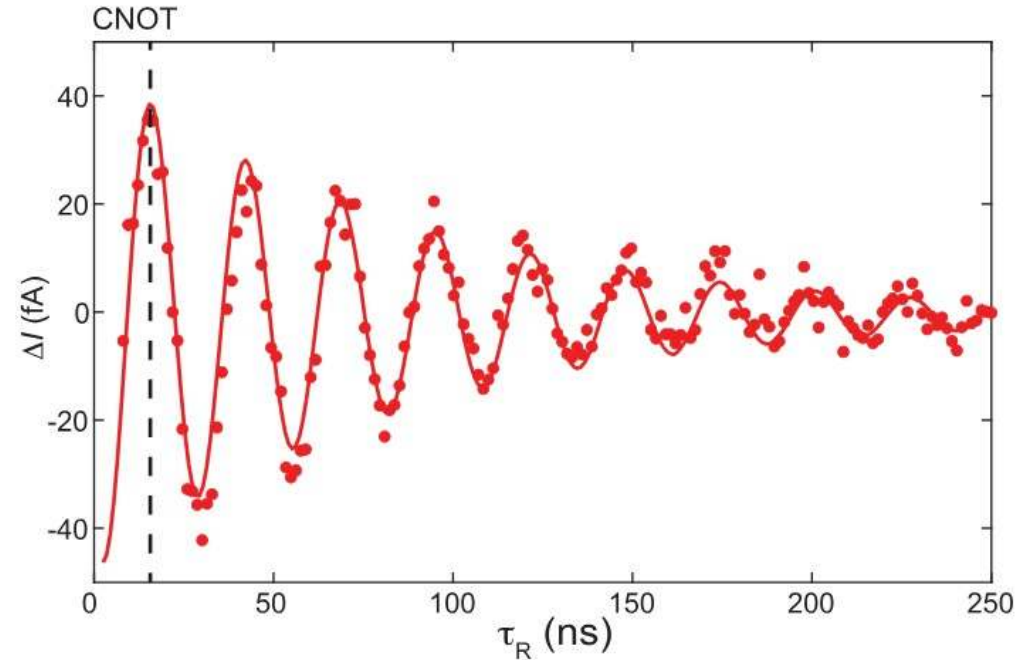


Qubit operations between individual atoms with STM

Atomic multiqubit platform

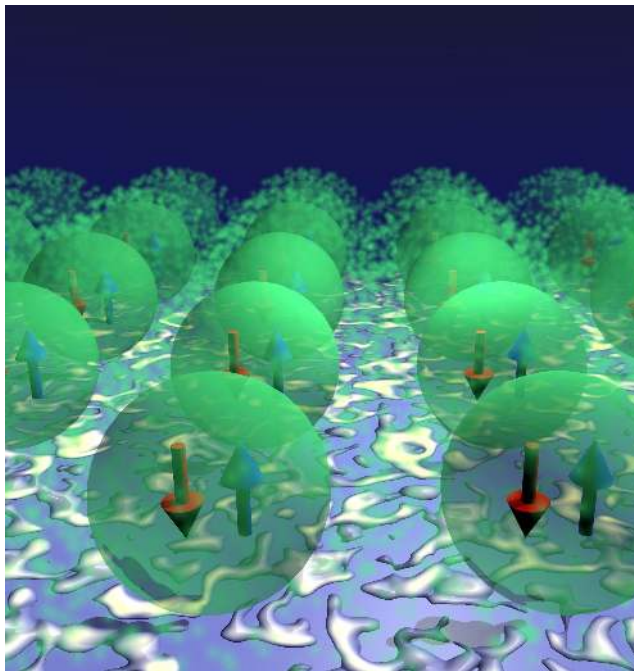


Conditional controlled qubit rotation



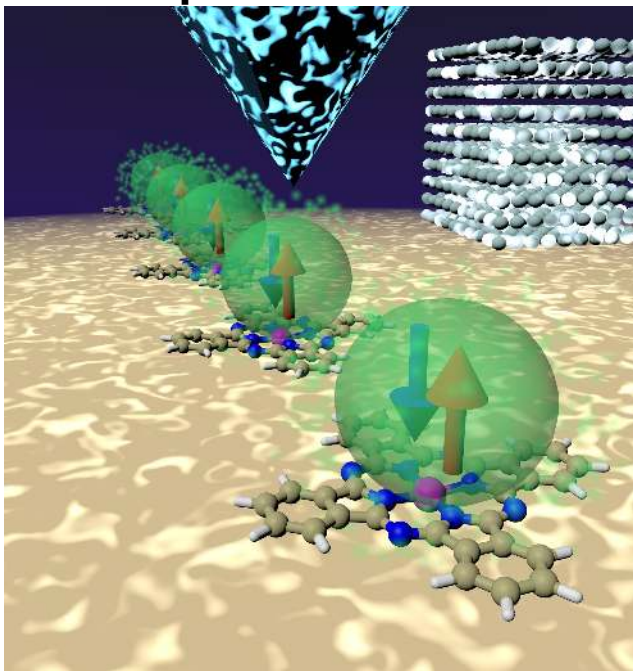
Engineering and Hamiltonian learning of an artificial quantum magnet

A triplon molecular quantum magnet



Phys. Rev. Lett. 131, 086701 (2023)

Hamiltonian learning from spin excitation



arXiv:2504.20711 (2025)

Robert Drost



Rouven Koch



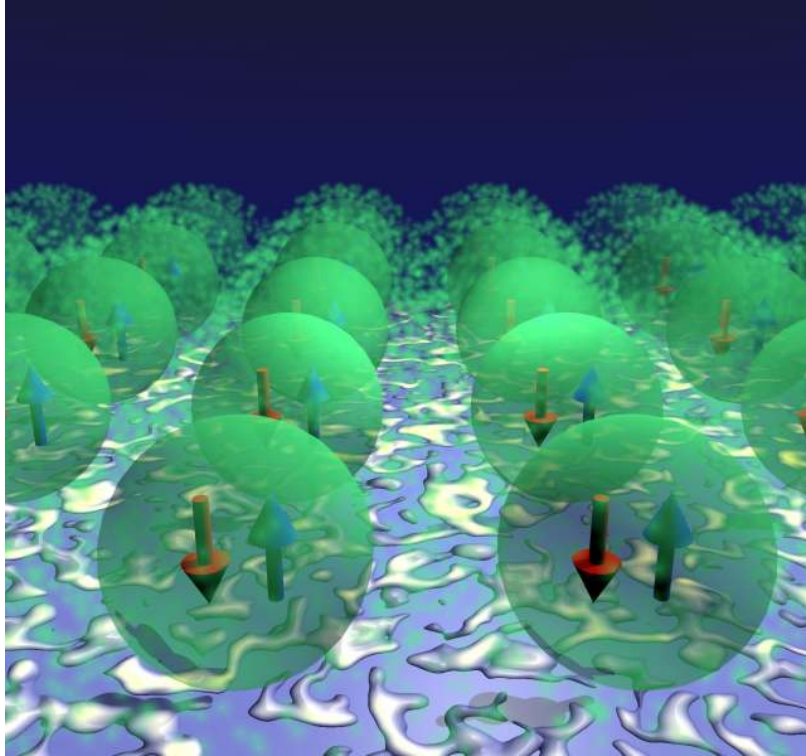
Shawulienu
Kezilebieke



Peter Liljeroth



Triplons in an artificial molecular quantum magnet



Phys. Rev. Lett. 131, 086701 (2023)

Triplons in an atomic-scale quantum magnet

Robert Drost



Shawulienu Kezilebieke



Peter Liljeroth



The quantum Heisenberg dimer

Let us take a minimal quantum magnet

$$\mathcal{H} = \vec{S}_0 \cdot \vec{S}_1$$

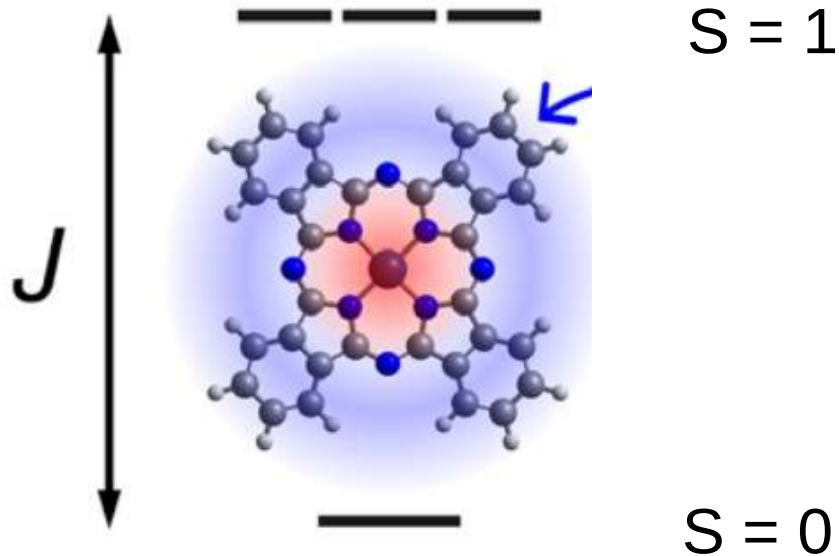
The ground state is unique, and does not break time-reversal

$$|GS\rangle = \frac{1}{\sqrt{2}} [|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle] \quad \langle \vec{S}_i \rangle = 0$$

The state is maximally entangled

How does this generalize to a macroscopic system, and how can we see its quasiparticles?

Triplet excitations in a molecule

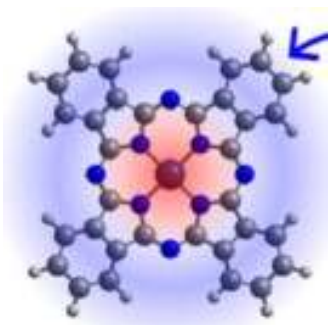


Internal antiferromagnetic coupling

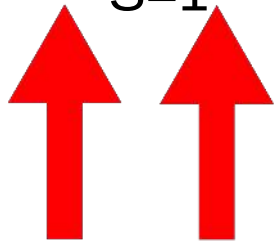
$$H = JS \cdot \mathbf{K}$$

A single molecule (CoPC) has a singlet-triplet transition

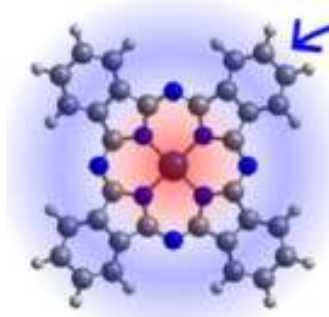
Triplons in a dimer



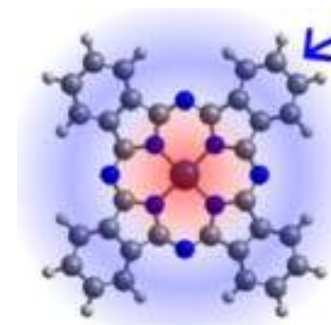
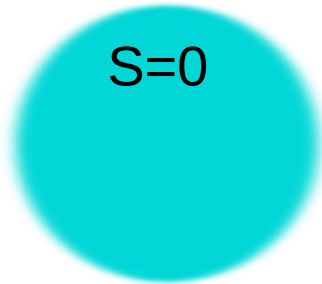
$S=1$



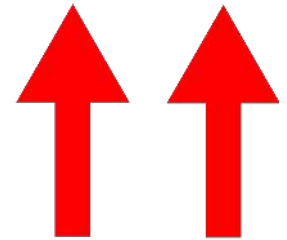
A triplet in the first molecule



$S=0$



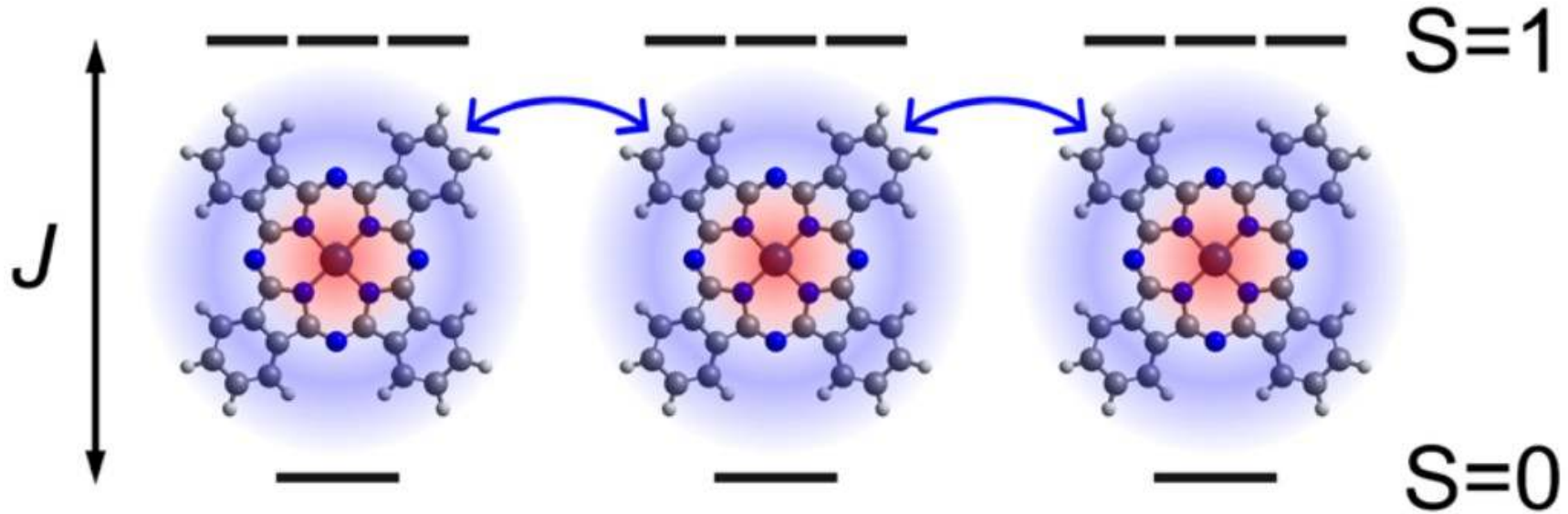
$S=1$



Can jump to the second molecule

Exchange coupling between molecules leads to propagating triplon excitations

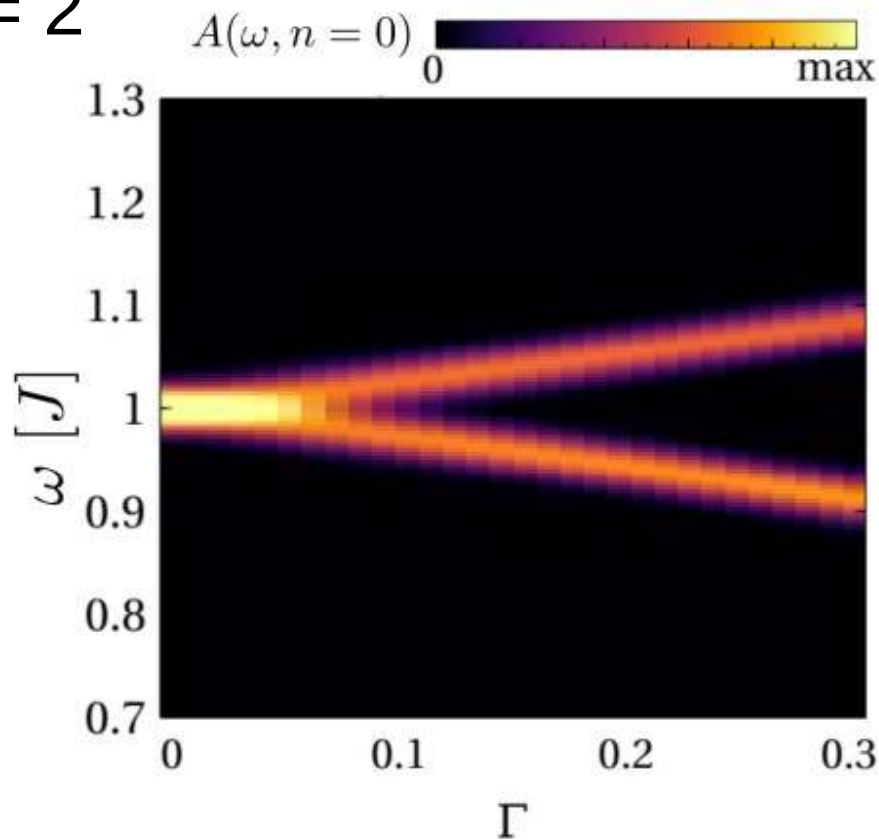
A chain of triplons



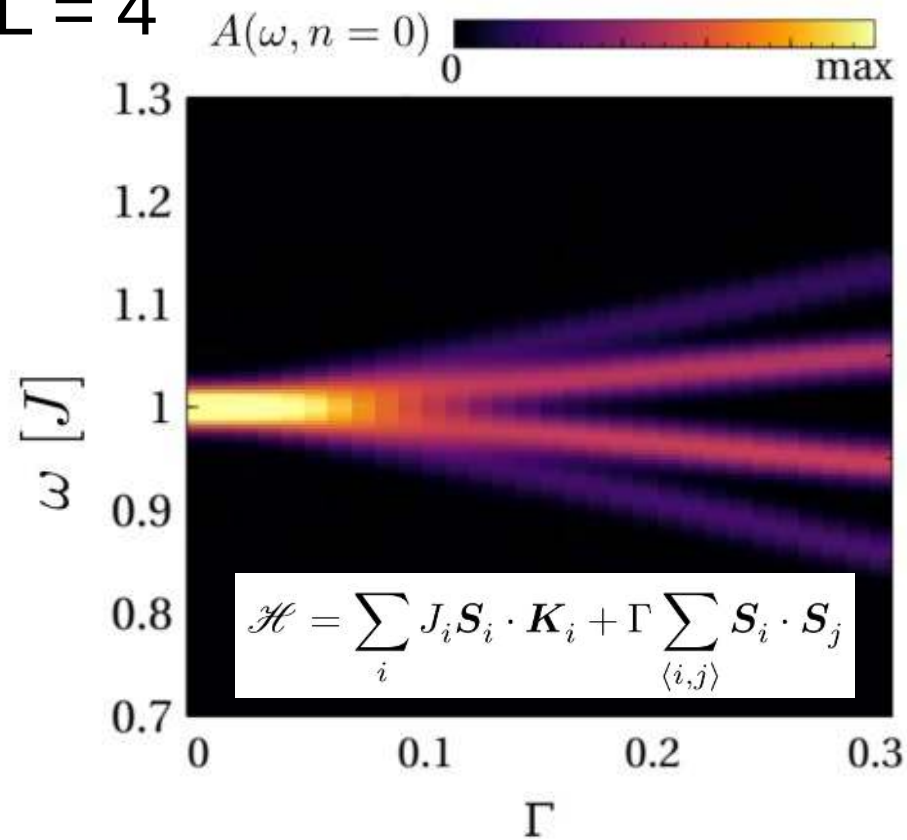
$$\mathcal{H} = \sum_i J_i \mathbf{S}_i \cdot \mathbf{K}_i + \Gamma \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

Chains of triplons

$L = 2$

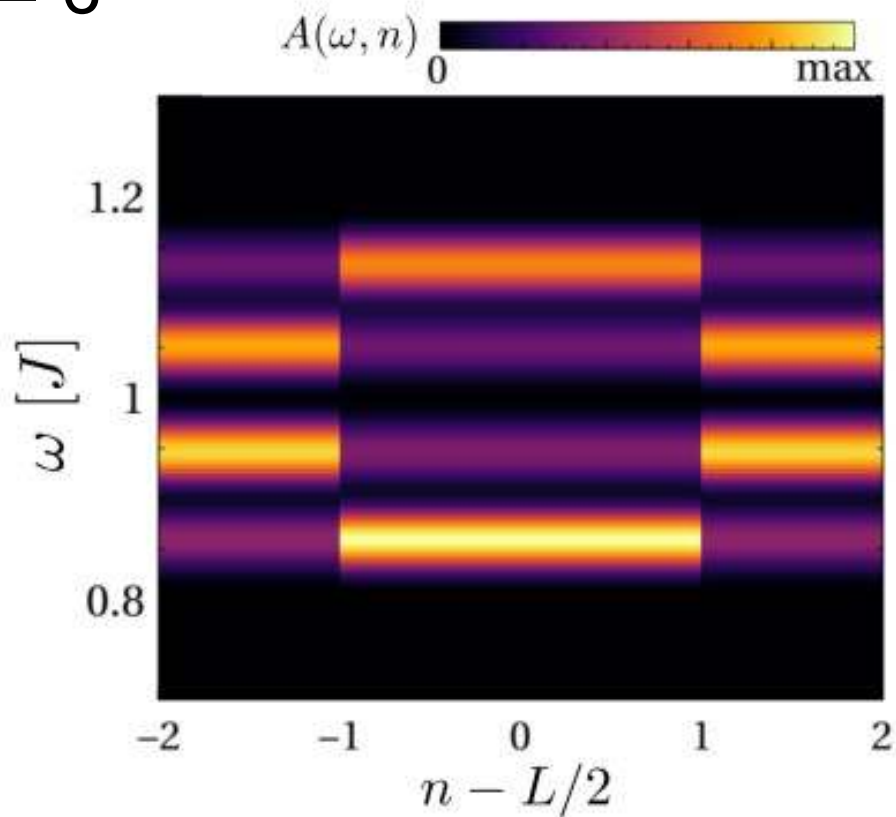


$L = 4$

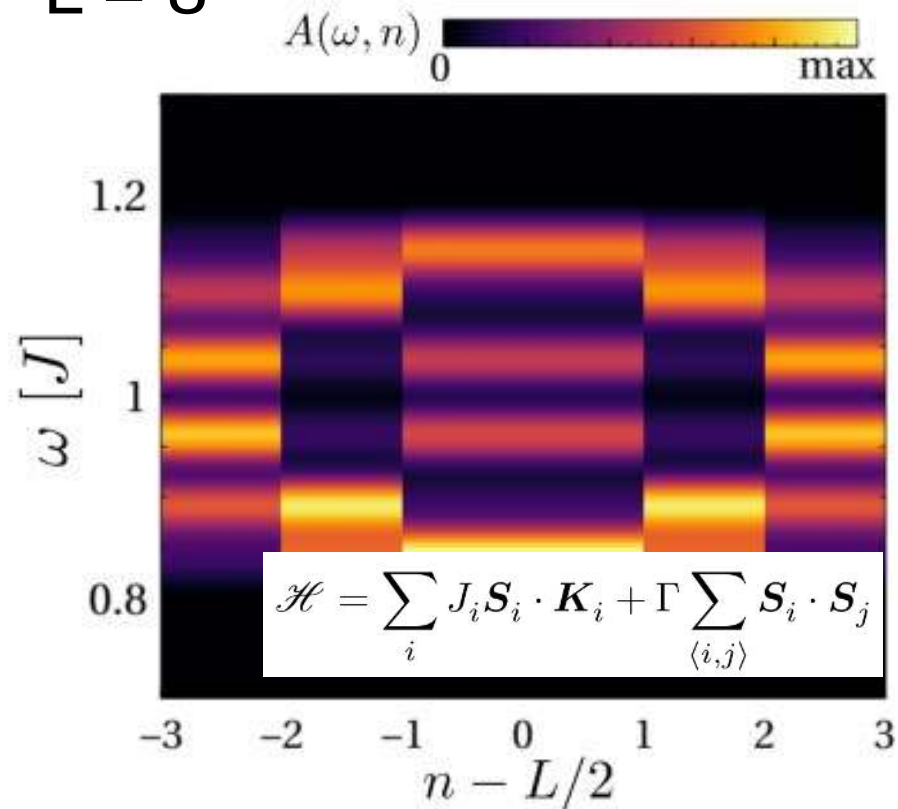


Chains of triplons

$L = 6$

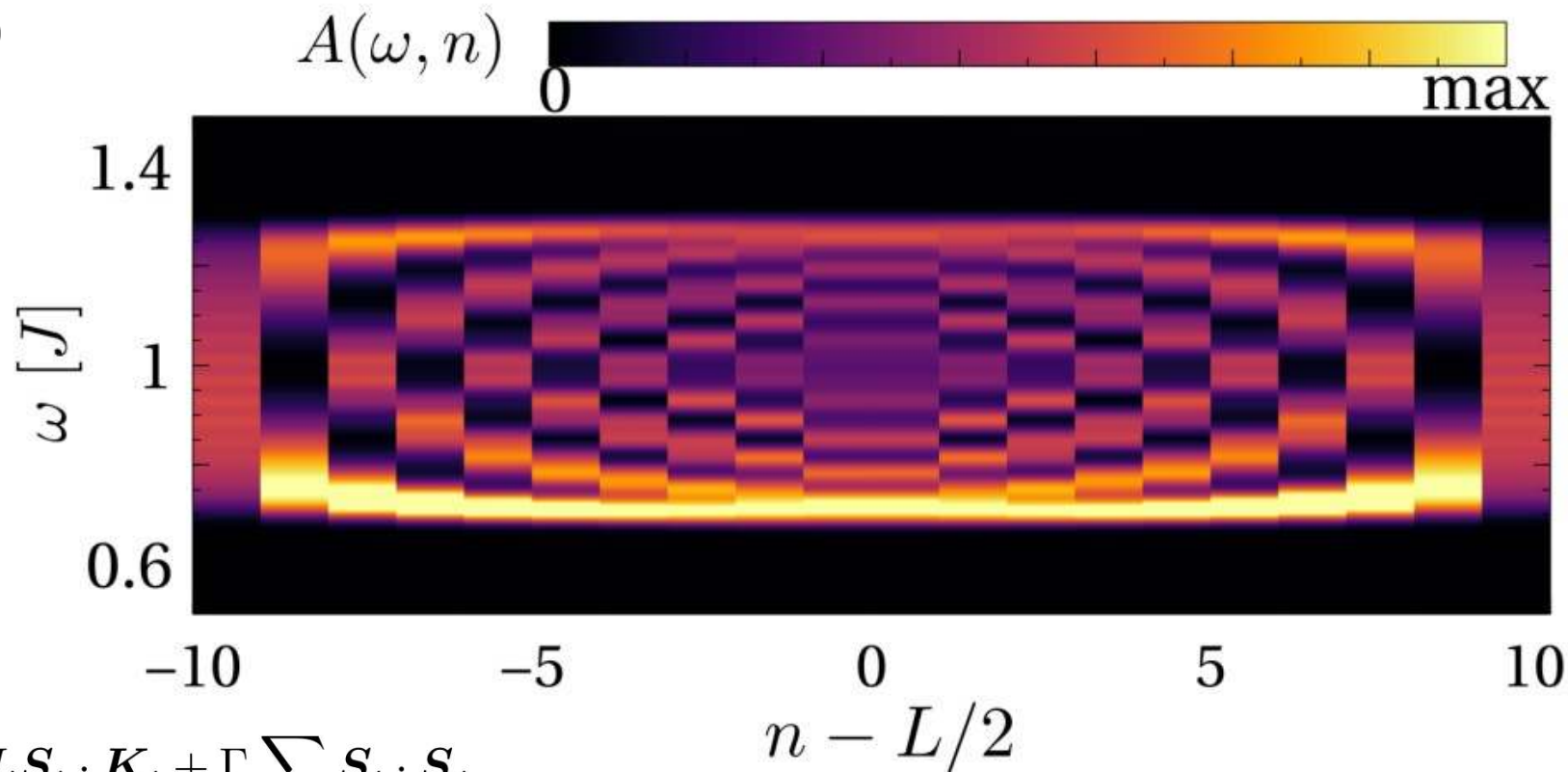


$L = 8$



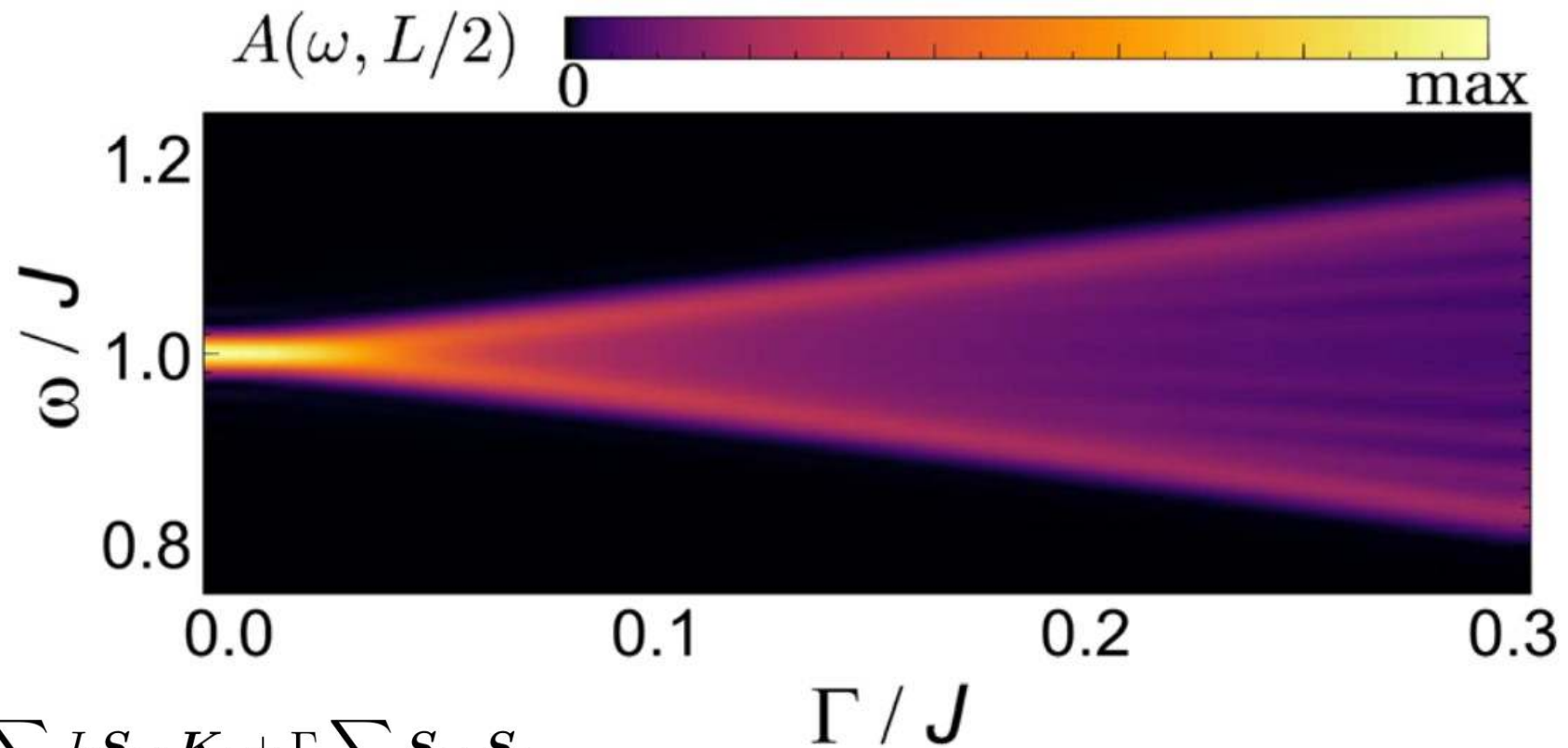
Chains of triplons

$L = 20$



$$\mathcal{H} = \sum_i J_i \mathbf{S}_i \cdot \mathbf{K}_i + \Gamma \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

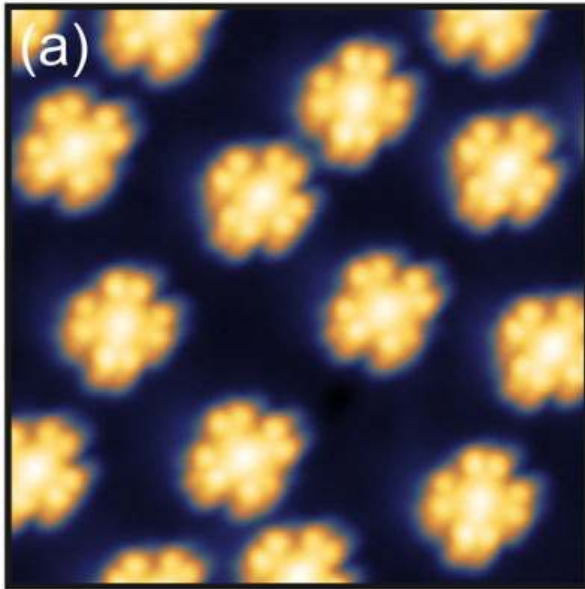
Chains of triplons



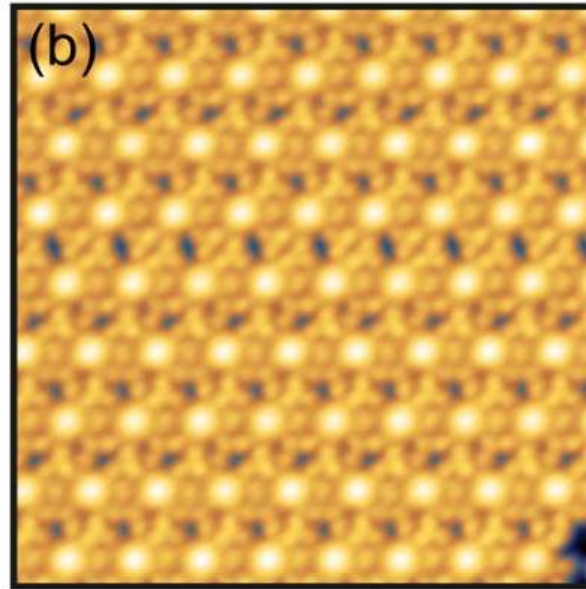
$$\mathcal{H} = \sum_i J_i \mathbf{S}_i \cdot \mathbf{K}_i + \Gamma \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

0D, 1D and 2D arrays

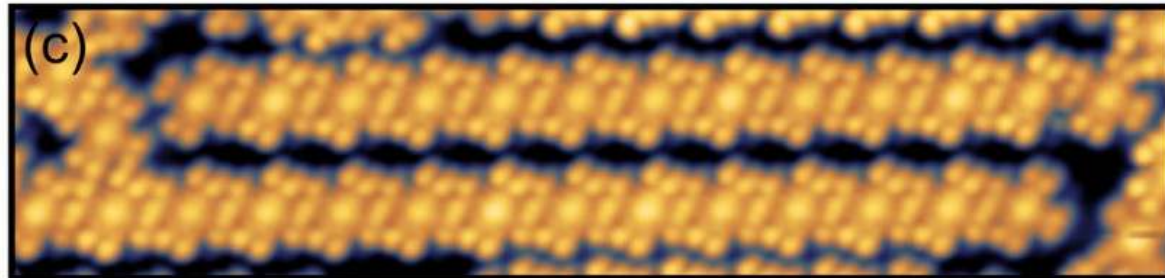
0D



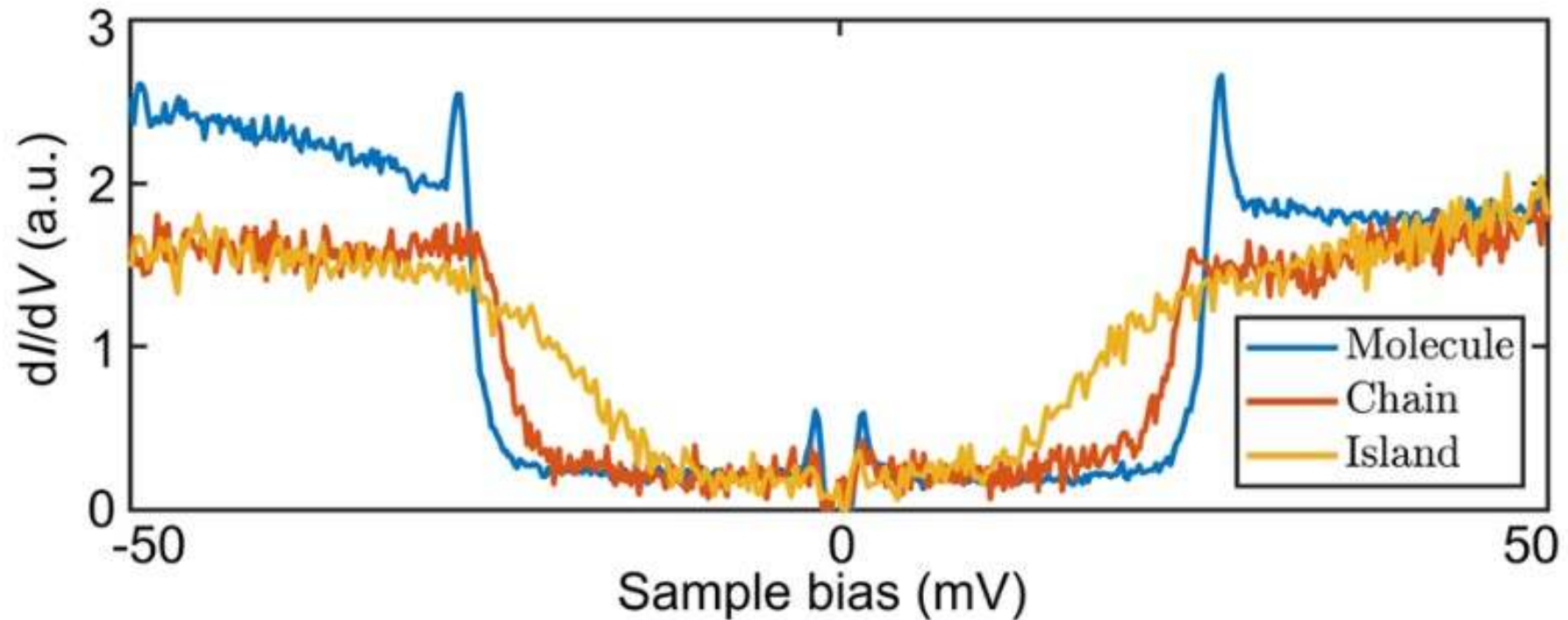
2D



1D

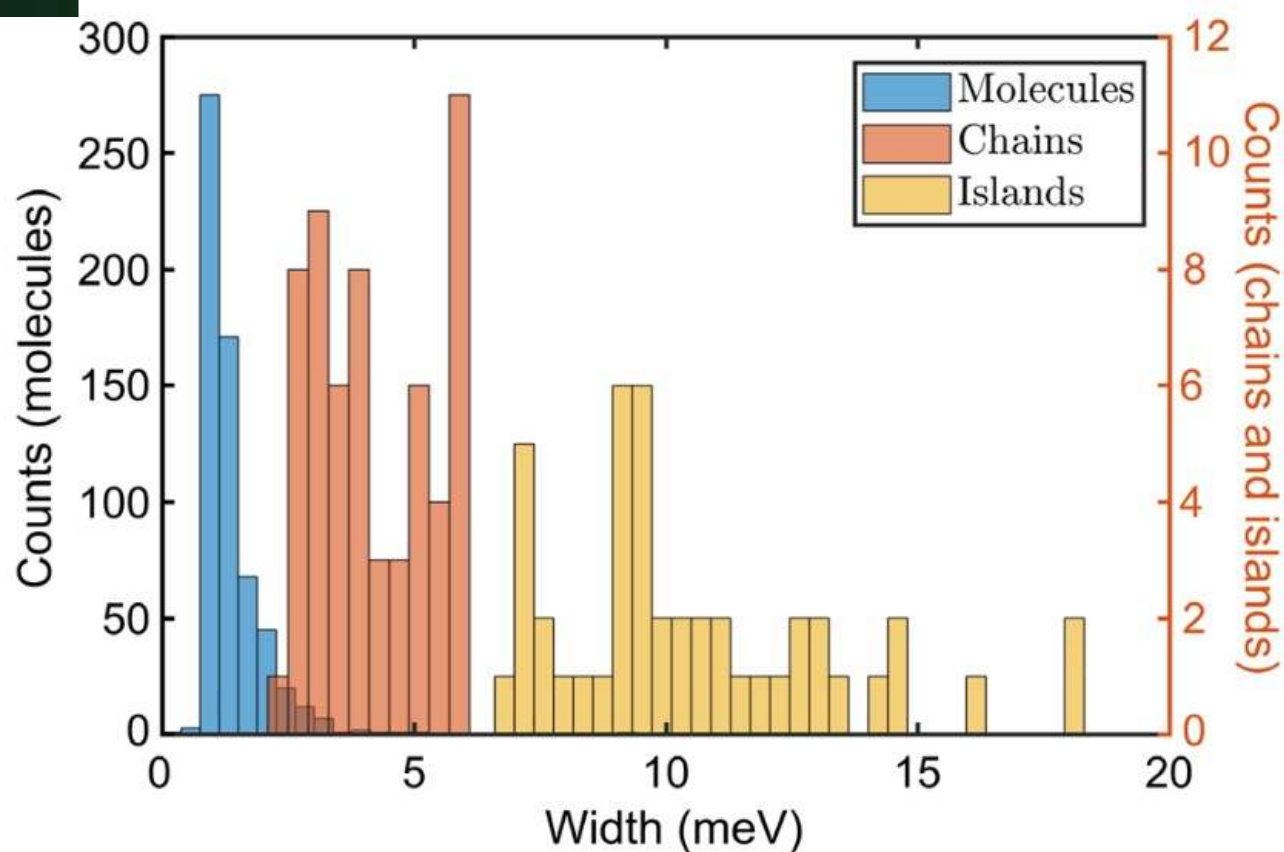


Inelastic excitations for 0D, 1D and 2D arrays



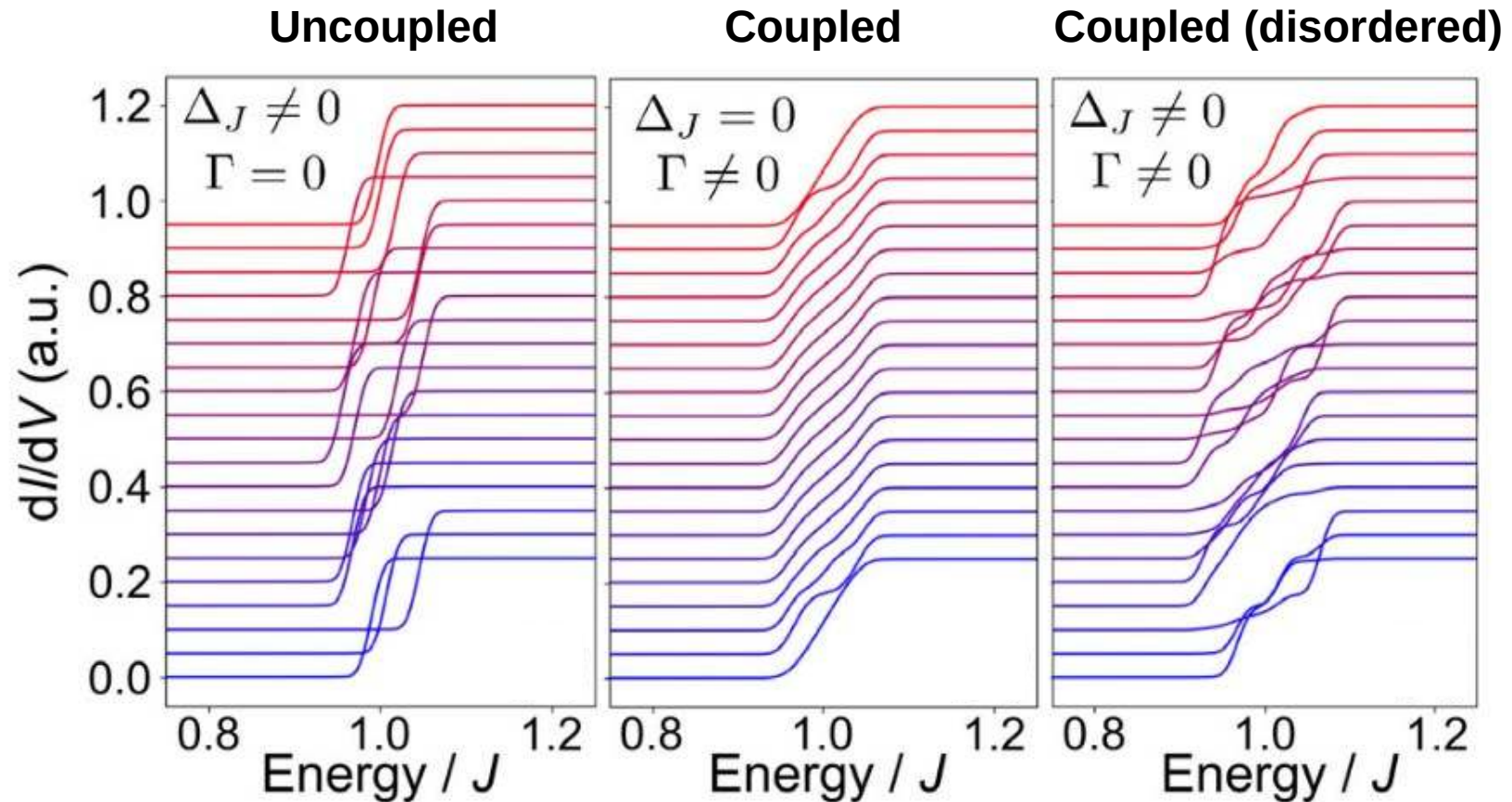
The width of the step is bigger for islands than for molecules

Statistic of the inelastic steps

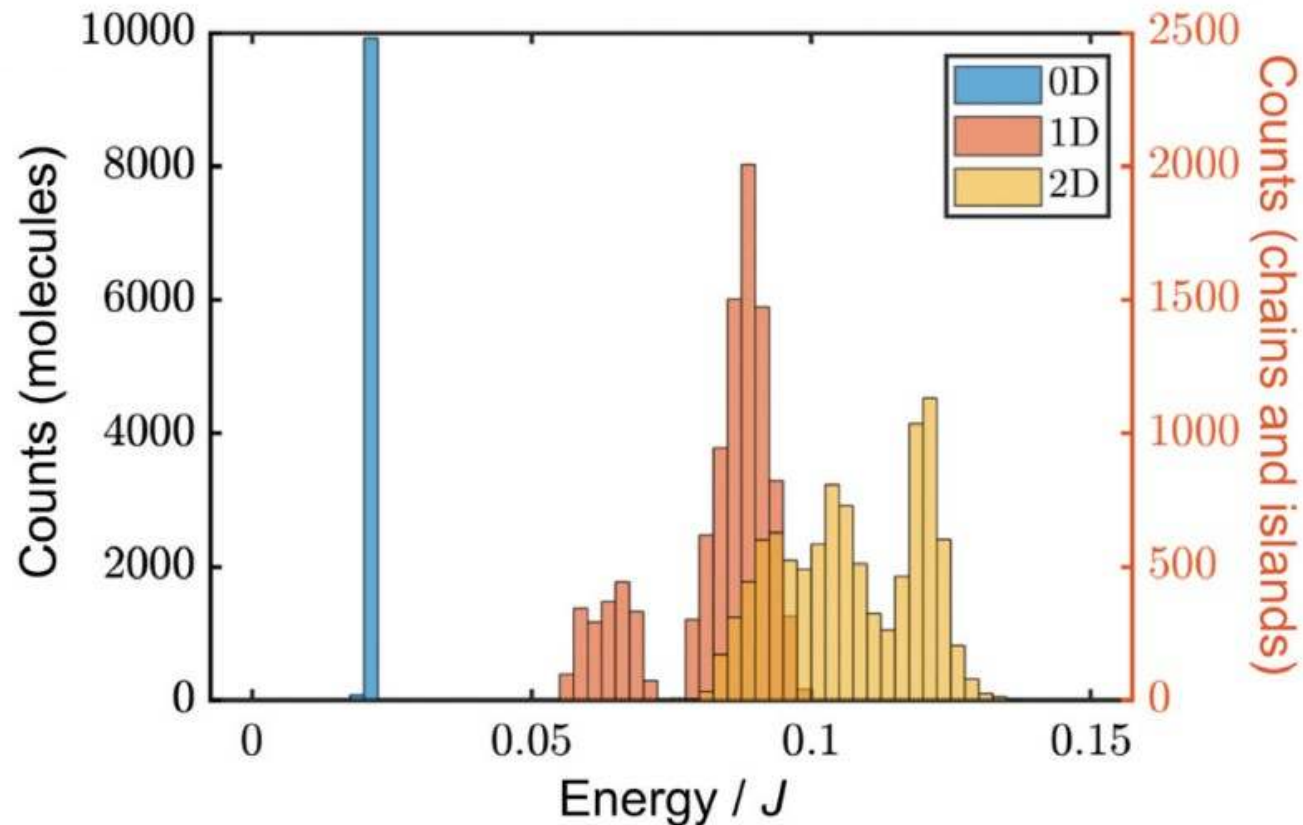


2D arrays have the largest width of the spin excitation

Coupled and uncoupled inelastic excitation

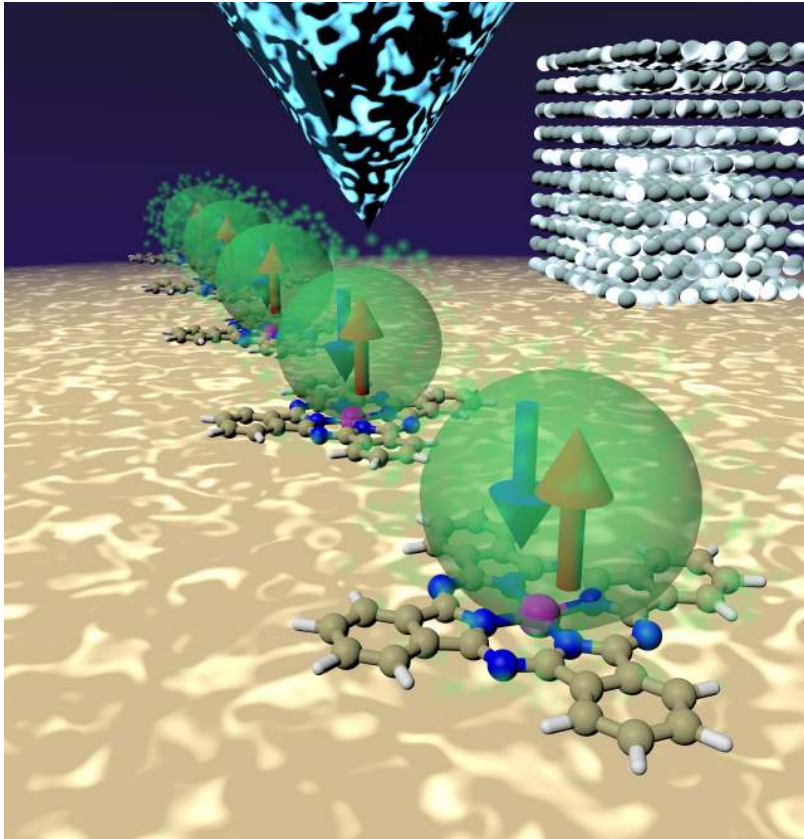


Statistics of the theoretical model



2D arrays have the largest width of the spin excitation

Hamiltonian learning triplons in a molecular quantum magnet



Rouven Koch



Robert Drost



Peter Liljeroth



arXiv:2504.20711 (2025)

The problem of Hamiltonian learning

If we measure a quantum magnet, how can we obtain its microscopic description?

Conventional workflow in theoretical physics

Solve the Hamiltonian

H

Compute an observable

$\langle \Psi | O | \Psi \rangle$

The problem we often find experimentally

Measure an observable

$\langle \Psi | O | \Psi \rangle$

Determine its Hamiltonian

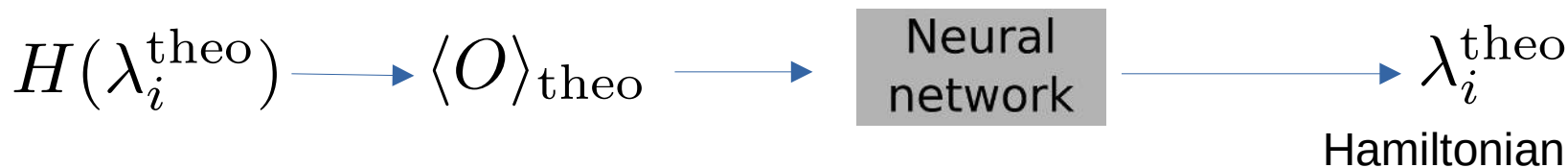
H

How can we extract a complex Hamiltonian from its physical observables?

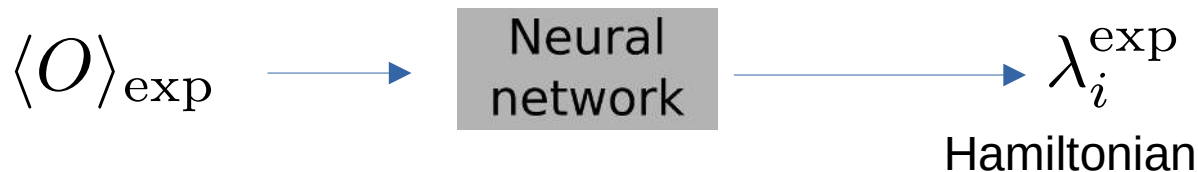
Hamiltonian learning: combining theory and experiment

Train a machine learning algorithm with simulated quantum many-body theory data

so that the parameters of each Hamiltonian are explicitly known

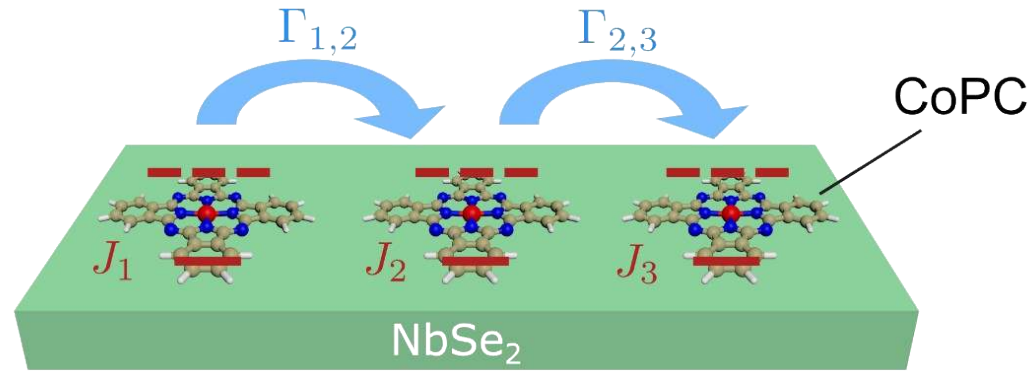


With the trained algorithm, evaluated in experimental data to extract the parameters



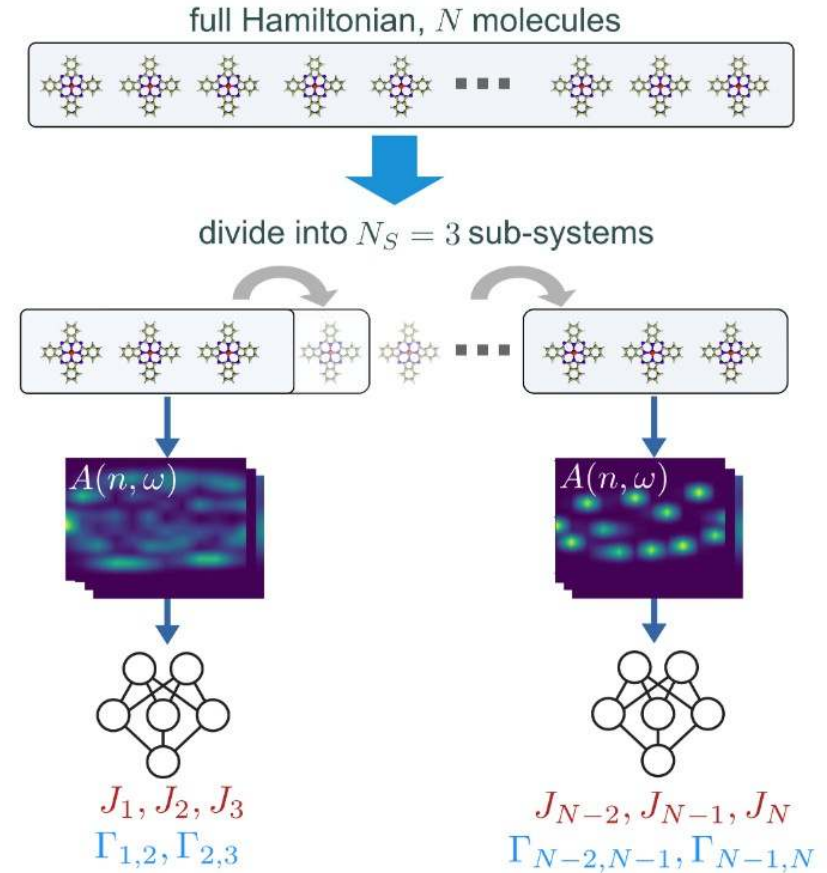
This strategy combines quantum many-body solvers, machine learning, and experimental measurements

A strategy for machine learning quantum magnets



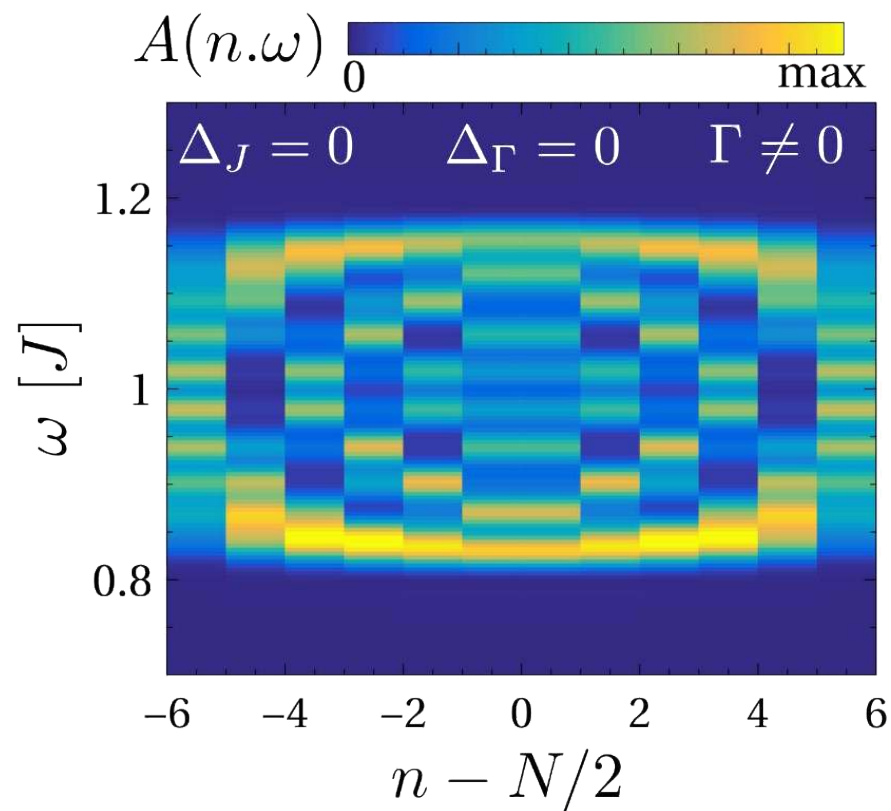
$\Gamma_{n,n+1}$: intermolecular exchange
 J_n : intramolecular exchange

$$H = \sum_n J_n \mathbf{S}_n \cdot \mathbf{K}_n + \sum_n \Gamma_{n,n+1} \mathbf{S}_n \cdot \mathbf{S}_{n+1}$$

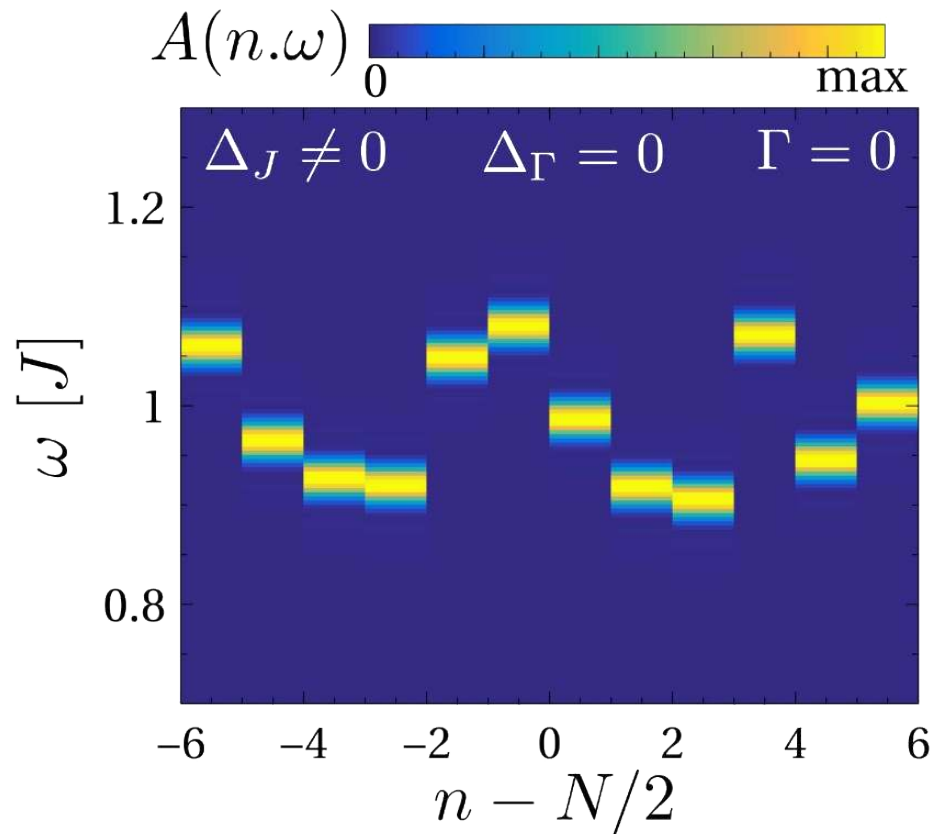


Pristine triplon excitations

Pristine and coupled molecules

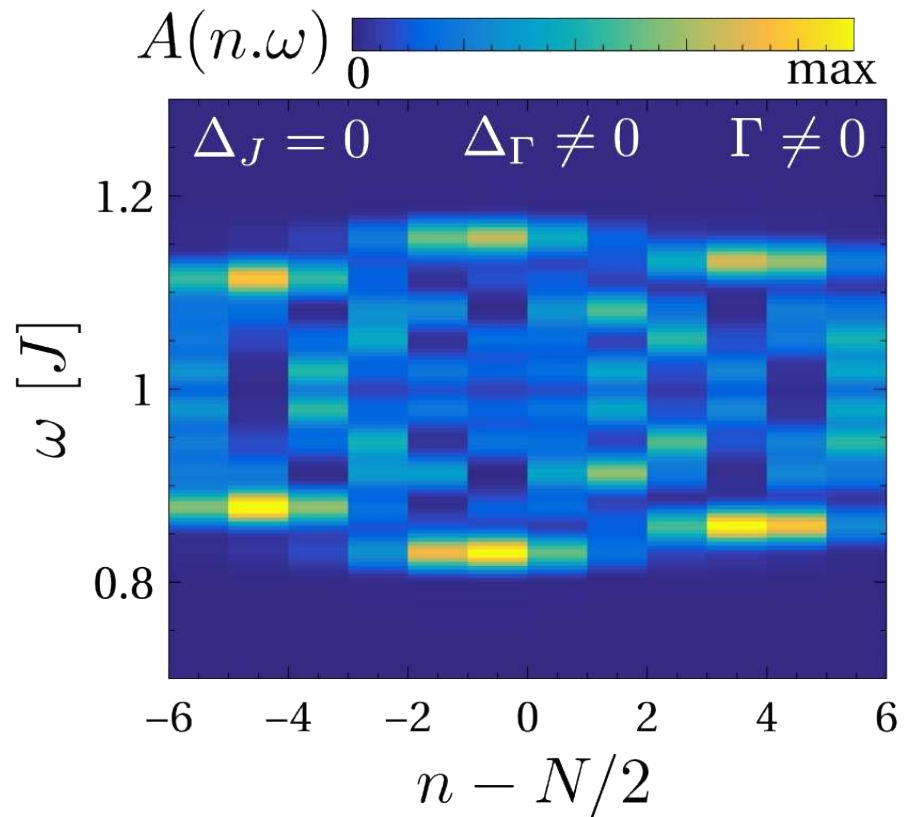


Decoupled and disordered molecules

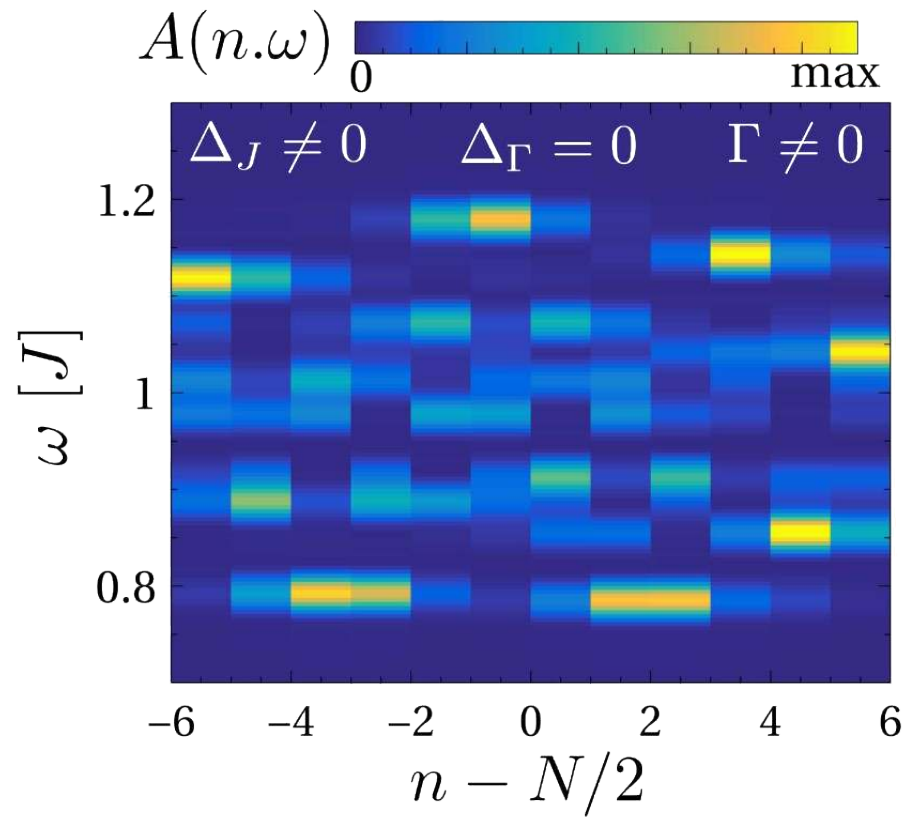


The different impact of molecular disorder

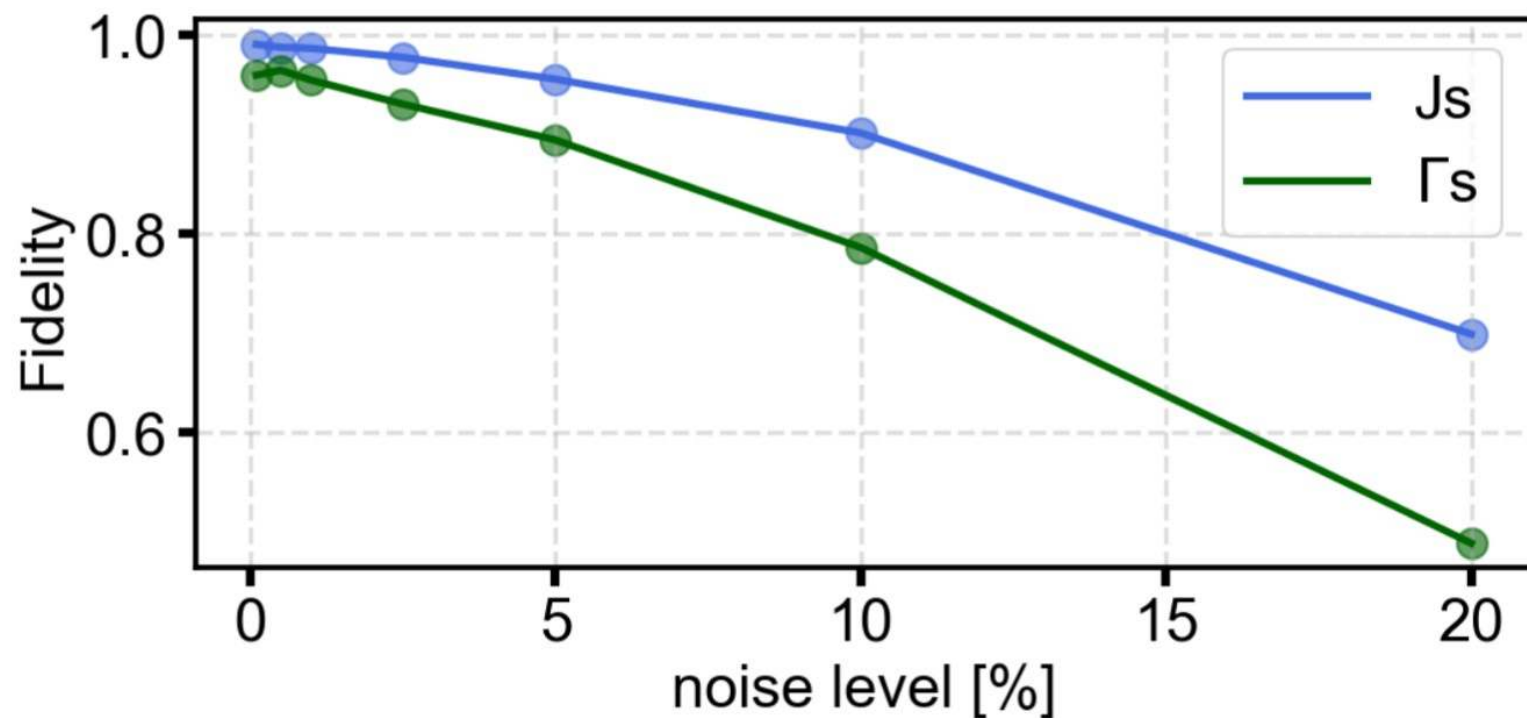
Disorder in inter-molecule coupling



Disorder in intra-molecule coupling

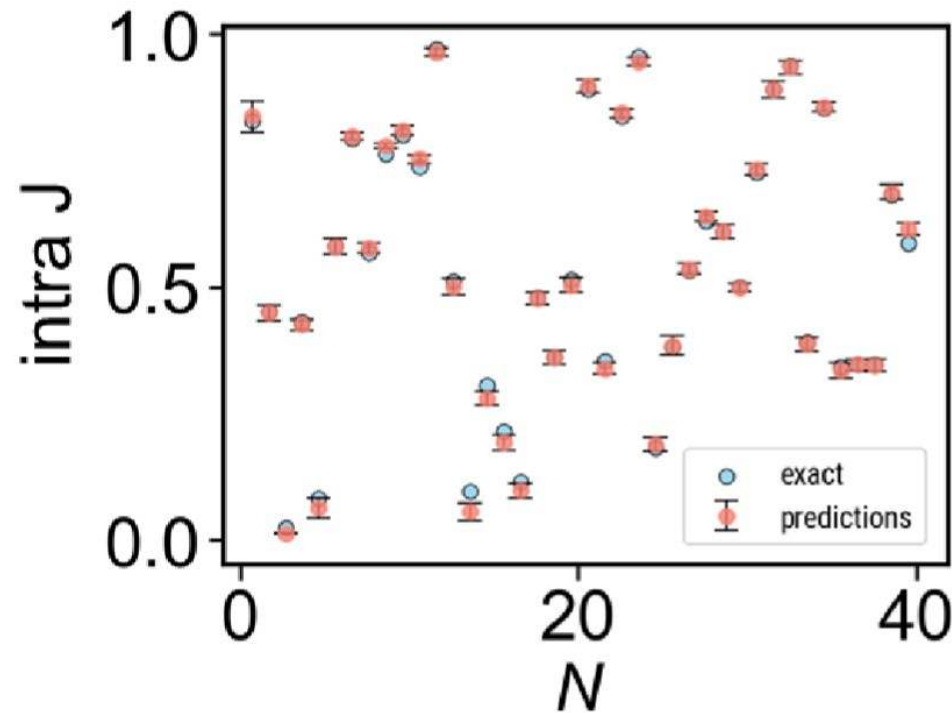


Resilience to noise of Hamiltonian learning

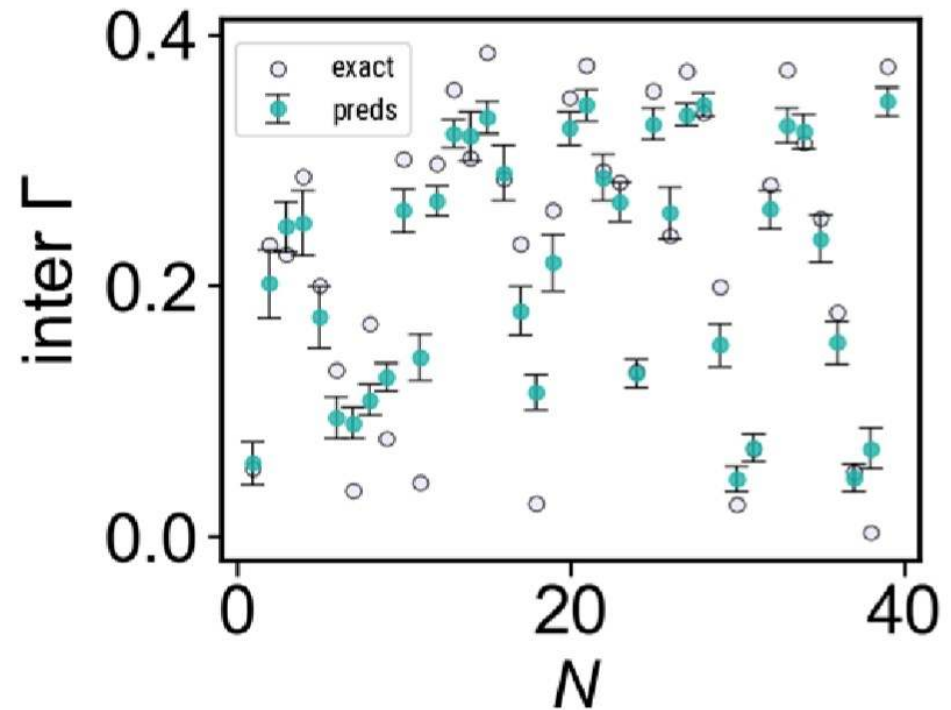


Noise model $\left(\frac{dI}{dV}\right)_{\text{noise}} = \left(\frac{dI}{dV}\right)_{\text{data}} + \eta \cdot \mathbf{R}$

Inferring spatially dependent exchange couplings

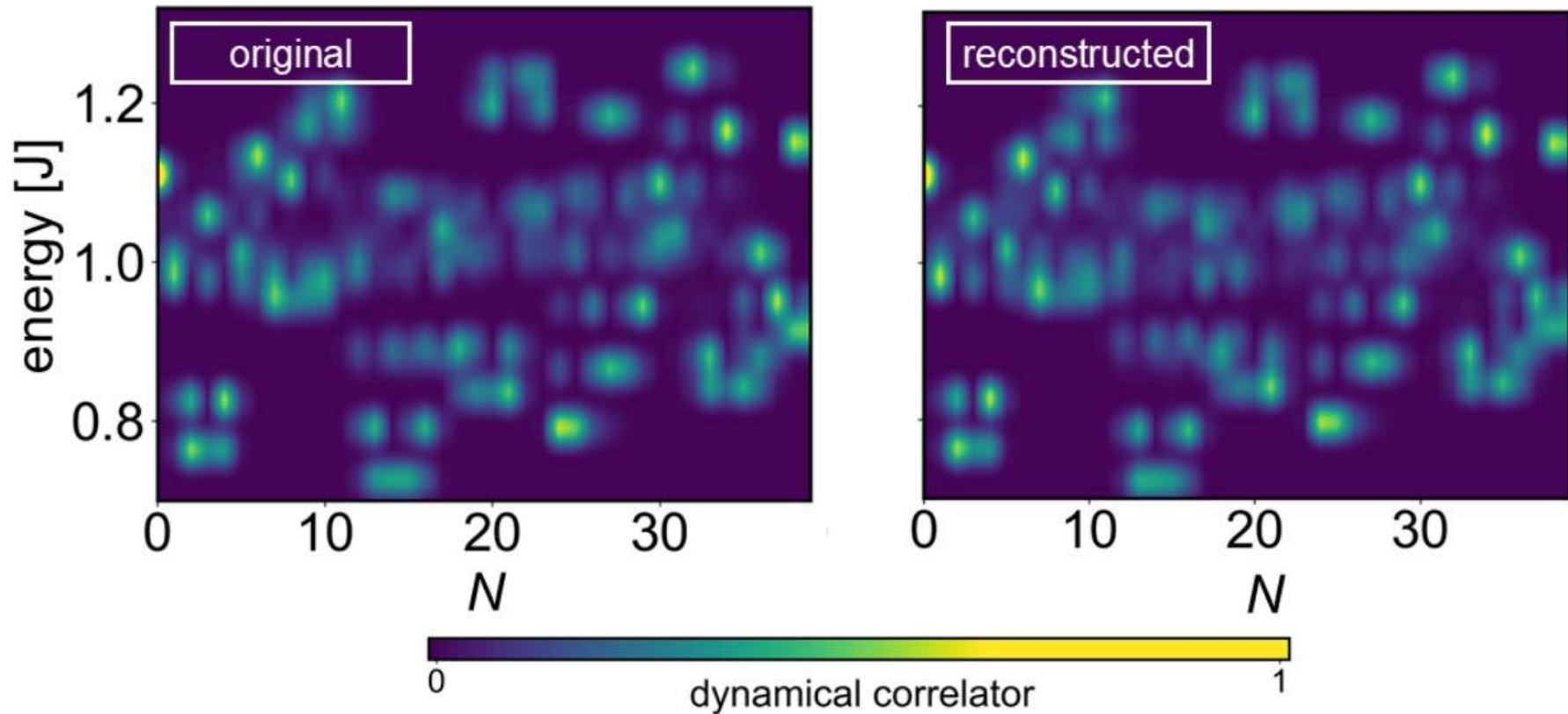


Intramolecular
Intermolecular

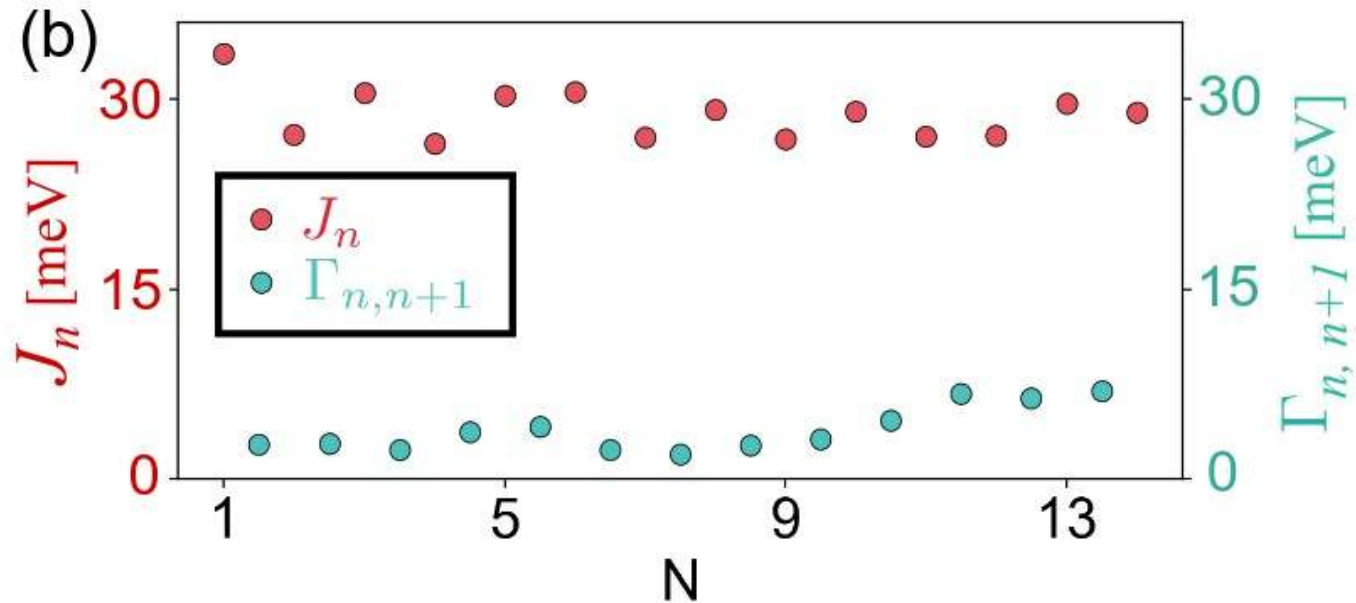


exchange variations can be predicted robustly
exchange variations are more subtle to predict

Reconstructing a quantum magnet Hamiltonian

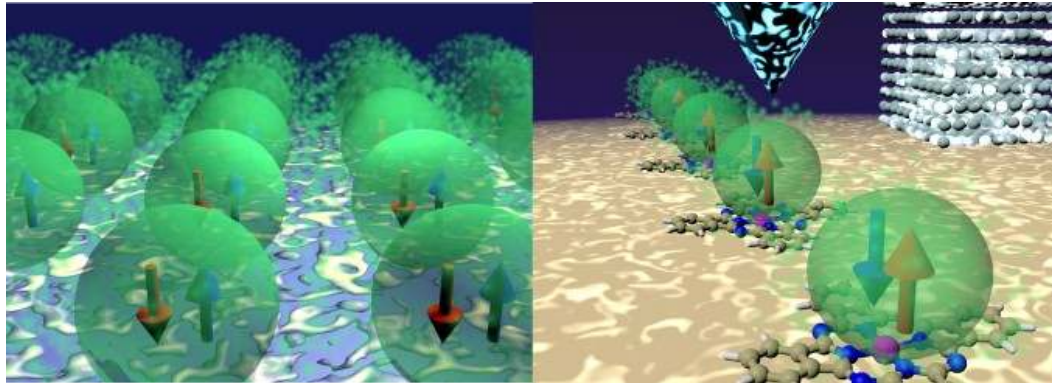


Machine learning triplons experimentally



Take home

STM sensing allows learning the Hamiltonian of complex collective quantum magnets, by leveraging machine learning trained with quantum many-body methods



Phys. Rev. Lett. 131, 086701 (2023)
arXiv:2504.20711 (2025)

